

## RESEARCH ARTICLE

# A Fast Startup 38.4-MHz Crystal Oscillator Achieving 99-nJ Startup Energy With Adaptive Chirping

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**ABSTRACT** Shortening the startup time of a crystal oscillator is critical in improving the energy efficiency of the Internet of Things (IoT) sensor nodes. This article presents an investigative study on the application of adaptive chirping using zero-phase cross-detection to reduce the startup time of a crystal oscillator that is robust to voltage, temperature and even process variations without any costly trimming. Post-layout simulations on a 38.4 MHz crystal resonator with 1.0 V supply and 65-nm CMOS process confirms the feasibility of this approach and its ability to effectively reduce the startup time by correcting both phase and frequency mismatches. The results showcase the promising potential of zero-phase adaptive chirp as a viable variation-tolerant technique allowing for a high frequency mismatch tolerance of  $1.1 \times 10^6$  ppm- $\Delta f$ .

**INDEX TERMS** CMOS, Internet-of-Things (IoT), crystal oscillator (XO), startup time, chirp injection, motional current, adaptive chirp injection, energy injection.

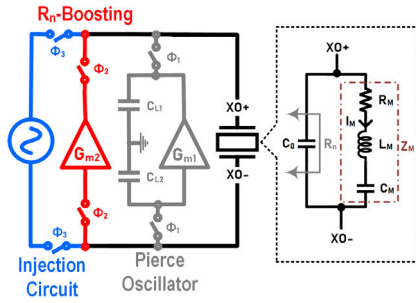
## I. INTRODUCTION

In the era of wireless sensor networks and the Internet of Things (IoT), wireless sensor nodes play a vital role by efficiently collecting and transmitting data for analysis purposes. These battery-powered nodes for extended operation are deployed where battery replacement is inconvenient and costly, making power optimization a momentous concern. Hence, there is a continued scientific interest in power optimization at the device and network levels [1].

The power-hungry blocks within the wireless sensor nodes, such as low-noise amplifiers, phase-locked loops, data converters, etc., will intermittently operate to reduce the average power consumption. While such duty-cycling

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technique can extend the battery life, it requires the circuit block to wake up instantly to avert latency and reduce startup energy loss. One crucial circuit block that requires a long startup time ( $T_{\text{start}}$ ) is the crystal oscillator (XO). The XO can provide a stable and precise frequency reference for the wireless sensor nodes. However, due to the high quality factor of the quartz crystal, it takes milliseconds for an MHz-range XO to reach the steady state after enabling the XO. This long  $T_{\text{start}}$  unavoidably constrains the latency of sensor nodes and correlates with increased depletion of energy. Reference [2] demonstrates that this startup loss accounts for a significant portion, approximately 42%, at each operational cycle of a transceiver at heavy duty cycling of 0.1%. Consequently, achieving additional power savings through heavy duty cycling necessitates a reduction in  $T_{\text{start}}$ .



**FIGURE 1.** Crystal resonator connected in Pierce oscillator (grey), negative resistance boosting startup (red), and Injection startup (blue).

For a traditional XO in Pierce oscillator configuration (Fig. 1), using the linear oscillator model featuring a crystal resonator, the motional current in the core of the quartz crystal ( $i_M$ ) grows exponentially – provided that the amplifier can promote the growth of oscillation – and is approximated as [3]:

$$i_M(t) \propto I_M(t=0) \cdot e^{\alpha t}, \quad (1)$$

$$\alpha \approx \frac{|R_n| - R_m}{2L_m}, \quad (2)$$

where  $R_m$  is the crystal's motional resistance,  $R_n$  is the negative resistance seen by the crystal's motional branch, and  $L_m$  is the motional inductance of the crystal. Then, we can obtain the  $T_{start}$  as:

$$T_{start} = \frac{Q}{\pi f_0} \cdot \frac{1}{\alpha} \cdot \ln\left(\frac{0.9 I_M(t=t_{ss})}{I_M(t=0)}\right) \quad (3)$$

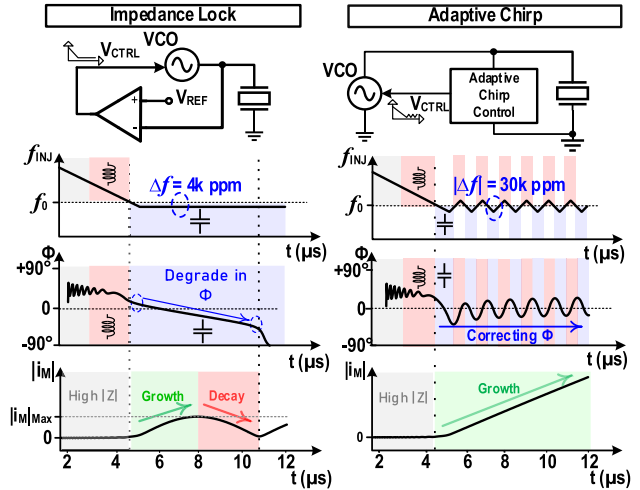
where  $Q$  is the quality factor of the quartz crystal and  $f_0$  is the resonance frequency of the crystal.

There are two main methods to reduce  $T_{start}$ . The first one is  $R_n$ -boosting [2], [3], [4], [5], [6]. As presented in (1) and (3), boosting  $R_n$  increases the growing rate of  $i_M$ , and thus shortens  $T_{start}$ . The second method is energy injection. Its underlying principle is to inject energy into the crystal core using an auxiliary source (e.g., RC oscillator) and raise the initial  $i_M$  of the crystal before connecting the crystal with the core amplifier to reduce the time for the  $i_M$  to build up from thermal noise. However, the phase of the injecting source must be in-phase with the  $i_M$  throughout the injection to effectively raise  $i_M$ . Due to the high  $Q$  of the crystal, it requires a typical injection duration above hundreds of cycles to raise, which in turns requires an auxiliary oscillator with a frequency close to  $f_0$ . Hence, recent publications on energy injection techniques focus on safeguarding the robustness of the XO's startup amid the process, temperature, and voltage (PVT) variations [7], [8], [9], [10], [11]. Although techniques like two-step injection [8] have strong tolerance against VT, they require costly post-fabrication calibration, such as trimming, to compensate for the process variation.

Chirp injection on the other hand guarantees the performance of the startup by sweeping the injection frequency to cover the frequency deviation of the auxiliary oscillator due to voltage, temperature and even process variation. It does not

**TABLE 1.** Crystal resonator and injection parameters used for analysis.

$f_0$	$L_m$	$C_m$	$R_m$	$C_0$	$V_{INJ}$
38.4 MHz	2.865m H	6f F	60 $\Omega$	2p F	1 V



**FIGURE 2.** Comparison of operation between Impedance lock and the proposed adaptive chirp (fixed peak-to-peak frequency toggling of 30,000 ppm around the resonance) in terms of frequency, phase and motional current envelop profiles.

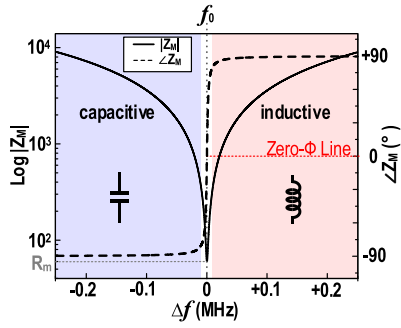
require specific trimming on the startup circuit to guarantee the operation, thereby significantly reducing the manufacturing and operating costs of the XO. A drawback of the typical chirp injection method is low energy injection efficiency; since the power is spread to a wide frequency band, the energy injected into the crystal core is low compared with constant frequency injection [7] or dithering [4], [11].

To this end, this article proposed a frequency and phase self-correcting energy injection method to expedite the XO startup that is robust against PVT variations. Dubbed Zero-Phase Adaptive Chirp (ZAC), it dynamically corrects the frequency and phase of the injection source by detecting the phase difference between the injection source and  $i_M$ . Simulated in the CMOS 65-nm process, the XO achieves a startup time of 170  $\mu s$ , with a startup energy of 99 nJ. It tolerates a frequency deviation of +110% (from nominal value).

The rest of this article is as follows. Section II outlines the conceptual utilization of adaptive chirping to continuously correct the injection phase. Section III elaborates on the circuit implementation of adaptive chirp using zero-phase cross-detection. The circuit architecture of proposed startup technique is described in section IV. Section V presents the post-layout simulation results and benchmarks this work with the state-of-the-art. Finally, Section VI concludes the article.

## II. FAST-STARTUP XO USING ADAPTIVE CHIRP

Crystal resonator is remarkable to offer an exceptionally high  $Q$  for excellent spectral purity. Yet, a drawback to this high  $Q$  is its slow startup behaviour. For a fast startup solution using



**FIGURE 3.** Frequency response of motional branch ( $Z_M$ ) of crystal resonator, exhibiting both capacitive ( $\Delta f < 0$ ) and inductive ( $\Delta f > 0$ ) characteristics.

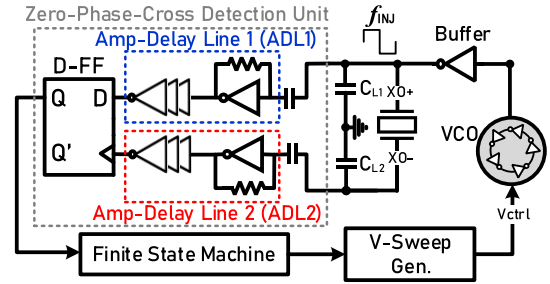
energy injection, a stringent requirement on the frequency mismatch between the resonance and injection frequency to promote robust growth on the  $i_M$  [12] is necessary. In light of this, PVT-tolerant energy injection techniques such as chirp injection are attractive as they allow the auxiliary injection circuit to inject energy at resonance frequency to increase  $i_M$  and reduce  $T_{start}$  amid PVT variation without any trimming after manufacturing. Yet, as the injection energy is distributed to a wide band for covering PVT variation, without feedback to lock the auxiliary oscillator, the energy delivered to the crystal is limited where the XO needs additional time to reach the steady state after the injection. Alternatively, resonance searching techniques such as impedance-guided injection [7] have been proposed to lock the injection frequency near the resonance frequency.

Herein, based on impedance-guided injection, we proposed adaptive chirping for fast XO start-up (Fig. 2). The operation principle is by comparing the frequency and phase of the chirping signal to that of the motional current from the crystal. In [7], the injection unit is locked at a low crystal impedance, where the exact impedance magnitude is determined by reference voltage,  $V_{REF}$ , and assumed to lock at  $f_0$ . However, to ensure an effective injection, the frequency mismatch after locking must be below 530 ppm to establish a maximum motional current for given crystal parameters in Table 1,  $C_{L1} = C_{L2} = 14$  pF, and  $V_{XO+} = 0.9$  V as obtained below [12]:

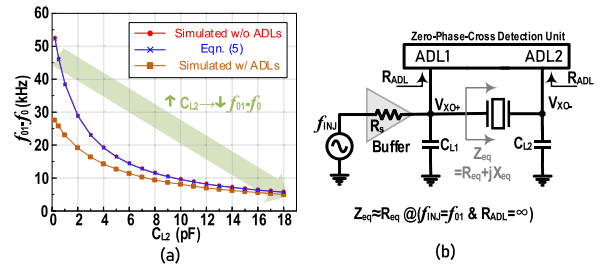
$$\frac{\Delta f}{f_0} = \mp \frac{4V_{INJ}/\pi}{V_{XO+}} \cdot \frac{C_M}{(C_{L1} + C_{L2})/2 + C_0}, \quad (4)$$

where  $V_{XO+}$  is the steady-state oscillating amplitude,  $V_{INJ}$  is the amplitude of the injection signal,  $C_{L1}$  and  $C_{L2}$  are the load capacitances to two ends of the crystal, and  $C_0$  is crystal's shunt capacitance. Locking within 530 ppm of  $f_0$  requires a delicate design consideration accounting in the chirping sweep rate, VCO noise, comparator offset,  $V_{REF}$ , and delay in resonance detection under PVT variations. Consequently, due to  $\Delta f$  and uncorrected phase error ( $\phi$ ), the injection has to be disabled by a short duration when the maximum  $i_M$ ,  $|i_M|_{Max}$ , has been reached.

In light of this challenge, we proposed adaptive chirp to relax such strict design requirements and enable  $\phi$  correction



**FIGURE 4.** Proposed zero-phase XO startup system block diagram highlighting the zero-phase cross-detection unit.

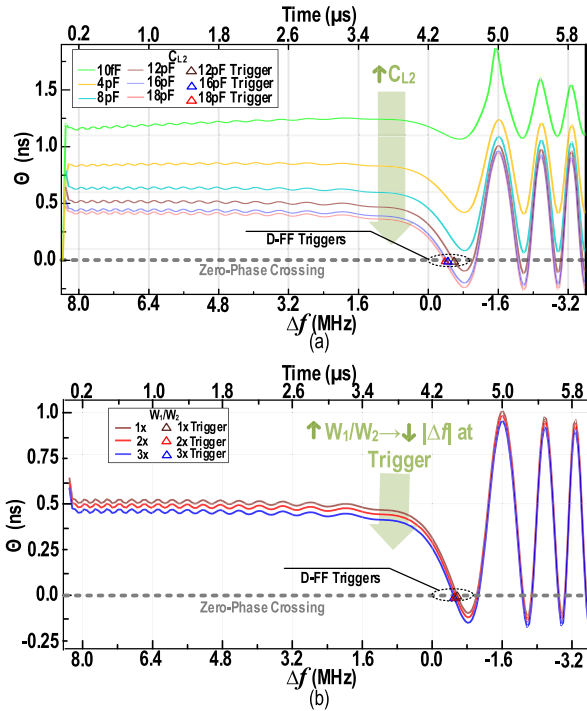


**FIGURE 5.** Panel (a) compares relative resonance frequencies between Equation (5) and simulations with/without ADL loading for different  $C_{L2}$  values. The simulations use the test-bench in panel (b) and parameters from Table 1, including Amp-Delay Lines from Figure 4.

under PVT variation. In principle, the chirping circuit sweeps the output signal's frequency ( $f_{INJ}$ ) up and down about  $f_0$  to correct the  $\phi$  accumulated due to injection with  $\Delta f$  (Fig. 2). Initially, the chirping signal starts its operation with a frequency significantly higher than  $f_0$  to guarantee it will cross  $f_0$  amid PVT variations. Due to the frequency mismatch, the impedance of the resonator ( $Z_m$ ) is determined mostly by the reactive components (either capacitive or inductive), which is higher than its impedance at  $f_0$  ( $\approx R_m$ ), as shown in Fig. 3. Hence, energy cannot be injected to the resonator initially and  $i_M$  maintains a weak magnitude. As  $f_{INJ}$  approaches  $f_0$  (entering low- $Z_m$  region), due to the positive  $\Delta f$  (initial  $f_{INJ} > f_0$ ), the resonator exhibits an inductive impedance, resulting in a pulling of  $\phi$  towards the positive. As  $f_{INJ}$  continues to decrease below  $f_0$ , the resonator transfers from inductive to capacitive. This transition leads to a pulling of  $\phi$  towards the negative. Therein, the  $\phi$  crosses the zero value. In principle, we can detect this zero-crossing point and invert the chirping direction (from down-chirping to up-chirping) to effectively correct the new  $\phi$ . Similarly, upon reaching the up-chirping operation, the  $\phi$  also crosses zero again, where we can again detect this point and invert the chirping direction. Hence, the circuit will generate an alternative up- and down-chirping sequence centred around  $f_0$  to enable a continuous correction of  $\phi$ . This coordinated approach facilitates sustained  $|i_M|$  growth.

### III. ZERO-PHASE LOCK AND ADAPTIVE CHIRP ANALYSIS

In this section, we will outline the implementation of both adaptive chirping and constant frequency lock injection



**FIGURE 6.** The  $\theta(t)$  versus  $\Delta f$  for different (a)  $C_{L2}$  and (b) delay line sizing. We set a chirping rate of 2 MHz/ $\mu$ s in the simulations.

techniques through the utilization of a Zero-phase Cross-detection unit, referred to as Zero-phase Adaptive Chirping (ZAC) and Zero-phase Lock (ZL) injections, respectively.

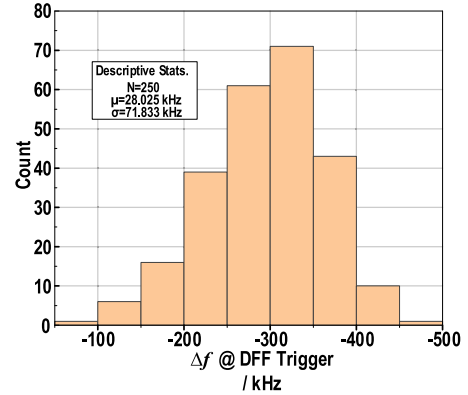
Fig. 4 shows the proposed system configuration and the Zero-phase Cross-detection unit to detect the  $\phi$  between  $i_M$  and  $V_{INJ}$  by using the transient phase difference between  $V_{XO+}$  and  $V_{XO-}$ . The proposed system has an added capacitance,  $C_{L2}$ , in series to the crystal, resulting the resonance frequency seen by the injector into the crystal to shift as given by:

$$f_{01}^2 \approx f_0^2 \frac{C_0 + C_{L2} + C_M}{C_0 + C_{L2}}, \quad (5)$$

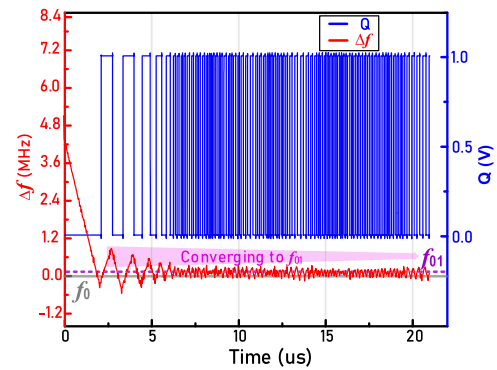
where  $f_{01}$  is the new resonance frequency seen into the crystal terminal in the presence of  $C_{L2}$  given no amp-delay line (ADL) loading (derived in Appendix). (5) evinces that  $f_{01}$  is independent of  $C_{L1}$  and  $R_s$  thereby it is independent of buffer loading and buffer driving strength respectively. Additionally, when  $C_{L2} \gg C_M$  then  $f_{01} \approx f_0$ . Utilizing the crystal parameters in Table 1, Fig. 5(a) shows that (5) gives close approximation of  $f_{01}$  for different  $C_{L2}$ .

Additionally, at  $f_{01}$ , since the combination of crystal and  $C_{L2}$  induce zero-phase and a grounded capacitor provides a  $-90^\circ$  phase shift, the steady-state phase  $\theta(\omega)$  at  $V_{XO-}$  with respect to  $V_{XO+}$  at  $f_{01}$  is  $-90^\circ$ , which can be validated by the relationship below:

$$V_{XO-}(j\omega) = V_{XO+} \cdot \frac{1/j\omega C_{L2}}{Z_{eq}(j\omega)}. \quad (6)$$



**FIGURE 7.** Monte-Carlo simulated  $\Delta f$  (at D-FF trigger) given  $C_{L2} = 14\text{pF}$ , and  $2x W_1/W_2$  under 2 MHz/ $\mu$ s down chirp.



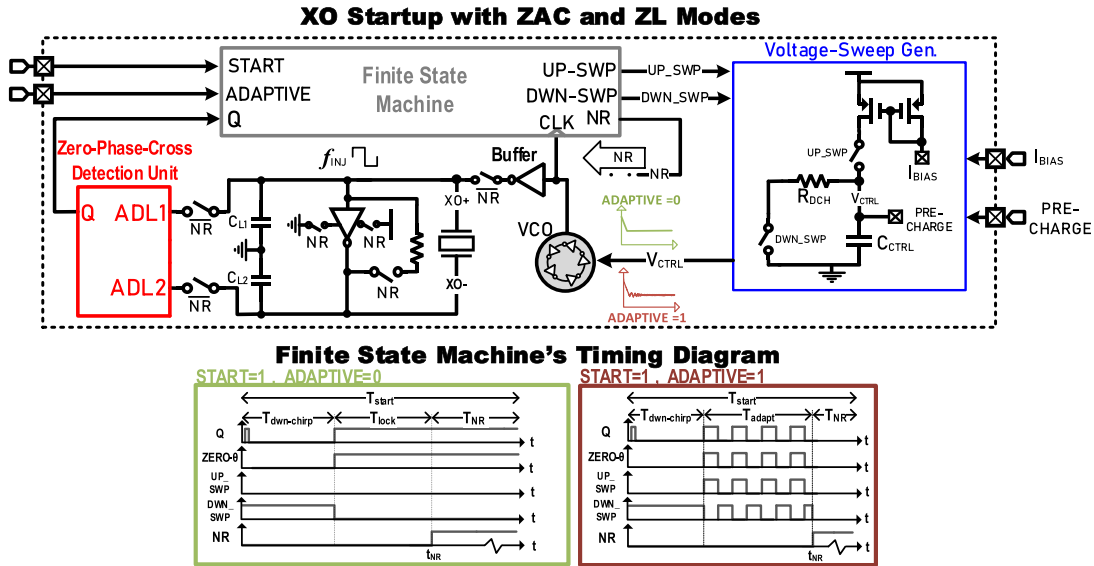
**FIGURE 8.** Frequency profile of the zero-phase adaptive chirp injection and the output from the D-FF.

From (6), we can obtain  $\theta(\omega)$  as:

$$\begin{aligned} \theta(\omega) &= \angle V_{XO-}(j\omega) - \angle V_{XO+}(j\omega), \\ \theta(\omega) &= \angle \frac{1}{j\omega C_{L2}} - \angle Z_{eq}(j\omega), \end{aligned} \quad (7)$$

where  $Z_{eq}$  is the impedance seen by the buffer looking into the crystal, and it is equivalent to the impedance of crystal added to the impedance of  $C_{L2}$  given the assumption of no Amp-Delay Line (ADL) loading. Given no ADL loading equation (7) shows that when  $\angle Z_{eq}(j\omega) = 0^\circ$ , which occurs when  $\omega \approx \omega_{01} (= 2\pi f_{01})$ ,  $\theta(\omega)$  becomes  $-90^\circ$ . Furthermore, the phase difference below and above  $\omega_{01}$  is  $-180^\circ$  and  $0^\circ$ , respectively.

The behavior mentioned above is valid in the steady state with no ADL loading, which forms a part of the transient response of the voltage difference between  $V_{XO+}$  and  $V_{XO-}$ . This article hence proposes the detection of  $f_{01}$  by exploiting the change in the transient phase,  $\theta(t)$ , when the chirping  $f_{INJ}$  crosses  $f_{01}$ . Note that the complete solution to such dynamic system's differential equation in providing transient phase includes both steady-state and transient components along with the ADL loading. The transient component with loading resists immediate aligning with the steady-state phase of  $-90^\circ$ ; however, it is nevertheless pulled toward the steady-state component, as depicted with Fig. 6's transient result with ADL loading. The transient phase crossing behaviour



**FIGURE 9.** System architecture of the proposed crystal oscillator with zero-phase lock and zero-phase adaptive chirp startup modes.

when chirping through  $f_{01}$  is utilized to represent  $f_0$  as opposed to low impedance thresholds in [7] and [9]. This behaviour is robust in detecting  $f_0$ , as it involves a single phase crossing point rather than a range of low voltage values. Moreover, such discrete crossing values can be exploited to implement adaptive chirping.

We utilized a zero-phase cross-detection unit to detect when the transient phase becomes zero. In principle as only the phases polarity of the signals is relevant, we can employ D-flip flop (D-FF) to detect the sequence of the transitions and thus the phase polarity. To guarantee the amplitudes of signals are capable in driving the logic circuits under PVT variations, we utilized a self-biased AC-coupled amplifier to amplify the signal swings, facilitating reliable triggering of the D-FF.

Detection of  $f_0$  using zero-phase cross detection requires the design of  $C_{L2}$  and the relative delay difference between amplifier delay lines 1 and 2, as they will affect the zero-crossing point. To ensure zero-crossing with tolerance to D-FF's metastability and PVT variations, given crystal resonator parameters given in Table 1,  $C_{L2}$  of 14 pF is selected as the transient phase goes well below zero-phase. Furthermore, delay difference between the lines is adjusted via width ratio ( $W_1/W_2$ ) and it is sized as 2x by a trade-off between consistent PVT operation and reduced footprint. Indeed, we utilized these two parameters to achieve coarse and fine-tuning on  $\theta(t)$  such that  $\theta(t) = 0^\circ$  when  $\Delta f = 0$  [Fig. 6(a) and (b)]. The Monte-Carlo simulated  $\Delta f$  (at D-FF trigger) in Fig. 7 proves zero-phase cross-detection unit's core functionality robustness against process variation.

The ZL technique can be realized by locking the injection frequency from the oscillator once the D-FF's output becomes high, corresponding to the first zero-phase cross-detection. This locking makes ZL, just like other resonance lock injections, slightly process sensitive. Notably, when compared to

the analogue based resonance searching technique in [7], which utilized a comparator, the proposed zero-phase searching promises a superior power consumption due to its digital implementation.

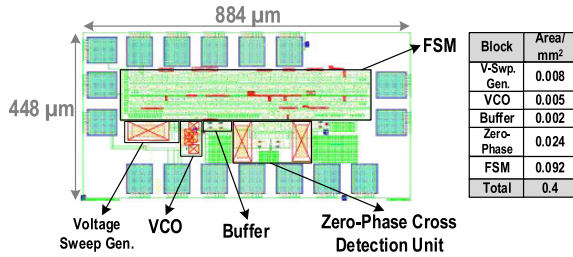
Similarly, the ZAC can be realized by inverting the chirping direction upon inversion of the D-FF's output, Q. Fig. 8 illustrates the operation of zero-phase cross-detection for ZAC showcasing how upon adaptive inversions of Q the frequency converges about  $f_{01} (\approx f_0)$ . This results in a lower frequency mismatch error compared to ZL. The adaptive frequency correction property also reduces the impact of variations on the injected frequency and crystal excitation. Additionally, the proposed adaptive injection provides phase correction, further enhancing variation tolerance (where phase correction is further discussed in Section V).

#### IV. CIRCUIT ARCHITECTURE

Fig. 9 illustrates the circuit architecture of the crystal oscillator with ZL and ZAC for fast startup. It includes primarily a digital CMOS finite state machine, voltage sweep generator, voltage-controlled ring oscillator, zero-phase cross-detection unit with digital CMOS D-FF, a weak buffer (implemented as a CMOS inverter) and self-biased inverting amplifier for Pierce oscillator configuration [13]. Whereby  $C_{L1} = C_{L2}$  considering a conventional Pierce oscillator. The startup sequence and operation of ZL and ZAC injection (Fig. 9) is as follows:

1) The initial voltage of the sweep voltage generator,  $V_{CTRL}$ , is set by charging its capacitor,  $C_{CTRL}$ , using the PRECHARGE port.

2) When the startup is enabled, the  $V_{CTRL}$  starts to drop according to  $\Delta V = V_{CTRL}(t = 0) \times \exp(-\Delta t / R_{DCH} C_T)$ , where  $C_T$  represents the aggregate of  $C_{CTRL}$ , parasitic capacitances, and the input capacitance of the VCO, with  $C_{CTRL}$



**FIGURE 10.** Layout of XO with the proposed ZL and ZAC fast startup injection techniques.

being the predominant capacitance contributor. Hence, the VCO's frequency starts to decrease.

3) The zero-phase cross-detection unit detects the relative phase between  $V_{XO+}$  and  $V_{XO-}$ . To overcome the error due to the random noise, the FSM only accepts transition when the results from 3 periods are identical. The FSM uses the VCO's clock signal to count this duration using a digital counter.

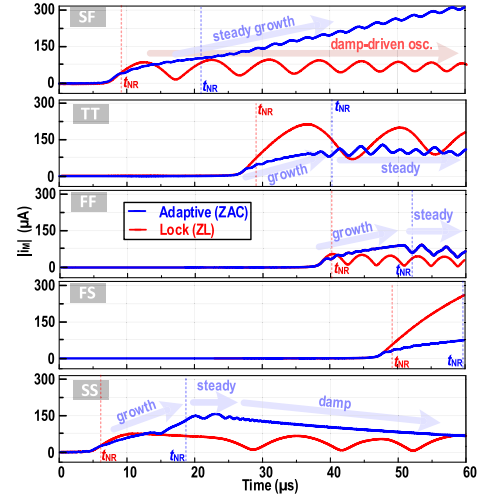
4) When the  $f_{INJ}$  exceeds  $f_{01}$ , the zero-phase cross-detection unit detect this zero-phase point.

5) For ZL injection, when the crossing point is detected (zero- $\theta$  becomes 1 V), the DWN\_SWP state turns to 0 V and consequently locks the injection frequency. The injection then continues until it reaches a preset duration,  $T_{lock}$ . For the ZAC injection, upon inversion of zero- $\theta$ , the DWN\_SWP and UP\_SWP states become 0 V and 1 V, respectively, leading to the charging of  $C_{CTRL}$  according to  $\Delta V = I_{BIAS} \times \Delta t / C_T$ . This results in an up-sweep of  $V_{CTRL}$  and an up-chirp of  $f_{INJ}$ . Another inversion of zero- $\theta$  inverts UP\_SWP and DWN\_SWP states again, causing a down-chirp of  $f_{INJ}$ . This process continues for a preset injection duration,  $T_{adapt}$ . After the injection durations for both ZL and ZAC injections, the NR state changes to 1 V at time  $t_{NR}$  (where this value will be discussed in Section V).

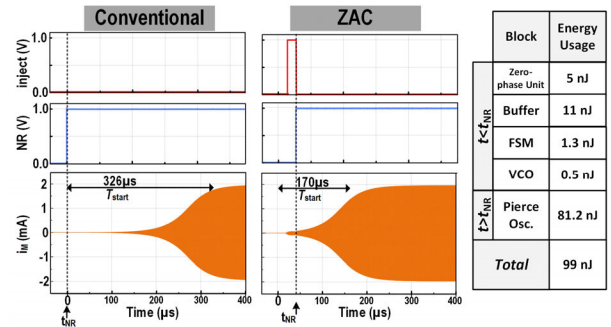
6) With NR state at 1 V, the injecting unit and zero-phase cross-detection unit are disabled, while the core amplifier unit, made from a self-biased inverter, is in place to sustain the oscillation in the steady state.

## V. POST-LAYOUT SIMULATION RESULTS

To prove the core functionality of the proposed fast startup XO, we implemented and simulated with post-layout parasitics in the 65-nm CMOS process. Fig. 10 shows the system layout which occupies a total area of 0.4 mm<sup>2</sup>. The capacitive loads of 14p F and voltage supply of 1 V is used for the oscillator. Furthermore, we have selected a crystal resonator with  $f_0$  of 38.4 MHz driven by a small buffer ( $V_{XO+} = 150$  mV peak-to-peak). To account for the PVT variations, the initial  $V_{CTRL}$  is set such that VCO's initial frequency is above  $f_0$  in all corners using the same  $V_{CTRL}$  value. This is achieved by conducting a parametric sweep simulation of  $V_{CTRL}$  against VCO's frequency and ensuring that the resulting frequency for the selected  $V_{CTRL}$  is higher than  $f_0$  across all PVT corners. Furthermore, the ring-VCO exhibits the fastest and slowest frequency down-sweep rates of 1.6 MHz/ $\mu$ s and 0.8 MHz/ $\mu$ s for the FF and SS corners, respectively.



**FIGURE 11.** The growth of  $|i_M|$  using ZAC and ZL. We set  $t_{NR}$  corresponding to 519 cycles and 70 cycles after the first zero-phase detection respectively.



**FIGURE 12.** Oscillation signal with a 38.4 MHz resonator, displayed without (left) and with (right) the proposed fast startup. The accompanying table indicates energy consumption with the proposed startup.

Fig. 11 shows  $|i_M|$  for all process corners to highlight the phase-error correction property of ZAC relative to ZL. While ZL only exhibits damped driven oscillation, ZAC shows continuous growth (at SF), growth and sustain (at TT and FF) and growth, sustain and gradual damping (at SS) corresponding to broad effectiveness of adaptive chirp. Leveraging the closed-loop property and aiming to enhance the average  $|i_M|$  across process variations, the selection of  $t_{NR}$  for ZAC and ZL, as inferred from studying Fig. 11, is 519 and 70 cycles respectively. These values equate to approximately 13.5  $\mu$ s and 1.8  $\mu$ s, based on a 38.4 MHz clock, following the initial zero-phase cross-detection. The extended injection duration of ZAC relative to ZL demonstrates its higher capacity for process variation tolerance. On the other hand, open loop injections, such as chirp injections, have to preset  $t_{NR}$  with respect to the start time. The startup time of the XO with ZAC and without startup aid is shown in Fig. 12. With ZAC, the startup time reduces from 326  $\mu$ s to 170  $\mu$ s. Furthermore, ZAC has reduced total startup energy from 152 nJ to 99 nJ. Whereby ZAC's energy before  $t_{NR}$  constitutes the consumptions in FSM (1.3 nJ), zero-phase cross detection (5 nJ), buffer (11 nJ), VCO (0.5 nJ).

TABLE 2. Performance summary and comparison with prior art.

Technique	JSSC'19 [8]	ISSCC'23 [9]	JSSC'22 [7]	JSSC'16 [10]	Proposed #	
	2-step PLL Injection	Automatic Phase-Error Correction	IGCI + Boosted $R_n$	Chirp + Boosted $R_n$	Zero-Phase Lock	Zero-Phase Adaptive Chirp
Technology (nm)	65	40	22	180	65	
Resonator Frequency (MHz)	54	16	38.4	39.25	38.4	
Steady state Voltage/ Supply voltage (V)	0.5/1.0	0.25/1.0 <sup>†</sup>	0.8/1.0	1.4/1.5	0.95/1.0	
PVT Tolerant Injection	NO	NO	YES	YES	YES	YES
Frequency Mismatch Tolerance (ppm- $\Delta f$ )	$5 \times 10^3$	$10^4$	$4.7 \times 10^4$ <sup>†</sup>	$1.25 \times 10^6$ <sup>*</sup>	$1.1 \times 10^6$ <sup>*</sup>	
$\Delta T_{start}/T_{start}$ over Temp <sub>range</sub>	1.25%	4.5%	26% <sup>†</sup>	7%	3% <sup>x</sup>	3.8%
Load Cap (pF)	6	6	3.75	6	14	
$T_{start}$ ( $\mu s$ )	19	17.5	58	158	175	170
$E_{START}$ (nJ)	34.9	9.2	45.6	349	107	99
Steady State Power ( $\mu W$ )	198	84	800	181	450	
Temperature Range ( $^{\circ}C$ )	-40 to 85	-20 to 85	-40 to 40 <sup>†</sup>	-30 to 125	-20 to 80	

# Simulated Results.

<sup>†</sup> Value obtained from visual inspection of figures.

<sup>\*</sup> Value obtained for corner case with highest value.

<sup>x</sup> Achieved low variation at cost of lower  $|i_M(t_{NR})|$

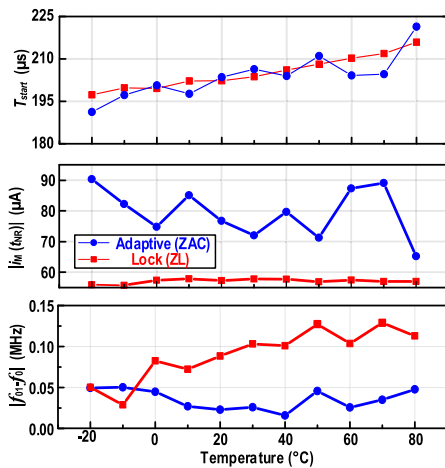


FIGURE 13. Simulated  $T_{start}$ ,  $|i_M(t_{NR})|$  and average frequency error against temperature variations.

The correlation between startup time and  $|i_M(t_{NR})|$ , as expressed in (3), informs the injection unit's efficiency. Fig. 13 illustrates  $|i_M(t_{NR})|$  under ZAC across  $-20$  to  $80$   $^{\circ}C$ , demonstrating a relative variation coefficient of 10.3%, with maximum and minimum  $|i_M(t_{NR})|$  of  $90$   $\mu A$  and  $71$   $\mu A$  at  $-20$   $^{\circ}C$  and  $50$   $^{\circ}C$ , respectively. These results correspond to a minimum and maximum  $T_{start}$  of  $191$   $\mu s$  and  $221$   $\mu s$  for  $-20$  and  $80$   $^{\circ}C$  with only a relative variation of 3.8%. In contrary, the constant frequency counterpart experiences a lower  $|i_M(t_{NR})|$  by 39% (from average of  $79$   $\mu A$  to  $57$   $\mu A$ ) and consequently a higher  $T_{start}$  from  $203$   $\mu s$  to  $205$   $\mu s$ , justifying ZAC's superior startup performance. Moreover, ZAC offers a lower average  $|f_{01} - f_0|$  by 157% with respect to ZL (from  $90$  kHz to  $35$  kHz). The improved frequency accuracy of ZAC across temperature variations is attributable to its longer convergence duration toward  $f_{01}$ . This extended convergence

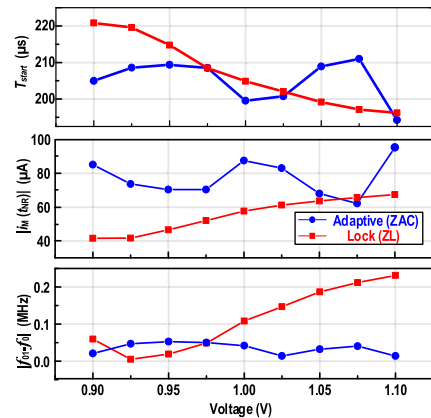


FIGURE 14. Simulated  $T_{start}$ ,  $|i_M(t_{NR})|$  and average frequency error against voltage variations.

duration significantly enhances ZAC's frequency detection variation tolerance by compensating for detection delays.

Fig. 14 illustrates the startup performance of the XO with ZAC and ZL against voltage variations ( $0.9$  V to  $1.1$  V). The XO with ZAC shows a 14% relative change in  $|i_M(t_{NR})|$  when the supply voltage varies by 10%, whereas that with ZL displays a higher relative variation of 18%. In addition to  $|i_M(t_{NR})|$ , the XO with ZAC also exhibits a lower average  $\Delta f$  of  $35$  kHz, in contrast to  $113$  kHz from that with ZL.

The  $|i_M(t_{NR})|$  response for the XOs with ZAC and ZL across process corners yields an average value of  $94$   $\mu A$  and  $54$   $\mu A$ , with  $219$   $\mu s$  and  $233$   $\mu s$  average startup times, respectively (Fig. 15); the fastest and slowest startup of the XO with ZAC are at the SF and SS corners ( $178$   $\mu s$  and  $308$   $\mu s$ ). At FS despite starting with the highest initial frequency of  $80$  MHz (corresponding to  $1.1 \times 10^6$  ppm- $\Delta f$ ) ZAC retains its correction property achieving  $221$   $\mu s$ .

Table 2 summarizes and compares the performance with prior architectures. The proposed adaptive zero-phase

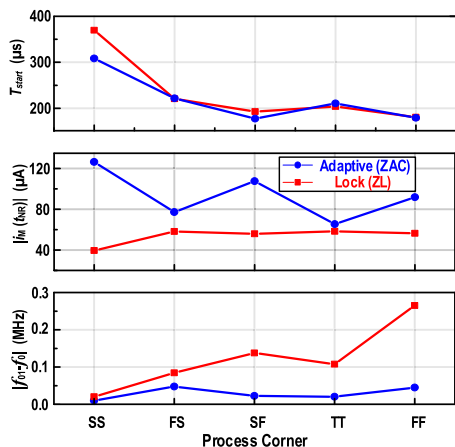


FIGURE 15. Simulated  $T_{start}$ ,  $I_M$  ( $f_{NR}$ ) and average frequency error against process corners.

injection technique supplements the PVT tolerant injection XO startup methods [7] and [10], enhancing them by providing an exceptional frequency mismatch tolerance of  $1.1 \times 10^6$  ppm- $\Delta f$ . While the proposed technique falls 12% short of the frequency mismatch tolerance achieved by the technique cited in [10], it compensates with superior  $3.5\times$  lower startup energy consumption ( $E_s$ ). The robustness of [10] and the proposed injection are attributed to the inherent chirping mechanism used. IGCI [7], also employing a chirping mechanism, enables quicker PVT-tolerant startup with respect to ZAC by utilizing a reduced  $C_L$  and lower steady-state voltages. However, this approach would entail an inherently higher oscillator phase noise cost [14], whereas the proposed design could demonstrate a superior performance. Moreover, owing to ZAC’s correction property, which mitigates the error induced by the resonance frequency detector, the startup variation across the temperature range is  $7\times$  superior to that of IGCI’s. Regardless of the higher power consumption and longer startup relative to other techniques, this investigative paper proposed and verified ZAC’s  $\Delta f$  and  $\phi$  correcting injection which showcases potential in PVT robust startup time reduction. For full exploitation of ZAC with lower steady-state power consumption and higher resonator excitation, optimization for a system with lower bias current Pierce oscillator and a higher injection voltage swing is planned respectively. For a given supply voltage, enhancing the injection voltage swing can be achieved by reducing the buffer output impedance,  $R_s$ , or by detaching  $C_{L1}$  during injection. Additionally, the verified operation with low injection voltage shows ZAC’s prospect in ultra-low voltage startup systems. This potential is further explored in future research.

VI. CONCLUSION

In this paper, we utilized chirp injection to enhance the startup performance of crystal oscillators. Specifically, we employed adaptive chirping to cover the frequency deviations caused by PVT variations without any specialized trimming of the startup circuit. This approach offers significant cost

savings in both manufacturing and operation. It provides a state-of-the-art startup time variation of merely 3.8% against temperature fluctuations, achieved without the need for trimming circuitry that is used for process compensation. The reduction in chip area, achieved by eliminating the need for trimming, along with the elimination of associated labor costs, leads to the savings in manufacturing expenses. Compared with conventional chirp injection methods that tend to have low energy injection efficiency due to their wide frequency band coverage, the proposed energy injection technique, zero-phase adaptive chirping, accelerates the crystal oscillator and also mitigates the impact of PVT variations. Post-layout simulations conducted in CMOS 65-nm process demonstrates the technique achieves a startup time of  $170 \mu s$  with a startup energy consumption of 99 nJ. Moreover, the injection oscillator for energy injection can withstand a frequency deviation range of  $+110\%$  from the nominal value, supporting the operation under different extreme operating corners. This innovative approach represents a promising advancement in crystal oscillator startup, offering improved robustness to PVT variations.

APPENDIX DERIVATION OF (5)

With reference to Fig. 5(b), the impedance seen by buffer looking into the crystal, with ignoring the loading by Amp-Delay Lines, is

$$Z_{eq} = Z_3 + 1/j\omega C_{L2},$$

where  $Z_3$  is the impedance of the crystal, and it is given as

$$Z_3 = 1/j\omega C_0 || (R_M + 1/j\omega C_M + j\omega L_M),$$

$$Z_3 = \frac{(aR_M + bd) + j(bR_M - ac)}{R_M^2 + c^2},$$

$$Z_3 = R_3 + jX_3,$$

where

$$a = \frac{\omega L_M - 1/\omega C_M}{\omega C_0},$$

$$b = -R_M/\omega C_0,$$

and

$$c = \omega L_M - 1/\omega C_M - 1/\omega C_0.$$

At  $f_{01}$  we have

$$Z_{eq} = R_{eq},$$

since at this frequency

$$X_{eq} = X_3 - 1/2\pi f_{01} C_{L2} = 0. \tag{1A}$$

Considering negligible  $R_m (= 0)$ , using (1A)  $f_{01}$  is derived to be

$$f_{01}^2 = \frac{C_{L2} + C_0 + C_M}{C_0 + C_{L2}} \left( \frac{1}{L_M C_M} \right). \tag{2A}$$

Using crystal’s series resonance frequency  $f_0 (= 1/2\pi \sqrt{L_M C_M})$ , (2A) gives  $f_{01}$  as (5).



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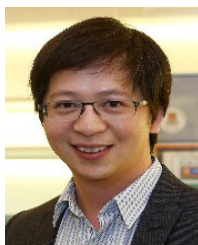
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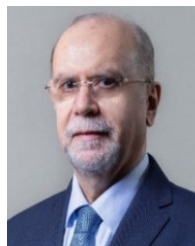
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