# Ultra-Broadband and Electro-Optical Tunable Absorption in Double-Walled Carbon Nanotubes

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*(Invited Paper)*

*Abstract***—Electro-optical modulators are critical elements in the rapidly developing data communication, optical interconnects, silicon-based photonic systems and terahertz technologies. The limited optoelectronic properties and complicated material growth in traditional semiconductors hinder the rapidly surging demand for modulator performance, energy efficiency, cost, etc. The emergence of two-dimensional materials and one-dimensional carbon nanotubes in recent decades has brought new opportunities with their tremendous selection degree of freedom for exceptional optoelectronic properties. In this article, we present ultra-broadband and electro-optical tunable absorption modulators by employing double-walled carbon nanotube films in a capacitor geometry, spanning the visible to terahertz spectra. The formation of supercapacitors around the ionic gel electrolyte and carbon nanotube film interfaces accounts for the large carrier transition and optical conductivity change, which behaves a thickness dependent electroabsorption dynamics. Our findings not only broaden the understanding of low-dimensional material applications in electro-optics but also pave the way for future developments in high-performance broadband modulators.**

*Index Terms***—One-dimensional materials, double-walled carbon nanotube, electro-optical modulator.**

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## I. INTRODUCTION

PTICAL modulators, a critical component to manipulate the physical parameters of light, are indispensable in photonic and optoelectronic systems. The most developed optical modulators adopt electro-optical approach to convert electrical signal to optical signal, which thus can be used in telecommunication, optical computing, optical signal processing [\[1\],](#page-5-0) [\[2\],](#page-5-0) [\[3\],](#page-5-0) etc. These technologies are based on traditional semiconductors owning advanced optoelectronic behaviors. Representatives are III-V group compounds (e.g., GaAs, InP), lithium niobate, and silicon-based photonic platforms [\[4\],](#page-5-0) [\[5\],](#page-5-0) [\[6\].](#page-5-0) The increasing demand for high-efficiency, fast speed, low energy consumption and miniaturized modulators has directed huge research efforts in silicon-based photonic platforms because of their merits for wafer-scale integration and compatibility to the mature complementary metal-oxide semiconductor technology [\[7\].](#page-5-0) According to the parameters being affected by the electrical potential, one can modulate intensity (amplitude), phase or polarization of the light passing through the active material in the device [\[8\].](#page-5-0) Intensity modulation, a fundamental regime usually accomplished by changing the light absorption of a component in photonic system, has been an essential need for many applications. The realization of absorption modulation relies on modifying the charge carrier density of the sensitive material, which is often utilized as electrodes in the modulator. Although were drastically developed, the electro-optical absorption modulation with traditional bulk semiconductors are still restricted in operation bandwidth, and need large driving voltage, which is attributed to the large bandgap and low doping efficiency. The emergence of a number of two-dimensional (2D) nanomaterials, graphene [\[9\],](#page-5-0) [\[10\],](#page-5-0) [\[11\]](#page-5-0) and its analogue (e.g., transition metal dichalcogenides of  $MoS_2$ ,  $WS_2$ ,  $MoSe_2$ ,  $WSe_2$ ) [\[12\],](#page-5-0) [\[13\],](#page-5-0) [\[14\],](#page-5-0) [\[15\]](#page-5-0) finds alternative option for improved performance. The combination of gapless graphene and gap tunable transition metal dichalcogenides greatly extends the absorption range, but the insufficiency of density of states around the Fermi level confines their gate dependent doping concentration, and gives weak absorption effect along ultrathin atomic layer. Effective strategies (e.g., waveguide integration, heterostructures, plasmon oscillator) [\[16\],](#page-5-0) [\[17\],](#page-5-0) [\[18\]](#page-5-0) have been proposed to enhance the interaction intensity, but in the cost of limited bandwidth and complicated fabrication process.

<span id="page-1-0"></span>In quasi one-dimensional carbon nanotube, optoelectronic properties and band structures change dramatically as a result of quantum confinement and curvature, in comparison to 2D nanomaterials. The valence or conduction bands are composed of discrete levels called subbands, and have different azimuthal quantum members. Carrier transition of external excitation can be interband or intersubband. In addition, the absorption spectrum variation under electrical field is significantly enhanced by excitonic effects [\[19\].](#page-5-0) Compared to bulk semiconductors and graphene, the typical absorption spectrum of single walled carbon nanotubes (SWCNTs) behaves multiple characteristic peaks corresponding to the transition between van Hove singularities in the density of states  $[20]$ . Benefit from the development in material synthesis, SWCNTs have been widely studied as an active medium to acquire optical modulation from the near infrared to mid-infrared region [\[21\],](#page-5-0) [\[22\].](#page-5-0) Investigations also reveal the existence of plasmon conductivity peaks in metallic and doped semiconducting carbon nanotubes at the THz frequency range [\[23\],](#page-5-0) indicating that carbon nanotube is a promising building block for THz devices. In comparison with the massive research attention attracted by SWCNTs, the investigation about structure and properties of double walled carbon nanotubes (DWCNTs) is a relatively less touched area until the synthesis and separation of its high purity sample [\[24\],](#page-5-0) [\[25\].](#page-5-0) In a simplified model, the optical absorption spectrum of DWCNTs can be described as a superposition of the spectra of its inner and outer tubes, and a small shift of absorption peaks caused by van der Waals interaction between the two adjacent walls [\[26\],](#page-5-0) [\[27\],](#page-5-0) [\[28\],](#page-5-0) [\[29\].](#page-5-0) Experimental observation in our previous work shows that the absorption peaks contributed by SWCNTs are significantly suppressed if the catalyst concentration during the synthesis can be controlled appropriately [\[30\].](#page-5-0) Notably that the outer walls of DWCNTs are quite large in diameter and have wide distribution, while the inner walls are in a comparable size of small individual SWCNTs. As a result, bandgaps of DWCNTs are much smaller than their single walled counterparts and the optical absorption spectrum of DWCNTs can be drastically broadened.

In this work, we present our recent demonstration of ultrabroadband and electro-optical tunable absorption in DWCNTs. A pair of DWCNT films are utilized as active electrodes, and isolated by a thin layer of ionic gel electrolyte to form a capacitor device. Electrical gating introduced carrier concentration reconfiguration on the nanotube film contributes to the absorption variation from the visible to mid-infrared. Drude-like conductivity change of the DWCNT films is in charge of the THz wave modulation. In particular, our result shows a thickness-dependent electro-optical absorption change over the entire spectrum range, giving alternative solution for modified efficiency in modulation strength.

# II. ELECTRO-OPTICAL ABSORPTION MODULATOR WITH CARBON NANOTUBE

## *A. Material Synthesis and Characterization*

The synthesis of DWCNTs was carried out with the floating catalyst chemical vapor deposition (FC-CVD) method. A



Fig. 1. (a) Schematic setup of the FC-CVD method for DWCNT synthesis and film deposition. (b) A large-scale DWCNT thin film collected on a membrane filter. (c) Raman spectroscopy of the synthesized DWCNTs film (excited by a 633 nm laser). (d) SEM and (e) HR-TEM images of the DWCNT film. (f) Transmittance of DWCNT films at 550 nm as a function of their sheet resistance.

schematic experimental setup is illustrated in Fig.  $1(a)$ . Gasphase precursors, including ferrocene (FeCp<sub>2</sub>), sulfur,  $CH<sub>4</sub>$  and carrier gases of  $N_2$  and  $H_2$  were introduced to generate DWCNTs in the reactor at 1100 °C. Aerosol-like DWCNTs were covalently constructed in gas-phase, and collected as large area continuous network by dry filtration from the outlet of the reactor with a membrane filter at room temperature, as shown in Fig. 1(b). The deposited film thickness can be flexibly tailored by setting the collection time. During the reaction process, it is possible to control the ratio between SWCNTs and DWCNTs by adjusting the concentration of the catalyst ( $FeCp<sub>2</sub>$ ). As the increase of catalyst concentration, the produced nanotubes change from SWCNTs in large bundles to DWCNTs in small bundles. In experiment, the concentration of catalyst was determined by controlling its evaporation temperature. When the temperature is over 35 °C (selected in this work), the Van Hove transition peaks in the absorption spectrum of the nanotubes disappear completely, indicating high-concentration DWCNTs are synthesized. Statistical analysis result in our previous report [\[30\]](#page-5-0) shows that only ∼5% are SWCNTs, and the others are DWCNTs and their bundles. For the DWCNTs, most of the tube length is from ∼1 *µ*m to∼40*µ*m, giving a mean length of∼20(±11)*µ*m. The identification of DWCNTs was carried out by Raman spectroscopy (Horiba LabRAM HR 800) under the excitation of a 633-nm continuous-wave laser. Fig.  $1(c)$  shows the Raman spectrum with radial breathing mode (RBM) at  $\sim$ 221.8 cm<sup>-1</sup>, disordered carbons caused vibration mode (D band) at  $\sim$ 1325 cm<sup>-1</sup> and tangential oscillation mode (G band) at  $\sim$ 1588 cm<sup>-1</sup>. The single Raman peak of the G band with a Lorentzian line-shape indicates that the excitation is from semiconducting nanotubes [\[31\].](#page-5-0) A further nanoscale investigation was conducted by scanning electron microscope (SEM, ZEISS Sigma VP, 1.0 kV) and high-resolution transmission electron microscope (HR-TEM, JEOL JEM-2200FS, 200 kV). Experimental results are shown in Fig. 1(d) and (e), respectively. The SEM morphology manifests homogenous nanotube networks, which is a result of random



Fig. 2. (a) Schematic of the capacitive modulator. A pair of DWCNT films are employed as parallel board electrodes and spaced by ionic gel electrolyte. Supercapacitive fields were formed by electron-cation and anion-hole concentration around the DWCNT films when gate is applied. (b) Optical absorption spectra of the pristine DWCNT films in different thicknesses corresponding to ∼80% (black curve) and 50% (red curve) transmittances at 550 nm wavelength. (c), (d). Pictures of the prepared devices without ionic gel electrolyte (c) and with ionic gel electrolyte (d).

orientation during the collection process. In Fig.  $1(e)$ , the HR-TEM image verifies the formation of two embedded nanotubes. The black particles on the nanotubes are from the residual catalyst of ferrocene. We measured the pristine film transmittance at 550 nm laser and corresponding sheet resistance with different samples; experimental data are plotted in Fig. [1\(f\).](#page-1-0) The fitted curve shows that the film transmittance increases exponentially against the sheet resistance. This is attributed to the increased carrier concentration when the film becomes thicker (and vice versa). Note that the sheet resistance can be less than 100  $\Omega$ /sq when the transmittance is approximately below 83%, which proves the high quality of the DWCNT films.

## *B. Modulator Fabrication*

The modulator device was configured in a capacitor geometry by placing two DWCNT thin films in parallel and isolated by ionic gel electrolyte (1-Ethyl-3-methylimidazolium bis(trifluoromethyl sulfonyl)imide, i.e., EMIM-TFSI). A schematic diagram of the device is shown in Fig.  $2(a)$ . Two DWCNT films  $(2 \times 2 \text{ cm}^2)$  were transferred to transparent polycarbonate boards. A pair of copper foil stripes were connected to the edge of each DWCNT thin film on the substrates for voltage application. Two polycarbonate substrates with the DWCNT film side were put together to form a 125  $\mu$ m spacer. A completed device was acquired after injecting the ionic gel solution. For comparison, two devices utilizing 80% and 50% transmittance DWCNT films at 550 nm wavelength were prepared. The film thicknesses were evaluated as ∼29 nm for the 80% transmittance and ∼80 nm for the 50% transmittance with atomic force microscopy. Fig.  $2(b)$  shows the optical absorption spectra of the selected DWCNT films (carried by Perkin-Elmer Lambda 950, UV-VIS-NIR spectrometer) in the wavelength range from 280 nm to 2100 nm. The absorption spectra exhibit smooth intensity distribution, which is a comprehensive result of two factors. One

is the wide outer tube diameter distribution ranging from ∼2 nm to ∼6 nm in our DWCNT films, as shown in the Supporting Information of our previous work [\[30\].](#page-5-0) The large outer walls of the DWCNTs form very small bandgaps. The other one comes from the excitonic transition energies of both the inner and outer walls when the outer tube is semiconducting type (i.e., the inner tube can be semiconducting or metallic). The combination of different charge transfer behaviors between the tubes also contributes to the broadband optical absorption of the DWCNTs. Fig.  $2(c)$ presents the two modulators utilizing DWCNT films with 80% (left) and 50% (right) transmittance at 550 nm wavelength before the injection of ionic gel electrolyte. The grey and dark grey areas are the DWCNT electrodes with different transparencies. The white membrane under the nanotube film is an isolation layer to prevent short-circuit, which turns transparent after injecting the ionic gel, as shown in Fig. 2(d).

# *C. Electro-Optical Modulation From the Visible to Infrared Light*

In silicon-based modulators, electrical field is generated between electrodes spaced by dielectric, this limits the formation of large concentration carrier doping in the active material because of the dielectric breakdown issue [\[32\].](#page-5-0) It is therefore challenging to realize modulation at high frequency range around near infrared and visible light. In comparison, ionic gel electrolyte with large electrochemical window has been demonstrated viable, alternative to extend the electroabsorption modulation range [\[33\].](#page-5-0) Cations and anions in the electrolyte offer a regime to form supercapacitor at the DWCNT film surfaces, which is confined to nanometer thickness. Together with the huge surface to volume ratio of DWCNT film, the strong localized field built by the supercapacitor enables storing drastically large density carrier at the film interfaces. This mechanism ensures the electro-optical modulation broadening to visible wavelengths under gate of only several volts. On the other hand, the interband absorption in the mid-infrared range can be Pauli blocked by natural doping to the nanotube films. This influence was suppressed by baking the device at 150 °C for 30 minutes before injecting the ionic gel, and the remaining fabrication was completed in a glove box filled with nitrogen gas to prevent ionic gel oxidation and device degradation. Fig. [3](#page-3-0) presents the change of transmission as a function of gate voltage between −2.1 and 2.1 V. The transmission spectra were normalized with the spectrum measured at open circuit voltage. All these transmission spectra have subtracted the interference from the substrate materials (i.e., a reference device without DWCNT electrodes), ensuring that the gate dependent transmission variation is solely contributed by the DWCNT films. Fig.  $3(a)$  shows a symmetric mapping of transmission change against voltage from the thinner sample (i.e., DWCNT film of 80% transmittance at 550 nm), but with the symmetry center slightly shifted to 0.3 V, due to charge impurities in the ionic gel and substrates. The maximum change in the transmission is about 6.6% at 2.1 V gate voltage, which presents a blue-shift as the increase of electrical gating, see Fig. [3\(b\).](#page-3-0) This phenomenon is attributed to the large interband transition around the Fermi energy when the gate voltage is increased.

<span id="page-3-0"></span>

Fig. 3. (a) Normalized transmission change of the device (with 80% transmittance of DWCNT films) plotted as functions of wavelength and gate voltage. (b) The variation of normalized transmission as a function of wavelength and positive voltage (with 80% transmittance of DWCNT films). (c) Normalized transmission change of the device (with 50% transmittance of DWCNT films) plotted as functions of wavelength and gate voltage. (d) The variation of normalized transmission as a function of wavelength and positive voltage (with 50% transmittance of DWCNT films).

Nevertheless, the transmission change in longer wavelengths starts to saturate from 1.2 V, and the transmission comes to a reverse change tendency at 2.1 V, which is probably a behavior caused by the quantum confined Stark effect in the semiconductor nanotubes. A further extension to the gate voltage intensity will induce irreversible degradation of the electrolyte. Thus, the amplitude modulation is also limited. Note that further increase in modulation intensity is not constrained by the nanotube material, whereas the film thickness might be an additional degree of freedom for more efficient modulation at the same gate voltage.

A comparison experiment was carried out in another device by replacing the DWCNT electrodes with thicker films (i.e., DWCNT film of 50% transmittance at 550 nm). The mapping graph and transmission change against gate voltage are illustrated in Fig.  $3(c)$  and (d). The transmission change as a function of the gate turns asymmetric, which can be attributed to a memory effect caused by multiple measurements. In experiment, we first carried out the positive gate voltage dependent transmission spectra for several times to identify its reproductivity. These multiple measurements introduced residual electrical doping to the DWCNT films when the gate is unapplied. The followed transmission spectra measurement in negative gate voltage drops in efficiency because part of the potential was used to remove the residual doping in reverse gate. This memory effect looks more pronounced in thicker DWCNT films compared to the mapping in Fig.  $3(a)$ . In contrast to the spectrum distribution in Fig.  $3(b)$ , the result in Fig.  $3(d)$  gives an absorption modulation increase to ∼24.4% at maximum, identifying that the modulation is thickness-dependent in our devices. For clarification, the tiny spikes of the spectra in Fig. 3(b) are from the random background noise of the white light source in the UV-VIS-NIR spectrometer rather than the device.



Fig. 4. Schematic representation of the THz time domain spectroscopy and detection system. M, reflection mirror; L1, L2, L3, L4, collimation and focus lenses. HWP1, HWP2, half-wave plates; P1, P2, polarizers.

## *D. Electro-Optical Modulation at Terahertz Wave*

Electro-optical response of the DWCNT modulator was also characterized with a THz time-domain spectroscopy and detection system in the frequency range of ∼0.2–1.05 THz. We measured the time-domain spectrum variation as a function of the gate voltages within −2.1 to 2.1 V. The experimental setup is drawn in Fig. 4. THz wave was acquired from a GaAs photoconductive emitter under the nonlinear excitation of femtosecond pulses (800 nm, 150 fs, 78 MHz) generated in a mode-locked Ti: sapphire laser. The incident laser pulses are guided to a compressing/expanding telescope (L1, L2) to get  $\sim$ 3 mm beam size in diameter. Half-wave plate HWP1 and polarizer P1 are utilized to split the laser beam for pumping of the THz emitter. The power can be adjusted by tuning the angle of HWP1. The split beam is then guided to an optical delay line and a fast delay line, both are based on hollow retro-reflector. The retro-reflector of the fast delay line is mounted on a fast translation voice coil stage, which enables 10 times moving per second. The beam is then focused by lens L3 to the THz emitter as pump excitation. On the other hand, the half-wave plate HWP2 and polarizer P2 are used to adjust the beam power pumped to the THz detector. Lens L4 focuses the beam to the THz detector (photoconductive antenna). The time-domain spectrum test can be obtained by placing the sample between the THz emitter and detector. In addition, a Teflon lens is used after the THz emitter to get a focused beam at the sample position. The entire setup shown in Fig. 4 was placed in a closed chamber pumped with nitrogen gas to eliminate the influence of ambient humidity.

Time-domain electrical field variation of the transmitted THz signal as a function of gate voltage is shown in Fig.  $5(a)$ . Here, the measurement is from the modulator with 80% transmittance DWCNT films. These spectra were then converted to transmission change through the Fourier transform and normalized to the spectrum under the 0 V gate, which are illustrated in Fig. [5\(b\).](#page-4-0) It is seen that the maximum change of transmission spectrum does not appear at 0 V. We attribute this to the memory effect introduced from the previous measurements in the visible and infrared range. The peaks with ∼0.5 THz separation in the change of transmission spectrum come from a Fabry-Perot resonance, which is formed by the polycarbonate surfaces.

<span id="page-4-0"></span>

Fig. 5. (a) Time-domain electrical field variation of the THz wave as a function of gate voltage, measured with the 80% transmittance of DWCNT films. (b) Normalized transmission change as a function of THz frequency converted from time-domain spectra (in (a)) by Fourier transform. (c) Time-domain electrical field variation of the THz wave as a function of gate voltage, measured with the 50% transmittance of DWCNT films. (d) Normalized transmission change as a function of THz frequency converted from time-domain spectra (in (c)) by Fourier transform.

 $35C$ 

 $0.4$ 

 $0.6$ 

Frequency (THz)

 $0.8$ 

 $1.0$ 

200

250

300 Time (ps)

Over the entire spectrum range, the maximum difference in the gate-dependent transmission change is approximately 32% at 0.75 THz. In comparison, we also characterized the THz timedomain spectrum variation as a function of gate voltage with the device employing 50% transmittance at 550 nm, the result is shown in Fig.  $5(c)$ . Again, it is shown that the largest electrical field appears at 0.6 V, which means the nanotube films return to their minimum conductivity. The Fourier transform spectra in Fig. 5(d) shows similar behavior against the result in Fig. 5(b). But, the maximum difference of the transmission change is only 7% at 0.75 THz. Note that the change of transmission spectrum corresponding to gate voltage in Fig.  $5(d)$  is much smaller than that in Fig. 5(b) corresponding to the same voltage. This observation also demonstrates a thickness-dependent modulation in terms of intensity in our supercapacitive modulators. Compared to the infrared and visible spectral results, due to the different mechanisms of optical absorption (film conductivity) in THz range, we observed higher modulation efficiency in the device with thinner DWCNT films rather than the device with thicker film. This can be attributed to the larger conductivity change in a less conductive (thinner) film against a more conductive (thicker) one. Apart from the intensity modulation, time response also matters the performance of the device. The charging time of the carbon nanotube capacitor, defined as RC time constant, is used to evaluate the response frequency. It has been known that ion diffusion imposing a limit to the response time of carbon nanotube capacitors in the range from microsecond to millisecond range [\[34\].](#page-5-0) Combine the experimentally measured RC time constant in an analogue graphene-ion liquid capacitor [\[33\],](#page-5-0) the response frequency of our device can be estimated in several to a few tens Hz range.

In this section, we demonstrate an electro-optical absorption modulator by employing DWCNT nanofilms as electrodes in a capacitive device geometry. Aerosol-like DWCNTs were synthesized in an FC-CVD stove and deposited as nanofilms with controllable thicknesses. Ionic gel electrolyte with large electrochemical window was utilized as spacer in the capacitor for gate-induced supercapacitor at the nanotube interfaces, which was demonstrated as an effective way for large concentration carrier doping under low gate voltages. Our modulators perform gate-dependent transmission modulation from the visible to infrared wavelengths (500–2100 nm) through interband transition and optical conductivity modulation at THz wave  $(\sim 0.2 - 1.05$  THz) by intraband transition. The gate-controlled transmission change in our DWCNT films shows thickness-dependent behavior, which increases as a function of thickness in the visible to infrared range while decreases in THz wave. Our results identify that DWCNT is a promising active electro-optical medium for ultra-broadband light manipulation with the combination of controllable intensity, low gate voltage, cost efficiency and easy fabrication.

#### III. SUMMARY AND PERSPECTIVES

In summary, we have demonstrated an ultra-broadband electro-optical absorption modulator by using FC-CVD synthesized DWCNTs as active material. The comparison experiment reveals that the intensity modulation of the DWCNT films is thickness-dependent over all the measured spectrum range. In particular, the ultra-broadband absorption modulation is mainly attributed to the formation of electrical double layers at the ionic gel and DWCNT film interfaces. The exponential potential drop cross ultrathin Helmholtz layers (*<*1nm) allows high-density charge accumulation [\[35\]](#page-5-0) near the DWCNT films, which dominates the absorption extension toward the visible range. It is notable that the electrical double layer is almost independent to the thickness of the ionic gel layer [\[36\].](#page-5-0) Nevertheless, in contrast to the fast carrier mobility in DWCNTs, the slow ion diffusion in ionic gel limits the time response of the device. Alternatively, if the ionic gel is replaced by solid dielectric (e.g.,  $\text{Al}_2\text{O}_3$ ,  $\text{SiO}_2$ ), the device will function as an electrostatic capacitor. In this model, the charge density scales linearly against the applied voltage and the dielectric constant, and it is inversely proportional against the thickness of the spacer between the DWCNT electrodes [\[37\].](#page-5-0) In order to obtain broader absorption modulation, three approaches can be used, i.e., select the dielectric material with higher permittivity, reduce the thickness of the spacer, and increase the gate voltage. However, the improvement is limited because of breakdown issue. Therefore, one can hardly achieve absorption modulation in the visible spectrum or even near infrared range with these modulators. Considering that the optical conductivity of DWCNTs is also determined by the charge density, the modulation efficiency in THz wave will decrease. On the other hand, electro-optical modulators with solid dielectric are competitive in applications where high speed response is critical (e.g., ultrafast signal processing). But complicated device design and nanofabrication process bring huge challenges to overcome the RC time constant <span id="page-5-0"></span>bottleneck for high speed operation. Here, we emphasize the ultra-broadband electro-optical absorption modulation (from visible to THz wave) in semiconducting DWCNTs by using ionic gel as dielectric material. Our findings open new avenues in the field of electro-optical modulators, highlighting potential alternatives of using one-dimensional carbon nanotubes, and broadening the spectrum of materials available for advanced photonic applications.

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