# A Lightweight Spherical Generalized Luneburg Lens Antenna With Low Cross-Polarization Over a Wide Range in Azimuth and Elevation 

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#### Abstract

In this paper, a novel dual-slant polarized three-dimensional (3D) periodic Luneburg lens with a diameter of $390 \mathrm{~mm}\left(4.6 \lambda_{0}\right.$ at 3.55 GHz$)$ is presented. Copper-plated cubes with truncated corners are placed in a body-centered cubic (BCC) lattice and held together with layers of Rohacell foam. The sizes of the cubes are varied to realize the gradient refractive index (GRIN) profile of a generalized Luneburg lens at a low cost, with a low weight and loss. The designed Luneburg lens operates from 3.3 to 3.8 GHz with a total weight of 1 kg . The quasi-isotropic response of the proposed periodic structure allows a wide angle coverage. Measured results show that the lens antenna can achieve a peak gain of 22 dBi with a scanning loss of less than 0.4 dB in a wide angular range in both the azimuth and the elevation plane. Importantly, the proposed lens antenna design achieves a cross-polarization level below peak gain of less than -19 dB at all angles for the two slant polarizations of the feed, while comparable designs report up to -11 dB at given angular directions. This high-gain multi-beam antenna with a commercially viable design is suitable for wireless communications at low microwave frequencies.


INDEX TERMS Body-centered cubic (BCC) lattice, gradient refractive index (GRIN) material, quasiisotropic periodic structure, multi-beam Luneburg lens antenna, wide angle coverage.

## I. INTRODUCTION

FUTURE generations of communication systems are expected to meet the demands for faster mobile broadband speeds, higher network capacity, and lower latency [1], [2]. In this regard, the design of antennas with dual polarization and highly directional multiple beams is of critical importance [3]. Dual-polarized antennas enable polarization diversity, and multiple beams provide spatial domain multiplexing, thus leading to increased system capacity. Different realizations of beamforming devices such as Butler matrices [4], [5], Rotman lenses [6], [7], planar and three-dimensional (3D) Luneburg lenses [8], [9], [10], [11], [12], [13] have been introduced in recent years. Planar quasi-optical beamforming structures can deliver high-gain scanned beams at low cost [14], among which, planar

Luneburg lens antennas are of interest for their wide scanning range [15]. Different designs of parallel-plate cylindrical Luneburg lenses have been produced with the ability to create high directive beams in their H-plane [16], [17], [18], [19], [20]. However, such planar structures provide the communication systems with only one-dimensional (1D) steering, whilst in many practical applications, 2D beam steering is required.

Spherical Luneburg lenses are simple and attractive passive steerable antennas for wireless communication systems. The spherically symmetric structures enable two-dimensional (2D) beam steering over a wide-angular range without performance degradation [8]. The traditional method for the realization of such 3 D lenses is the layered approach, where the lens is divided into a series of shells with
different refractive indexes (RIs) [21], [22]. The fabrication and assembly of several spherical shells with different RIs can be complex and costly. Alternatively, the Gradient Refractive INdex (GRIN) profile can be realized using quasiperiodic structures, resulting in Luneburg lenses with a gradual variation in their effective material properties and, consequently, low reflections within the lens.

Various quasi-periodic Luneburg lenses have been reported. For instance, Luneburg lenses based on planar layers with periodic holes drilled in dielectric slabs were reported in literature [23], [24], [25]. Recently, additive manufacturing has been used to produce Luneburg lenses operating above 10 GHz [12], [26]. This technique offers a high manufacturing resolution. However, the dielectric material and methods used in these lens designs may result in heavy and costly devices at low frequency (below 10 GHz ). In [27], a low-weight, low-cost, layered structure for the construction of spherical Luneburg lenses at lower microwave frequencies was described. This method, while suited for low-weight designs, suffers from high crosspolarization at wider scan angles. A modified structure was consequently introduced in [28] where the lens layers were designed to achieve low anisotropy in the azimuth plane, at the expense of degraded performance in the elevation plane. However, the main application of the spherical Luneburg lens is to fully steer the beam in both azimuth and elevation planes, which motivates the reported research.

This paper reports a novel layered design of a periodic 3D Luneburg lens consisting of Rohacell foam with conducting inclusions placed in a novel lattice. Specifically, the periodic inclusions are arranged in a body-centred cubic (BCC) lattice, and we demonstrate that this lattice arrangement allows for a quasi-isotropic response. We use this periodic structure to design the first reported light-weight multibeam Luneburg lens antenna with wide angular coverage in both the azimuth and elevation planes, where high crosspolarization isolation and efficiency are maintained in all angular directions. Importantly, compared with previously reported Luneburg lens antennas with the inclusions placed in simple cubic [27] or hexagonal lattices [28], the presented lens antenna maintains its attractive radiation properties for scanning in the whole angular space, demonstrating the benefits of the proposed lattice arrangement for lens designs. Specifically, the comparative analysis reported here demonstrates an improvement in cross-polarization of up to 8 dB and a reduction in side lobe level (SLL) of up to 4 dB . The use of the BCC lattice is inspired by glide-symmetric structures [29], and the observed quasi-isotropic response is similar to the previously observed response in planar glidesymmetric lenses [10], [20]. The proposed quasi-isotropic periodic structure greatly simplifies the construction of spherical lenses by virtue of planar layers instead of shells, whilst keeping the RI profile consistent. Besides, the use of foam as the supporting material makes the design light, and thus suitable for applications in bands below 10 GHz , where the weight of conventional Luneburg lenses can be


FIGURE 1. (a) Calculated RI of the lens when the focus is moved to $1.3 R$ ( $R$ is radius of the lens, $r$ is the radial position.) (b) Ray tracing of the generalized Luneburg lens.


FIGURE 2. Configuration of the unit cells (left) and definition of the points of high symmetry delimiting the irreducible Brillouin zone of interest (right). The inclusions are shown in a simple cubic (SC) and in a body-centred cubic (BCC) lattice, where $p$ is the period, and $s$ is the side length of the inclusion.
prohibitive. It is worth to note that the volume and mass of the lens increases cubically with the radius, meaning that this is increasingly important for lower frequency and higher gain.

The high gain multi-beam lens antenna presented in this work can find applications in wireless communication systems such as fourth-generation (4G) and fifth-generation (5G) networks, specifically in the so-called frequency range 1 (FR1) corresponding to below 6 GHz frequency band allocations. The proposed design features low fabrication cost while enabling polarization diversity suitable for high capacity communication systems.

## II. DESIGN OF THE LUNEBURG LENS ANTENNA

This section describes the lens antenna design. First, a quasi-isotropic structure for the realization of the effective RI profile of the lens is introduced. Then, the lens is implemented and tested with a square waveguide feed.

## A. REALIZATION OF THE QUASI-PERIODIC LENS

A generalized Luneburg lens with a focal point displaced $0.3 R$ from the lens surface, where $R$ is the radius of the lens, is designed [30], [31], [32]. The corresponding RI distribution and a ray tracing representation for this lens are provided in Fig. 1. The required RI profile of the lens is


FIGURE 3. Effective refractive index of the SC and BCC lattice for propagation along the symmetry directions. (a) The structures have $p=26 \mathrm{~mm}$ and different $s$. (b) The structures have $s=12.5 \mathrm{~mm}$ and different $p$.
realized using a quasi-periodic structure, where the index of refraction is controlled by varying the size of the inclusions.

A 3D periodic quasi-isotropic BCC structure is introduced, analyzed, and its response is compared with that of a structure arranged in a simple cubic (SC) lattice. The SC structure was used in the design of the lens in [27]. It is worth noting that a dielectric BCC structure was proposed in [33] and was demonstrated to provide wide bandwidth and robust realization using additive manufacturing. The unit cells of the two structures studied here are shown in Fig. 2, where the gray color corresponds to metal and the green color corresponds to the host medium, which is assumed to be a vacuum in the simulation models. The BCC structure is inspired by previously reported 2D glidesymmetric structures [20], [34].

In a lens antenna based on these structures, the field within the lens can be expressed as a combination of these two modes with the relative amplitude between the modes depending on the local environment in the lens, the source type and source position. This means that, since the response is different for the different modes, we expect that the lens introduces some polarization impurity and phase errors, and that this performance degradation depends on the feed position.

Here, the dispersion diagrams for the two structures are obtained using the CST eigenmode solver. The results are shown in Fig. 3 for the propagation along $\overline{\Gamma X}, \overline{\Gamma M}$, and $\overline{\Gamma \mathrm{R}}$. At the considered frequency range, the structures support


FIGURE 4. Comparison of the effective refractive index for propagation along $\Gamma \rightarrow R$ at 3.5 GHz with and without chamfered corners. The dimensions are $c=1 \mathrm{~mm}$ and $p=26 \mathrm{~mm}$.
propagation of waves in two different modes, and these are indicated with different line markers. It is worth noting that in a lens antenna based on these structures, the field within the lens can be expressed as a combination of these two modes with the relative amplitude between the modes depending on the local environment in the lens, the source type and source position. This means that, since the response is different for the different modes, we expect that the lens introduces some polarization impurity and phase errors, and that this performance degradation depends on the feed position. With this said, we note that the structures appear identical for the two modes for propagation along $\overline{\Gamma X}$ and $\overline{\Gamma R}$, but not for $\overline{\Gamma M}$. In Fig. 3(a), the period is the same in the structures and the side length $(s)$ is tuned so that the effective RI in the two structures at low frequencies are the same. We observe that the bandwidth of operation is significantly wider in the BCC structure. Furthermore, the response of the BCC structure for propagation along $\overline{\Gamma M}$ is more similar for the two modes, compared to the SC structure. As a result, we expect the performance to be more stable with scanning angle in a lens based on the BCC structure, compared to a lens based on the SC structure. It is noted that the bandwidth in the SC structure can be increased by reducing the period, as demonstrated in Fig. 3(b). However, we observe that the SC structure in this case still remains slightly more anisotropic and that a lens implemented with the SC structure requires a larger number of inclusions in a given volume, which can result in increased manufacturing complexity.

As required by the manufacturer to avoid issues related to tolerances, the corners of the cubic inclusions are truncated. The inclusion of chamfered corners for use in the fabricated prototype is shown in the inset of Fig. 4. Here, $c$ is the chamfering of the corners, which is chosen to be $c=1 \mathrm{~mm}$ to comply with the manufacturing requirements. In Fig. 4, we observe that the chamfering only has a small impact on the effective refractive index. This effect is considered for the lens design.

Considering this unit cell configuration allows choosing higher periodicity with lower dispersion, the unit cell size in this design is chosen to be equal to $26 \times 26 \times 26 \mathrm{~mm}^{3}$, corresponding to $p=26 \mathrm{~mm}$, which is $0.3 \lambda_{0}$ at mid-band frequency. The side length of the cubic inclusions with 1 mm


FIGURE 5. Lens antenna in full-wave simulator. (a) The lens antenna structure. (b) E-field distribution at mid band frequency. The lens radius is 195 mm ( $2.3 \lambda_{0}$ at 3.55 GHz).

TABLE 1. Radial distribution of inclusions versus side length.

| Side length, $s$ | Radial range | Quantity |
| :---: | :---: | :---: |
| 7.5 mm | $196-199.5 \mathrm{~mm}$ | 930 |
| 9 mm | $191-196 \mathrm{~mm}$ | 560 |
| 10 mm | $178-190 \mathrm{~mm}$ | 470 |
| 11 mm | $166-178 \mathrm{~mm}$ | 564 |
| 11.5 mm | $152-166 \mathrm{~mm}$ | 518 |
| 12 mm | $131-152 \mathrm{~mm}$ | 418 |
| 12.5 mm | $111-131 \mathrm{~mm}$ | 292 |
| 13 mm | $0-110 \mathrm{~mm}$ | 15 |

chamfering varies from 7.5 mm to 13 mm in minimum steps of 0.5 mm to approximate the effective RI distribution as required for this lens design. Eight different sizes of inclusions are used and their radial distribution is detailed in Table 1. From previous works on discretized Luneburg lenses, this is considered to provide sufficient sampling of the continuous RI profile [35]. Note that the inclusions at the center of the lens require a cube with a side of half-a-period, hence the benefit of using chamfered corners to avoid issues related to manufacturing tolerances.

## B. LENS ANTENNA DESIGN

A square waveguide feed excited by a PCB crosseddipole [27] is used to investigate the dual-slant polarized performance of the lens. The feed is placed at a distance of $0.3 R$ from the lens surface. The simulated model of the lens antenna system and E-field distribution is shown in Fig. 5. The simulated normalized radiation patterns of the lens antenna at mid-band frequency ( 3.55 GHz ) are compared with the ones in [27], [28] to demonstrate the superior performance of this new design. The results are shown in Fig. 6. Comparison of the plots shows that the proposed lens antenna offers the most stable performance across the analysed scanning range. For the lens with cubic lattice and cubic inclusions [27], high SLL and cross polarization discrimination (XPD) are observed at wider scan angles ( $\pm 45^{\circ}$ for this case) where defocusing causes larger phase error on the lens aperture. The modified design of hexagonal lattice with cylindrical inclusions introduced in [28] improves the scanning performance in the azimuth plane and delivers


FIGURE 6. Comparison of the simulated azimuth radiation patterns (RP) of this work (LL3) with the ones in [27] (LL1) and [28] (LL2) at 3.55 GHz at (a) $0^{\circ} \mathrm{Az} 0^{\circ} \mathrm{El}$ (b) $30^{\circ} \mathrm{Az}$ $0^{\circ} \mathrm{El}$ (c) $30^{\circ} \mathrm{Az} 30^{\circ} \mathrm{El}$ (d) $45^{\circ} \mathrm{Az} 0^{\circ} \mathrm{El}$ (e) $0^{\circ} \mathrm{Az} 45^{\circ} \mathrm{El}$ (f) $45^{\circ} \mathrm{Az} 45^{\circ} \mathrm{El}$.
the best figures for scanning in the azimuth plane. However, when it comes to the elevation scanning, the SLL and XPD becomes the worst, mainly because this design in not symmetric along the elevation plane. On the other hand, with the proposed lens design using BCC unit cells, the patterns repeat every $45^{\circ}$. This is due to the symmetries in the unit cell layout as shown in Fig. 5, which preserve the symmetries already discussed at unit cell level in Section II-A. As a consequence, the lens is invariant upon rotation of $90^{\circ}$ around any of the main axes. Further combined with reflection symmetries, the results reported here can be reduced without lack of generality to a scanning range of $\left[0^{\circ} 45^{\circ}\right]$ in both azimuth and elevation. In this lens antenna, the SLL and XPD stay below -15 dB and -19 dB , respectively, across all the angles. The gain, beamwidth, and aperture efficiency are similar to the other two lenses. The performance in SLL, and more generally beam shape, provide a good indication of the focusing characteristics of the lenses, while the XPD is a good indicator of the polarization sensitivity. The results reported demonstrate the improvement achieved using quasiisotropic unit cells.

## III. EXPERIMENTAL RESULTS

In this section, the fabrication process of the lens antenna is described. The prototype lens antenna is measured in


FIGURE 7. Lens construction. (a) Copper plated foams inclusions. (b) Foam layer. (c) Inclusions placed in a layer. (d) Prototype of the complete spherical Luneburg lens, including the reference coordinate system in which the azimuth and elevation angles are defined.
an anechoic chamber. Finally, the simulation and measured results are compared with the state-of-the art Luneburg lenses.

## A. LENS FABRICATION AND TEST

The presented Luneburg lens antenna has been fabricated and tested. The lens has a diameter of 390 mm in the vertical direction and a diameter of 400 mm in the horizontal to ensure that inclusions are not too close to the edge. It comprises 30 cylindrical layers of Rohacell (a lightweight structural foam with a dielectric constant close to that of air) and a thickness of 13 mm for each layer. Through square holes were cut from foam layers to accommodate the inclusions. The layers are stacked up in the vertical direction to form the complete sphere. A total of 3479 inclusions are required for the lens design, as detailed in Table 1. Two sets of inclusions were fabricated. The first three big sizes were made by 3D printed plastic material coated with copper. The corners are also chamfered as shown in Fig. 7 (a). The rest of the inclusions were cut from foam and electrolytically plated with copper. The reason for such a choice is to avoid contact of big inclusions, requiring more precise manufacturing, while smaller inclusions are less sensitive and can be manufactured using cheaper processes. Next, the inclusions are placed inside the layers' holes. One of the foam layers, inclusions, and a layer of lens are shown in Fig. 7. The complete lens prototype weights almost 1 kg , which is very lightweight for a lens of this size. Since the weight of the system scales cubically with the radius of the lens, the density of the material used for implementation must be minimized. A similar lens implemented with dielectric thermoplastic or polystyrene, having a density typically around $1 \mathrm{~g} / \mathrm{cm}^{3}$, would weight approximately 20 kg , which is significantly heavier than the proposed design. This estimate of the weight accounts for the filling factor varying with the refractive index profile. As an example, the additively manufactured half-Gutman lens in [33] weights 77.4 g for a radius of 4 cm , corresponding to an average density of $0.3 \mathrm{~g} / \mathrm{cm}^{3}$, which would scale up to 19.4 kg for a full lens with a radius of 20 cm similar to the one considered in this work.


FIGURE 8. Lens antenna prototype in the anechoic chamber. The azimuth is measured by rotating the lens about its axis. The elevation scanning is measured by rotating the feed along the lens elevation.


FIGURE 9. Simulated and measured reflection coefficients of the feed and of the lens with feed.

Far-field measurements of the lens antenna were made in the NSI anechoic chamber at UTS Techlab, NSW, Australia. To hold the lens on the positioner, a wooden fixture was constructed that also supports the feed at a distance of 60 mm $(0.3 R)$ from the lens surface. It allows to rotate the feed at different elevation angles, while the azimuth scanning performance is tested by rotating the lens along its polar axis. The lens antenna measurement setup is shown in Fig. 8.

## B. SIMULATION AND MEASUREMENT RESULTS

In this section, the simulated and measured results of the lens antenna are compared. First, an Agilent Network Analyzer N5225A has been used to measure the reflection coefficients of the feed alone and with the lens antenna. These results are given in Fig. 9. The feed rotation along the lens surface does not affect significantly the reflection coefficient of the

TABLE 2. Comparison with other reported spherical Luneburg lens antennas.

| Ref. | Realization | Feed Type | Frequency <br> Band $(\mathrm{GHz})$ | Scan Angle <br> Az/El (deg) | Gain <br> (dBi) | HPBW <br> Az/El (deg) | $\begin{array}{\|c\|} \text { SLL } \\ \mathrm{Az} / \mathrm{El}(\mathrm{~dB}) \\ \hline \end{array}$ | Radiation Efficiency (\%) | Aperture <br> Efficiency <br> (\%) | XPD at wide angles <br> (dB) | Pol. |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| [24] | Perforated <br> slices | Open-ended waveguide | 12-18 | $\pm 75$ <br> (Azimuth only) | 23.1-25.7 | 11/14 | -9/-22.8 | 80 | 47-93 | N.G. | LP |
| [26] | 3D Printing (cubes) | Open-ended waveguide | 8-12 | N.G. | 17.3-20.3 | 13/19 | -20/-25 | N.G. | 47-53 | N.G. | LP |
| [36] | Randomely distributed conductive fibres | Horn | 3-6 | Full range | 23-28.5 | 5.2-10.6 | $\begin{gathered} -21 \\ \text { (boresight) } \end{gathered}$ | N.G. | 50-56.2 | N.G. | $\begin{gathered} \text { LP } \\ (\mathrm{VP}) \end{gathered}$ |
| [27] | Conductive cubes in cubic grid | Open-ended waveguide | 3.3-3.8 | $\pm 60 / \pm 60$ | 21.9-23.4 | 13/14.5 | -12/-11 | 91 | 75-85 | -11 | Dual-LP <br> (Slant) |
| [28] | Conductive cylindrical in hexagonal grid | Open-ended waveguide | 3.3-3.8 | $\pm 60 / \pm 30$ | 21.3-22.8 | 12.4/13.9 | -16/-12 | 97 | 71-84 | -11 | Dual-LP <br> (Slant) |
| This <br> work | Conductive BCC in cubic grid | Open-ended waveguide | 3.3-3.8 | Full range | 21.2-23.3 | 13/14.2 | -15/-15 | 95 | 71.5-89 | -19 | Dual-LP <br> (Slant) |

* Az/El:Azimuth/Elevation; HPBW: half-power beamwidth; LP: linear polarization; N.G.:not given; SLL: sidelobe level; XPD: cross polarization discrimination.


FIGURE 10. Simulated (solid line) and measured (dash line) normalized co-polar and cross-polar radiation patterns of the lens antenna at $+45^{\circ}$ polarization for the beams at (a) $0^{\circ} \mathrm{Az} 0^{\circ} \mathrm{El}$. (b) $45^{\circ} \mathrm{Az} 0^{\circ} \mathrm{El}$. (c) $45^{\circ} \mathrm{Az} 30^{\circ} \mathrm{EI}$. (d) $45^{\circ} \mathrm{Az} 45^{\circ} \mathrm{EI}$. Frequency superimposed: $3.3 \mathrm{GHz}, 3.5 \mathrm{GHz}, 3.8 \mathrm{GHz}$. (Note: The lens has 4 -fold symmetry, so the patterns at $0^{\circ}$ Az $45^{\circ} \mathrm{El}$ are similar to the ones at $45^{\circ} \mathrm{Az} 0^{\circ} \mathrm{El}$.)
lens antenna system, so this figure is only shown for the feed at $0^{\circ}$ and $45^{\circ}$ azimuth. The reflection coefficient of the feed with the lens is less than -20 dB across the band. The measured port-to-port isolation is below -30 dB and is not reported here for brevity. Good agreement between the simulated and measured results is observed. Then, radiation property of the proposed 3D Luneburg is investigated in detail.

The simulated and measured radiation patterns of the lens antenna at normal incidence and some of the selected angles at three different frequencies across the band (3.3, $3.55,3.8 \mathrm{GHz}$ ) for one of the slant polarizations are shown
in Fig. 10. Since the lens patterns repeat after $45^{\circ}$ in azimuth and elevation as a direct consequence of the symmetries discussed in Section II-B, we only show the patterns up to this angle, and the reported results cover the worst case. The simulated and measured results agree well. Although the beam shape and beamwidth is stable across the band at different angles, higher SLL is observed in the measured results, which is attributed mostly to the test fixture. Fig. 11 compares the simulated and measured gain with the maximum ideal directivity, which is evaluated based on a uniform illumination of an aperture with area $A$ (i.e., $D_{\max }=4 \pi A / \lambda^{2}$ ), where $A$ is the cross section area


FIGURE 11. (a) Simulated (black lines) and measured (red lines) realized gain, and maximum ideal directivity. (b) Aperture efficiency versus gain.
of the lens. The plots show that the realized gain varies from 21.15 to 23.3 dBi at all angles with a scan loss of less than 0.4 dB in all cases. Comparison of the simulated and measured realized gain and maximum ideal directivity shows the aperture illumination efficiency is better than $70 \%$ across the band from 3.3 to 3.8 GHz , while the radiation efficiency stays better than $95 \%$. These plots are shown in Fig. 11 (b). It is to be noted that much of the losses are due to spillover from the feed, corresponding to the radiated power not being intercepting by the lens. This demonstrates the very low loss nature of the lens material resulting in a very high efficiency.

## C. COMPARISON AND DISCUSSION

The Luneburg lenses are primarily used in high frequencies and mm-wave applications. A direct comparison with the mm-wave designs is difficult because the technique suitable for such methods is not practical for low microwave frequencies. In Table 2, the performance of the Luneburg lens antenna reported in this work is compared with the state-of-the-art lens antennas. The perforated slices and 3D printing techniques used in [24], [26] are not applicable for the fabrication of the very large lens antennas at lower microwave frequencies. Such material has high density making the design not viable for commercial use. In [37], concentric spherical shells comprising lightweight foam with embedded randomly conductive wires are used, which
overcomes the weight problem. From a construction point of view, however, maintaining the consistency of the profile within each shell is difficult. The last three entries correspond to the layered lens designs at lower microwave frequencies, where all the designs use the same feed and have the same size and operating frequency within the band 3.3 GHz to 3.8 GHz , with the possibility to extend the bandwidth as described in [28]. Comparisons of the radiation pattern of these lenses as displayed in Fig. 6 show that all three lenses perform the same at boresight, while they differ in scanning performance. The design in [27] shows variation in SLL, and cross-polarization level with azimuth scanning. Although the hexagonal lattice used in [28] reduces the anisotropy in the azimuth plane and has lower fabrication complexity thanks to cylindrical inclusions and reduced number of layers, the performance deteriorates at larger elevation angles. The quasi-isotropic periodic structure in this work allows the best scanning performance among all other reported works, making the design suitable for full coverage high gain multibeam steerable antennas.

## IV. CONCLUSION

A multi-beam high-gain antenna has been presented using a 3D generalized Luneburg lens fed by a dual-polarized square waveguide operating at sub-6GHz bands. The GRIN material is realized by placing conductive inclusions in a BCC lattice supported by Rohacell foam layers. The design shows low loss and is highly efficient, and produces directive beams in the whole space without any notable performance variation when fed from different positions. This wide angle scanning capability is a consequence of the symmetry of the BCC lattice, which can be used to overcome the scanning limitations in previously reported lens antennas using similar inserts [27], [28]. The concept is demonstrated by an experimental prototype and the simulated and experimental results agree well. The presented work is a suitable high gain beam steering and multi-beam antenna for use in commercial microwave communications systems such as $4 \mathrm{G}, 5 \mathrm{G}$ sub- 6 GHz band, GPS, WiMax, etc.

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