Integrated Optical Fiber-Tip Cantilevers

Christopher Holmes, Alexander Jantzen, Alan C. Gray, Lewis G. Carpenter, Paul C. Gow, Stephen G. Lynch, James C. Gates, and Peter G. R. Smith

***Abstract*— A microcantilever at the end face of an integrated optical fiber is reported, fabrication is uniquely achieved using a precision dicing saw. The methodology is a single-step rapid process, capable of achieving trenches with high aspect ratio (>10:1). The platform on which fabrication is made is a monolithic, integrated optical fiber. This integrally fuses optical fiber to a planar substrate using flame hydrolysis deposition and high temperature consolidation (>1000 °C). This paper is the first report of a fiber-tip cantilever using the technique and this integrated platform. As an approach to quantify the optical response of such a multicavity arrangement, a method using Mason's rule is presented. This is used to infer the spectral responses of individual cavities formed and through physical actuation, an estimation of the cantilever's spring constant is made.**

***Index Terms*— Integrated optics, optical fiber devices, optical fibers, optical interferometry, optical sensors.**

I. Introduction

OPTICAL alternatives to Microelectromechanical Systems (MEMS) are becoming of increasing interest to science and technology due to their immunity to electromagnetic interference, inherent safety in flammable environments and compatibility with optical-fiber. Recent developments in forming cantilevers at the tip of optical fibers have opened-up exciting new opportunities for miniaturized optical sensors [1]–[3]. In particular measurements of fluidic flow [4], bio-mechanical characterization [5], [6], Atomic Force Microscopy (AFM) imaging [7], and chemical sensing [8] have all recently demonstrated using this format. Typically these systems use Fabry-Pérot (FP) interferometry for monitoring, largely due to the inherent cavity formed between the fiber and cantilever. The spectral response from such cavities generally has a low-finesse due to the weak Fresnel reflections from the silica-air interfaces. However, in some fabrication methodologies finesse can be enhanced through use of additional coatings. The physical fabrication approaches reported thus far have included picosecond-laser machining [4], focused ion beam [9], wire-cut micromachining [6] and photolithography [10].

This work reports cantilever fabrication solely through the use of physical micromachining [11]. The technique has the advantages of removing large amounts of material (mm^3) quickly (minutes) whilst still maintaining a vertical form factor. The platform chosen on which to demonstrate this technique is a novel Integrated Optical Fiber (IOF) [12], [13], illustrated in Figure 1. IOF uses Flame Hydrolysis Deposition (FHD) to robustly form a miscible alloy between the optical fiber and the planar substrate. IOF offers superior mechanical strength and has advantages associated with integration including thermal homogeneity and the ability to fabricate multiple components upon a single compact chip.

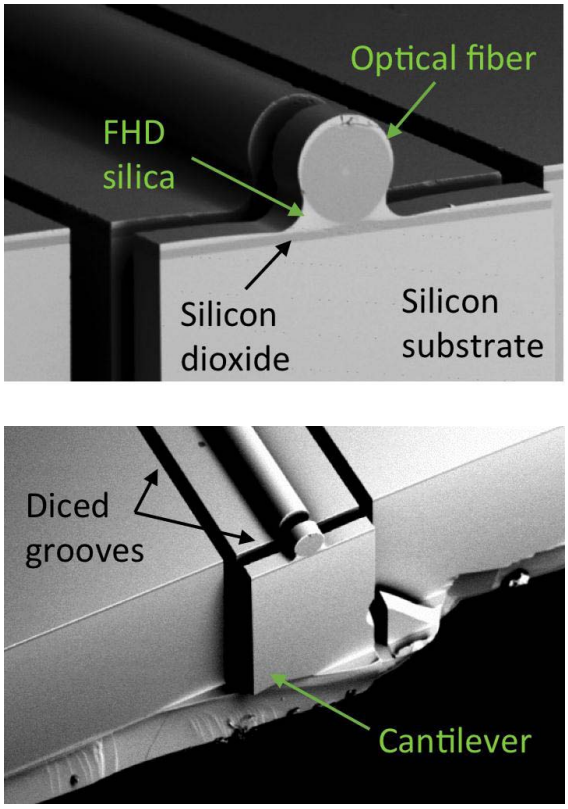

Fig. 1. Scanning electron microscope image of a physically micromachined integrated optical fiber-tip cantilever.

It must be stressed that the fabrication approach reported solely uses physical micromachining and is therefore different to other reports that use a combination of physical machining and wet-etching [14].

II. Fabrication Technique

IOF adapts the commercial glass deposition technique FHD, which is commonly used for the fabrication of Array Waveguide Gratings (AWGs) and other silica based Planar Lightwave Circuits (PLCs). The adaptation made involves the pre-layering of an optical fiber to a planar substrate. In this

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The authors are with the Optoelectronics Research Centre, University of Southampton, Southampton SO17 1BJ, U.K. (e-mail: christopher.holmes@southampton.ac.uk).

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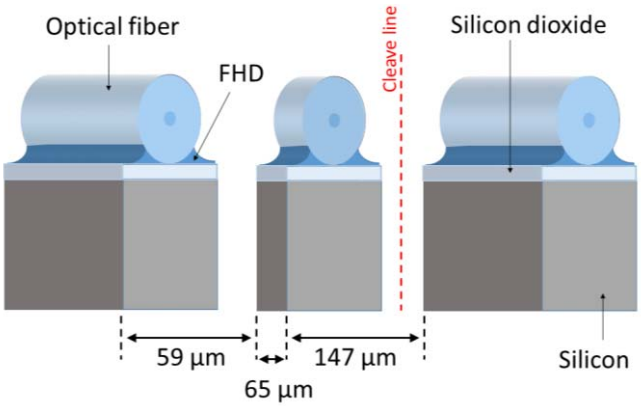

Fig. 2. Schematic side view of fabricated cantilever.

instance a 1 mm thick silicon wafer with a 15 μm thick thermal oxide is used as the substrate. Following glass soot deposition, a consolidation is undertaken at high temperature (1250°C), forming a miscible glass alloy between the optical fiber and thermal oxide. This results in a mechanically robust integrated platform that has the ability to optically guide both in the fiber and the FHD layer(s). It must be noted that in this demonstration the IOF is used solely for its mechanical properties, as all light is guided in the fiber.

Fabrication of the cantilever was achieved with the use of a precision dicing saw (Loadpoint Microace) and a nickel bonded synthetic diamond blade (DiscoTech ZH05-SD4800-NI-50-GG). The depth of cut chosen was 800 μm; rotation speed of 25 krpm and cutting speed of 0.1 mm/s. In total five cuts were made, two parallel and three orthogonal to the fiber. The orthogonal cuts formed air cavities of 59 μm (single cut) and 147 μm (dual cut) respectively and a glass cavity of 65 μm, as shown in Figure 2. To form a distal cantilever this structure was cleaved, which left only a single air and glass cavity.

The Sa surface roughness (which is an arithmetic mean height calculation) of the air-glass interfaces was measured to be 29 nm using a white light interferometer. This is larger than some previously reported results [15], which used a dicing blade containing much smaller grit (5000 grit size used not 4800). It must be noted that in order to achieve the depth-of-cut desired (800 μm) it was necessary to use the more course grit due to commercial availability at the time at which this script was written. Improved interface surface quality should be achievable without compromising sensitivity if further developments of commercial blades are made.

To understand the optical spectra resulting from these cavities a multi-cavity solution was formulated using signal flow graph analysis and Mason's rule.

III. Signal Flow Graphs and Mason's Rule

Optomechanical cantilevers with optical fiber readouts are nowadays routinely interrogated via reliable, low noise interferometers, which guarantee linear response over large displacements and accessing a frequency bandwidth of several tens of kHz without compromising the overall performance. The following approach uses spectrally broadband data and

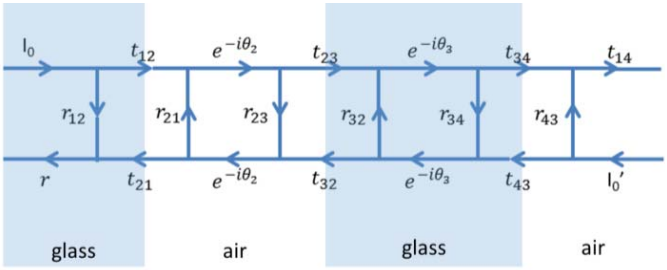

Fig. 3. Graphical representation of light transmitting and reflecting through a glass-air-glass-air cavity.

signal flow graph theory. It is not intended to be a method of optical measurement but rather a method of optical characterization and a route to infer spring constant of this and similar fiber-tip cantilever constructs, e.g. as part of design iteration and quality control assessment.

Signal flow graph analysis of layered media is considered to explain the spectral characteristics of the fabricated cavities. Using Mason's rule a graphical representation of a dynamic system of linear equations can be represented [16]. Figure 3 represents the pathway of light through a glass-air-glass-air interface. This technique can also be expanded to three cavities as in the pre-cleaved structure (shown in Figure 2) or indeed multiple cavities. With respect to nomenclature, t correspondence to transmission coefficient, r the reflection coefficient and θ is a phase term defined as:

$$\theta = \frac{2\pi (n_{eff} l)}{\lambda} \tag{1}$$

n_{eff} is the effective refractive index and l the propagation distance. Each node in the flow graph corresponds to the wave mode amplitude at an interface and the arrows define coupling terms from one node to another. The resultant closed loops being the cavity resonances. It is important to note that Masons rule yields the reflectance via inspection.

The amplitude transmitted through the structure can be obtained by summing all the different paths from source node to output node. Mason's Rule expresses this transfer function of the system as:

$$T = \frac{\sum_k P_k \Delta_k}{\Delta} \tag{2}$$

where the summation is made over all k-paths connecting the input to output node. Δ is the determinant of the system [16], P_k is the k^{th} forward path gain (product of gains found through traversing a succession of branches in the direction of arrows with no node passed more than once) and Δ_k is the determinant of the k^{th} forward loop.

IV. Results

The back reflected spectra of the fabricated cavity was measured using a broadband SLED (Amonics, ASLD-CWDM-5-B-FA) and an Optical Spectrum Analyzer (OSA) (Ando AQ6317B). The observed spectra prior to cleaving (treble cavity) and after cleaving (double cavity) is illustrated in Figure 4. The spectral resolution of the OSA was set to 0.1 nm.

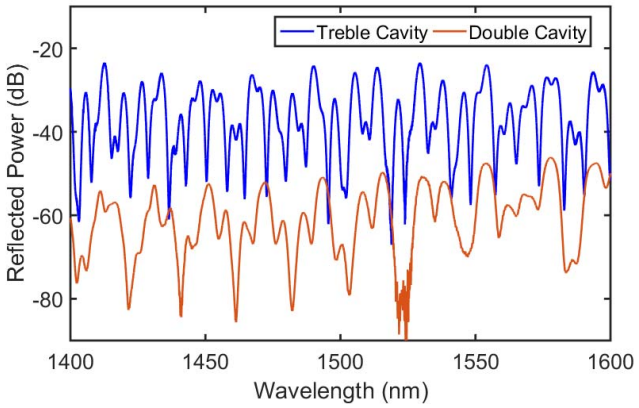

Fig. 4. Measured back reflection from the physically machined integrated fiber prior and post cleave.

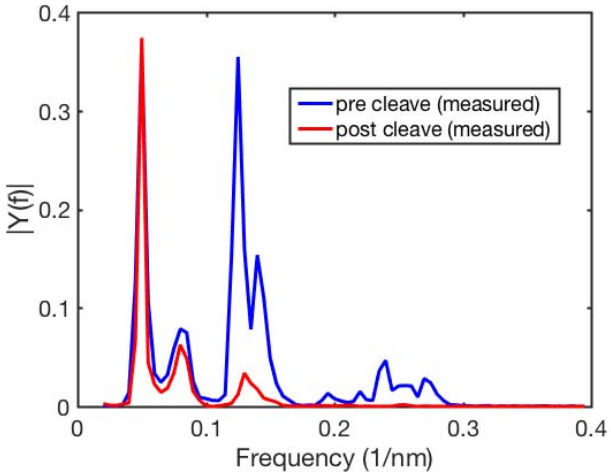

Fig. 5. The Fourier Transform of the measured spectral response pre and post cleave.

To highlight the periodic features from this optical spectra a Fourier transformation was made, shown in Figure 5, which compares the spectra before and after cleaving. This is compared to a theoretical frequency distribution calculated using Mason's rule, illustrated in Figure 6. It must be noted that the theoretical model also accounts for dispersion effects through use of accepted Sellmeier coefficients for silica, however the contribution of this term only marginally improved the fit.

From Figure 5 and Figure 6 it is evident that the main resonance features appear comparable in both the measured and theoretical model. It is noted that variation in amplitude, |Y(f)|, between the theoretical and measured model is consistent with diffraction effects and recoupling in the air cavity sections. Additionally, there are also scattering losses at the interfaces between cavities. These parameters were not directly accounted for, but it is understood that divergence at an air gap of ~60 μm reduces fringe visibility by ~25% [17].

A. Force Calibration

Lateral force calibration was made using a KLA Tencor P16 stylus profiler. The forces applied are presented schematically in Figure 7. In the arrangement the stylus was translated such that its leading edge was the only point of contact with the chip. The stylus was scanned in the direction of forward

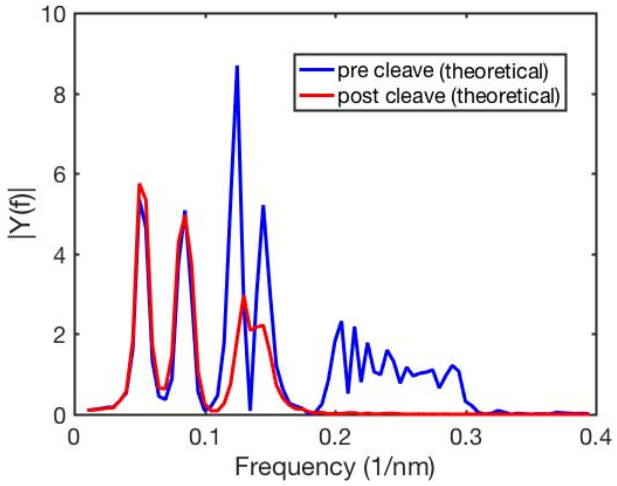

Fig. 6. The Fourier Transform of the theoretical spectral response pre and post cleave.

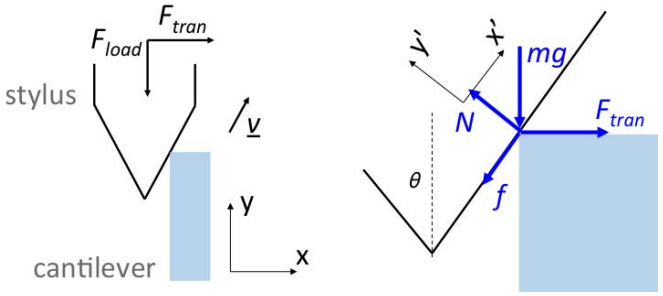

Fig. 7. Configuration of the stylus profiler and the components of applied force.

light propagation and off to one side of the fiber, such that it did not touch the fiber.

The following treatment assumes that the cantilever was level and the stylus does not deflect or twist to a significant degree. The velocity of translation is assumed to be constant and therefore the system to have no net force.

The two forces acting on the stylus are the vertical load, F_{load}, and horizontal applied load, F_{tran}, that enabled a constant positive velocity, $\underline{v}$, along the x' component. For each applied load the velocity was kept at a constant 10 μm/s. To vary the force in the lateral component the preset values of vertical load were varied. The spectral data for this was taken and interpreted into cavity length changes using Masosn's rule.

Variation in optical path lengths for the air and glass cavities were interpreted through fitting Mason's rule to the reflected spectra. In this case an ASE was used for interrogation (with a bandwidth ranging from 1530 nm to 1570 nm). The spectral response from these respective cavities with respect to applied load is shown in Figure 8. As expected, the air cavity has the greatest response to applied force, with a sensitivity of 0.17 ± 0.01 pm/nN. The error bars presented in Figure 8 are calculated from a 5 data point standard error for that load.

A reduced chi-square calculation, χ^2, of 7.64 (2.d.p) was observed for the fiber cavity trend. This equates to a chi-square per degree of freedom (df=9) of 0.76 (2.d.p). Interpretation of this relates to a 57% chance of obtaining a set of measurements at least this discrepant from the model, assuming the model to be true. As they form part of the same spectrum, it is not

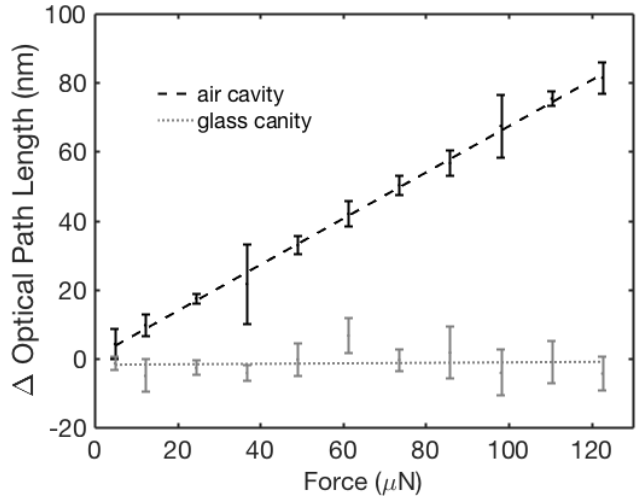

Fig. 8. Variation in optical path lengths of the air and glass cavities subject to mechanical actuation.

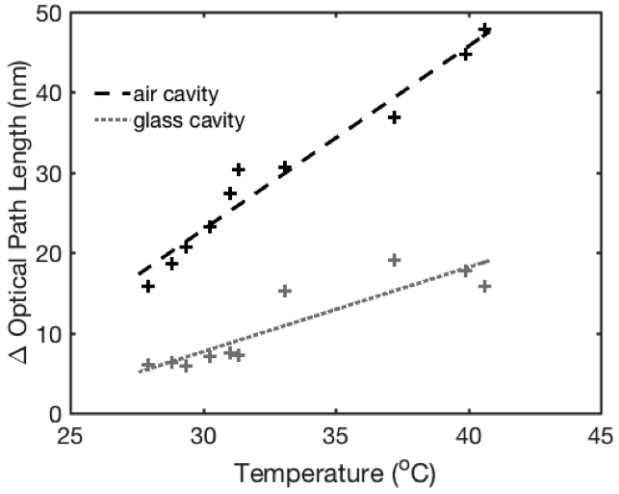

Fig. 9. Changes in optical path lengths of the air and glass cavities subject to thermal variation.

improbable that the measurements have an element of cross talk due to the nature of fitting.

An optical path length change (in the air cavity) and applied load can be interpreted in terms of the cantilever's end-tip displacement for a set lateral force. From these two values an estimation of spring constant can be made.

B. Estimating Spring Constant

The spring constant, k, for a cantilever is defined as the normal force, F_x, required for a unit normal displacement at the end-tip, Δ_x. The optical path length change in the air cavity, l_{air}, is approximately equal to the normal displacement of the cantilever's end-tip, for this considered geometry. The normal force along the x-component, N_x, can be derived as:

$$N_x = \frac{mg}{1 - \mu_k\sqrt{3}} \tag{3}$$

where μ_k is the constant for kinetic friction, approximated to be 0.115 [18] between silica (FHD layer) and diamond (profiler head), m is the set load mass and g is gravitational acceleration (9.81ms^{-2}). Using the definition of spring constant, the gradient measured in Figure 8 and Equation 3 the following can derived:

$$k = \frac{F_x}{\Delta_x} \approx \frac{1}{1 - \mu_k\sqrt{3}} \left(\frac{d\Delta_{air}}{dF_{load}}\right)^{-1} \tag{4}$$

This gives an estimation of 740 ± 40 N/m for spring constant.

Considering the simplified geometry of a purely silicon cantilever beam, without the optical fiber or FHD element. An analytical solution using the equation:

$$k = \frac{Ewt^3}{4L^3} \tag{5}$$

can be used, where Young's modulus, E, for silicon is taken to be 130 GPa; the width of the cantilever, w, 501 ± 1 μm; the thickness of the cantilever, t, 65 ± 1 μm and the cantilever length, L, 806 ± 5 μm. This calculated estimation gives a spring constant, k, equal to 850 ± 40 N/m.

This simplified approximation is an overestimation of the value inferred through measurement. It is noted however that the displacements measured at the fibre core are 62.5 μm (radius of fibre) above the position at which normal force is calculated, relating to an 8% underestimation of the measured value (horizontal difference trigonometrically equal to vertical difference 62.5/806 = 8%). This corresponds to a modified measured estimate of 790 ± 50 N/m, which overlaps the analytical approximation, with that inferred through profiler actuation.

C. Thermal Calibration

The thermal response of the cantilever is shown in Figure 9. The respective cavity responses were 2.3 nm/°C and 1.1 nm/°C for the air and glass cavities respectively.

It is noted that the glass cavity has a significantly smaller thermal response than the air cavity. The reason for this is two-fold. Firstly the air cavity is dominated by the thermal expansion of the underlying silicon, which is approximately an order of magnitude greater than that of silica. Secondly the change in optical path length brought about through thermal expansion is opposed by silica's strain-optic coefficient.

V. Discussion

Mason's rule enables an approach for multi-cavity analysis. Most authors simulate FP cavities by either analyzing first- and second-order reflections [19], or through use of the S/T-matrix approach that inherently calculates the infinite amount of reflections occurring at the boundaries [20]. Mason's method arguably sits between these, enabling effective calculation of the visibility factor of the cavities. It needs to be stressed that the interrogation method presented here is intended only to be a tool for spectral and physical calibration. Dynamic mode methods shall be the consideration of future work.

The empirical estimation for spring constant was 790 ± 50 N/m, which was comparable to that expected for a cantilever of this dimension. This may be considered slightly stiff for some applications, such as those associated with AFM, where typically spring constants range between 300 N/m to 0.01 N/m. It should be noted that the spring constant of such a cantilever could feasibly be reduced further through fabrication if required. For example, considering

Equation 5, thickness and length have a significant influence on the value. Using a combination of a thicker wafer and a blade capable of a deeper depth of cut, dimensions of up to 2 mm are feasible, which would give a $(2/0.8)^3 \approx$ 16-fold reduction from the 0.8 mm depth of cut made in this demonstration. The thickness of the cantilever could also be reduced by half, giving a further 8-fold reduction. The width of the cantilever is 501 μm, through using a smaller diameter fiber and dicing it smaller this can be reduced to 60 μm width, giving a reduction of approximately 8-fold. Furthermore silica could be used instead of silicon, which has a Young's modulus of 73 GPa and so would approximately half the spring constant. In combination spring constants below 1 N/m are entirely feasible, which is approaching the highest performance end of commercial AFMs.

It was shown that through using two distinct optical cavities a degree of thermal compensation can be made. However, in this particular configuration thermal compensation does not have the precision of alternative approaches such as Fiber Bragg gratings (FBGs) [12], [13]. Future work will therefore consider architectures that enhance this thermal response or use the technique in combination with FBGs. The solution dictated largely by the application e.g. at elevated temperatures the use of FBGs may not be possible.

VI. Conclusions

The first demonstration of a fiber-tip cantilever in an Integrated Optical Fiber (IOF) platform has been shown. IOF is proven to be a robust method capable of withstanding the mechanical rigors of physical micromachining. The microstructured cantilever is optically monitored and thus has immunity to electromagnetic interference and considered safe in flammable environments. This is the first reported demonstration that uses only physical micromachining to achieve microstructures, all other approaches reported have been in combination with other cleanroom toolsets such as wet etching.

The use of Mason's rule to interpret multicavity spectra was demonstrated. Through the interpretation deflection monitoring of the cantilever can be made. This is a useful tool to infer spectral shifts and to estimate the spring constant if the cantilever simultaneously undergoes known force actuation. A spring constant of 790 ± 50 N/m was empirically estimated using this technique. Through further development of design parameters the stiffness could be reduced further making it comparable to AFM gold standards.

It was demonstrated that thermal variation could be inferred through monitoring spectral features associated with the glass cavity. However, this may only be of significance over other techniques at elevated temperatures.

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Christopher Holmes received a master's (First Class) degree in physics from the University of Warwick in 2004, after a period in industry, returned to academia and received the Ph.D. degree with the University of Southampton in 2010. He is currently a Senior Enterprise Fellow at the Optoelectronics Research Centre, University of Southampton. His thesis was on silica integrated optics and CASE supported by the photonics company Stratophase Ltd. He continues to engage with industry and is currently the Theme Leader with the University of Southampton multidisciplinary initiative MENSUS (Monitoring of Engineered and Natural Systems Using Sensors).

Alexander Jantzen received the master's (First Class) degree in physics from the University of Southampton in 2015. He is currently pursuing the Ph.D. degree with the Optoelectronics Research Centre. His research interests are focused on designing, fabricating, and developing integrated micro opto-mechanical system components for sensing in harsh environments with a focus on silica devices aimed at aerospace use.

Alan C. Gray received the B.Sc. degree in experimental physics from Maynooth University in 2014, and the M.Sc. degree in applied physics from the University of Limerick in 2015. He is currently a Postgraduate Researcher with the Optical Engineering and Quantum Photonics Group, Optoelectronics Research Centre, University of Southampton. His research interests are in integrated photonic structures for quantum applications.

Lewis G. Carpenter received the bachelor's (First Class) degree in electrical engineering from the University of Southampton in 2009, and the Ph.D. degree from the Optoelectronics Research Centre (ORC), University of Southampton, in 2013, under the supervision of Prof. Smith and Dr. Gates. His Ph.D. focused on precision dicing (sawing) and micromilling techniques for silica photonic device fabrication, and particularly for the production of micro optomechanical sensors and ultra-smooth micron scale features for integrated optics. Working as a Research Fellow in 2014, he researched the integration of precision micromilled structures with photonic waveguides for the application to optofluidic evanescent field sensing. He is currently a Knowledge Transfer Partnership (KTP) Associate with Covesion Ltd., and ORC. Currently, as a KTP Associate, he develops ridge waveguides in periodically poled lithium niobate for high efficient wavelength conversion, with special interest in the application of these technologies to atom/ion trapping and single photon sources. He received the National Grid Award for the highest awarded dissertation mark.

Paul C. Gow received the Ph.D. degree in physics from the University of Southampton in 2016. He is currently with the Optical Engineering and Quantum Photonics Group, Optoelectronics Research Centre. His main research interests are in enabling quantum technologies, direct UV written planar photonics devices, and laser machining of silica photonic structures.

Stephen G. Lynch received the B.Eng. degree in electronic engineering from the University of Southampton in 2012. He is currently a Research Fellow with the Optoelectronics Research Centre. He is currently investigating planar Bragg-grating stabilized lasers and locking them to stable external resonators using novel feedback control techniques for his research.

James C. Gates received the M.Phys. degree and the Ph.D. degree in optoelectronics from the University of Southampton, in 1999 and 2003, respectively. His Ph.D. dissertation investigated the optical properties of various photonic devices at a nanometer scale using an interferometric near-field technique. He is currently a Principle Research Fellow with the Optoelectronics Research Centre (ORC). After his Ph.D., he moved back to the Department of Physics, Southampton, and worked in the field of nanophotonics, before returning to the ORC in 2006. He currently develops fabrication techniques for manufacturing photonic devices, principally planar-integrated optical devices. He has worked on applications in sensing, telecommunications, and quantum optics.

Peter G. R. Smith received the B.A. degree in physics and the D.Phil. degree in nonlinear optics from Oxford University, in 1990 and 1993, respectively. After a year spent as a Management Consultant, he joined the University of Southampton, where he is currently a Professor of Electronics and Computer Science with the Optoelectronics Research Centre. He has worked on a number of areas in optics research ranging from laser spectroscopy to polymer integrated-optics. He has published over 180 journal and conference papers in the fields of periodically poled materials and UV written devices at major conferences, including invited talks at national and international meetings.