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# Fermi Velocity and Effective Mass Variations in ZGaN Ribbons: Influence of Li-Passivation

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**ABSTRACT** The paper presents the structural stability and electronic properties of Zigzag Gallium Nitride nano ribbons (ZGaNNR) by considering the lithium (Li) atom by employing density functional theory (DFT). Li atom has been considered as a passivating element at various symmetric sites. By using Li atoms, a significant impact has been observed on the structural and electronic characteristics of ZGaNNRs. Bare@edges\_both structure emerged to be the most energetically stable among other structures. For Li-passivation@edge\_Ga structures, the minimum band gap has been noticed for III-V group family of nanoribbons. Interestingly, other structures of ZGaNNRs turn metallic nature irrespective of the Li site. Further, Li-bare@edge\_N structure possesses the highest Fermi velocity as compared to other structures. This is useful for designing high speed interconnect applications. Further, we investigated the effective mass of various Li-ZGaNNR structures using standard two probe models. The effective mass of H-bare@edge\_N structure reveals the highest effective mass in both valence and conduction bands. The proposed work proves the high capability towards the designing of the nano-scale devices.

**INDEX TERMS** GaNNR, passivation, Fermi velocity, effective mass.

## I. INTRODUCTION

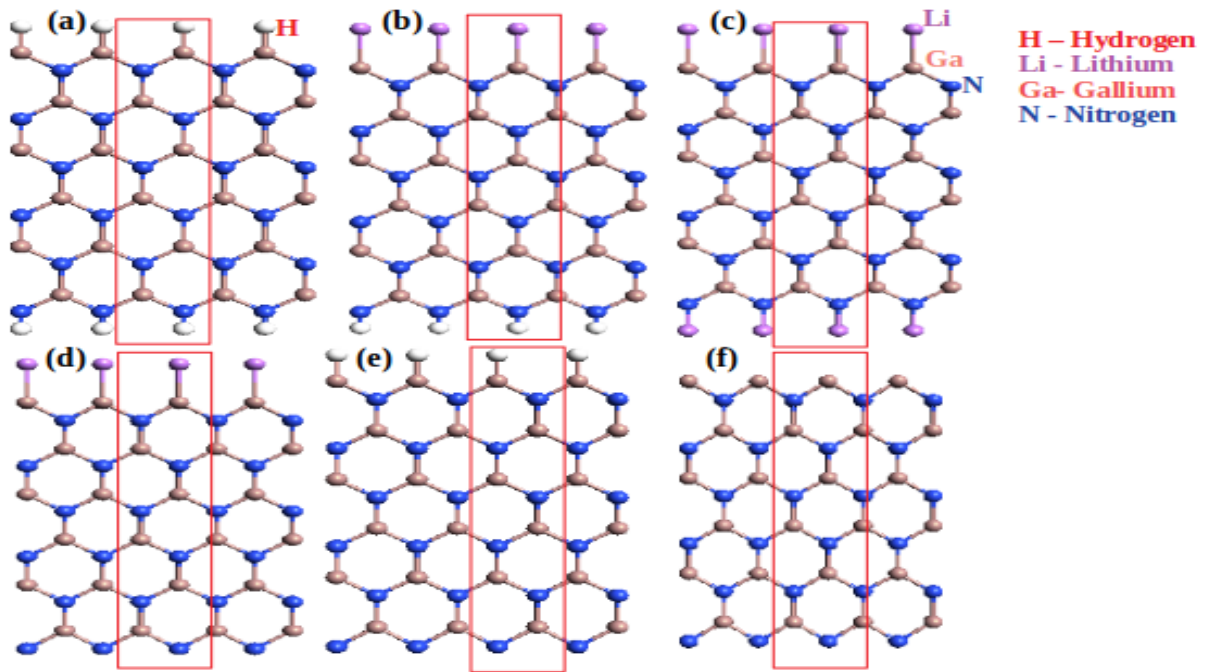
Considering the III generation era of semiconductor devices in the nano-scale regime, mainly nitrides (GaN, AlN and BN) have gained both technological and scientific attention in recent years [1], [2]. Many researchers explored the Gallium nitride nanoribbon (GaNNR) as a replacement of CMOS technology due to its wider band gap, steep reverse breakdown voltage, good thermal and chemical stability [3]–[10]. Yogi *et al.* investigated the electronic properties of Cl decorated AGaNNRs [11]. Switching phenomena and effect of NDR have been examined by Inge *et al.* using ZGaNNR by edge fluorination [12]. Ismil *et al.* and Xu *et al.* studied optical, electronic properties and also synthesized GaN [13], [14]. Sanjay *et al.* used the gallium particle trapping effect to synthesize GaN nanoparticles [15]. GaN nanoparticles are synthesized by using pulsed laser deposition, nonthermal plasma method and ammonobasic reactive sublimation [16], [17]. Thermal conductivity, phonon, analysis of charge states and noise properties have been investigated using GaNNRs [18]–[20]. Power amplifier, gunn diode, high temperature

applications have been designed using GaN [21]–[23]. Mohamed Azize *et al.* designed high electron mobility transistors using GaNNR [24]. Mina Rais *et al.* used GaN as an electromechanical material [25]. Other GaN nanocomposites such as nanoribbons [26], [27], nanobelts [28], and nanotubes [29] have been synthesized. Electronic properties are studied with edge hydrogenation for graphene nanoribbons [30]. Spin states are used for ZGNR to finding out the opto-mechanical properties [31]. Cu adsorption at different sites of AGNR structures exhibit metallic nature [32]. The qualitative characteristics of the band structures have been altered due to fluorine (F) edge functionalization [33]. The band gaps in graphene nanoribbons can be modulated with passivation techniques, and would also turn some structures metallic in nature [34]. On the other hand, exploring the effective mass and Fermi velocity of monolayer ZGaN would be useful for exploring nano-scale semiconductor devices.

## II. COMPUTATIONAL DETAIL

The paper investigated the structural and electronic properties of various Li-ZGaNNRs using first-principle methods by employing DFT. ATK-VNL tool is used for calculating all

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**FIGURE 1.** The schematic view of various Li-ZGaNNRs structures of (a) Pristine (b) Li-passivation@edge\_Ga (c) Li-passivation@edges\_both (d) Li-bare@edge\_N (e) H-bare@edge\_N and (f) Bare@edges\_both structures with  $N_z=6$ .

parameters of ZGaNNRs [35]. Nanoribbons with the width  $N_z=6$  is considered to examine the effect of Li interactions. Local-density approximations (LDA) is chosen to include the exchange-correlation [36]. For the geometrical optimization, the mesh cut-off energy of 100 Ry and  $1 \times 1 \times 100$  k-points sampling is used for the periodicity in the Brillouin zone. The transport characteristics are investigated under the NEGF framework. The current is obtained using a two probe model using Landauer Buttiker formula [37].

### III. RESULTS AND DISCUSSION

#### A. STRUCTURAL PROPERTIES

Various Li-passivation ZGaNNR structures are considered to analyse the thermostatic stability of GaNNRs. Figure 1(a)-(f) depict the optimized geometries of considered structures. All considered Li-passivation ZGaNNRs undergo structural changes after optimization and also noticed the change in bond length between the atoms. Table 1 summarizes the reduction or elongation of a bond length. In optimization,

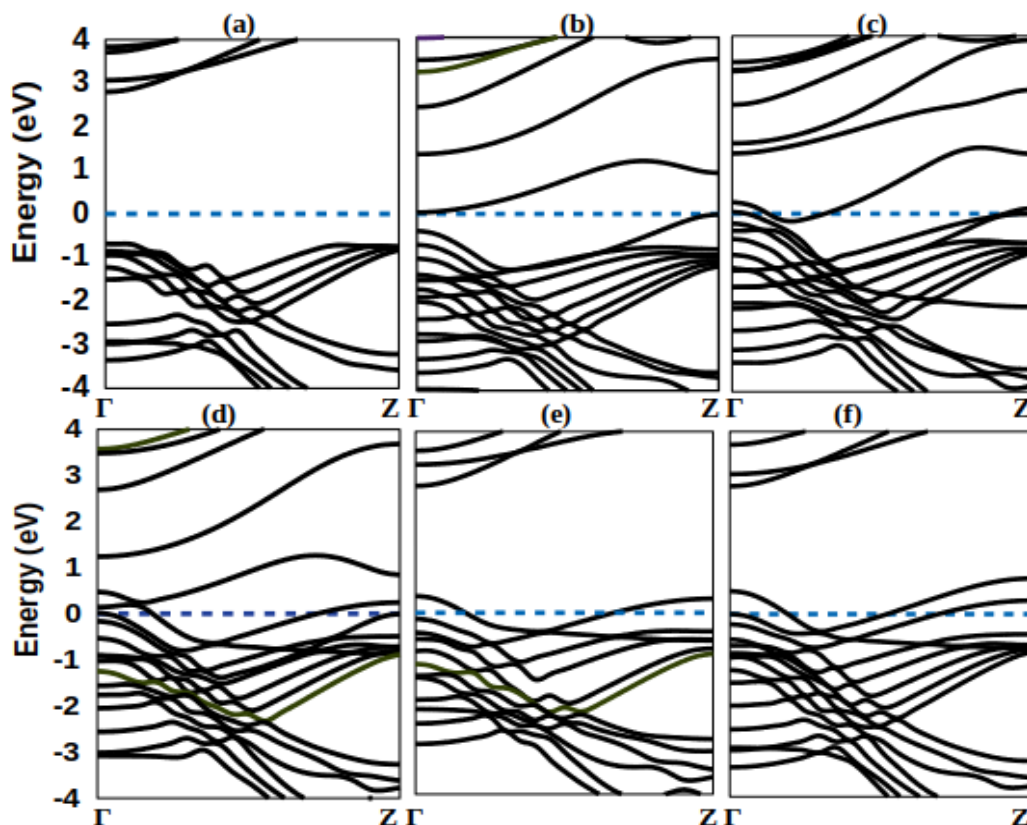
**TABLE 1.** The optimized bond length(before optimization) of the atoms for the Li-passivated ZGaNNRs. (All units are in Å).

Configurations	Ga-N	Ga-Li	N-Li	N-H	Ga-H
Pristine	1.81(1.61)	-	-	1.107(1.1)	1.54(1.1)
Li-passivation @edge_Ga	1.86(1.61)	2.48(1.6)	-	1.108(1.1)	-
Li-passivation @edges_both	1.87(1.61)	2.64(1.6)	1.78(1.6)	-	-
Li-bare@edge_N	1.85(1.61)	2.47(1.6)	-	-	-
H-bare@edge_N	1.81(1.61)	-	-	1.107(1.1)	1.54(1.1)
Bare@edges_both	1.84(1.61)	-	-	-	-

bond length analyses reveal that the highest edge restoration occurs at the ribbon's Ga edge with Li atom and it is increased from 1.6Å to 2.64Å for the  $N_z=6$ . To investigate structural stability, we examined the binding energy ( $E_{bi}$ ) per Li atom of considered structures. Table 2 summarizes the  $E_{bi}$  calculations.  $E_{bi}$  has been analyzed using a  $E_{bi} = [E_t - n_{Ga}E_{Ga} - n_N E_N - n_H E_H - n_{Li} E_{Li}] / N_t$ , where  $E_t$ ,  $E_{Ga}$  implies the total energy of Li-ZGaNNRs and single isolated Ga atom and  $E_{Li}$ ,  $E_H$  and  $E_N$  signifies the single atom energy of Li, hydrogen and nitrogen atoms respectively.  $n_{Ga}$ ,  $n_N$ ,  $n_H$ ,  $n_{Li}$  are the number of atoms of the respective energy. Further  $N_t$  is the total number of atoms in considered ribbon. We have varied the ZGaNNR width from 2 to 6 of various Li-ZGaNNRs to investigate binding energy and band gap. It is shown in Table 2 and Table 3. From Table 2, it is observed that width  $N_z=6$  structures possesses

**TABLE 2.** The Calculated binding energy ( $E_{bi}$ ) for the various Li-passivated ZGaNNRs.

Configurations	$E_{bi}(eV)$				
	2	3	4	5	6
Pristine	-5.49	-5.75	-5.99	-6.15	-6.25
Li-passivation @edge_Ga	-5.12	-5.48	-5.70	-5.91	-6.09
Li-passivation @edges_both	-4.74	-4.84	-5.19	-5.47	-5.86
Li-bare@edge_N	-4.88	-5.36	-5.63	-5.81	-5.99
H-bare@edge_N	-5.30	-5.67	-5.86	-5.99	-6.19
Bare@edges_both	-5.47	-5.60	-6.02	-5.99	-6.25



**FIGURE 2.** The band structures of various Li-ZGaN ribbon structures of (a) Pristine, (b) Li-passivation@edge\_Ga, (c) Li-passivation@edges\_both, (d) Li-bare@edge\_N, (e) H-bare@edge\_N and (f) Bare@edges\_both structures with  $N_z=6$ .

**TABLE 3.** The Calculated band gap ( $E_g$ ) for the various Li-passivated ZGaN ribbons, here M implies the metallic nature of that structure.

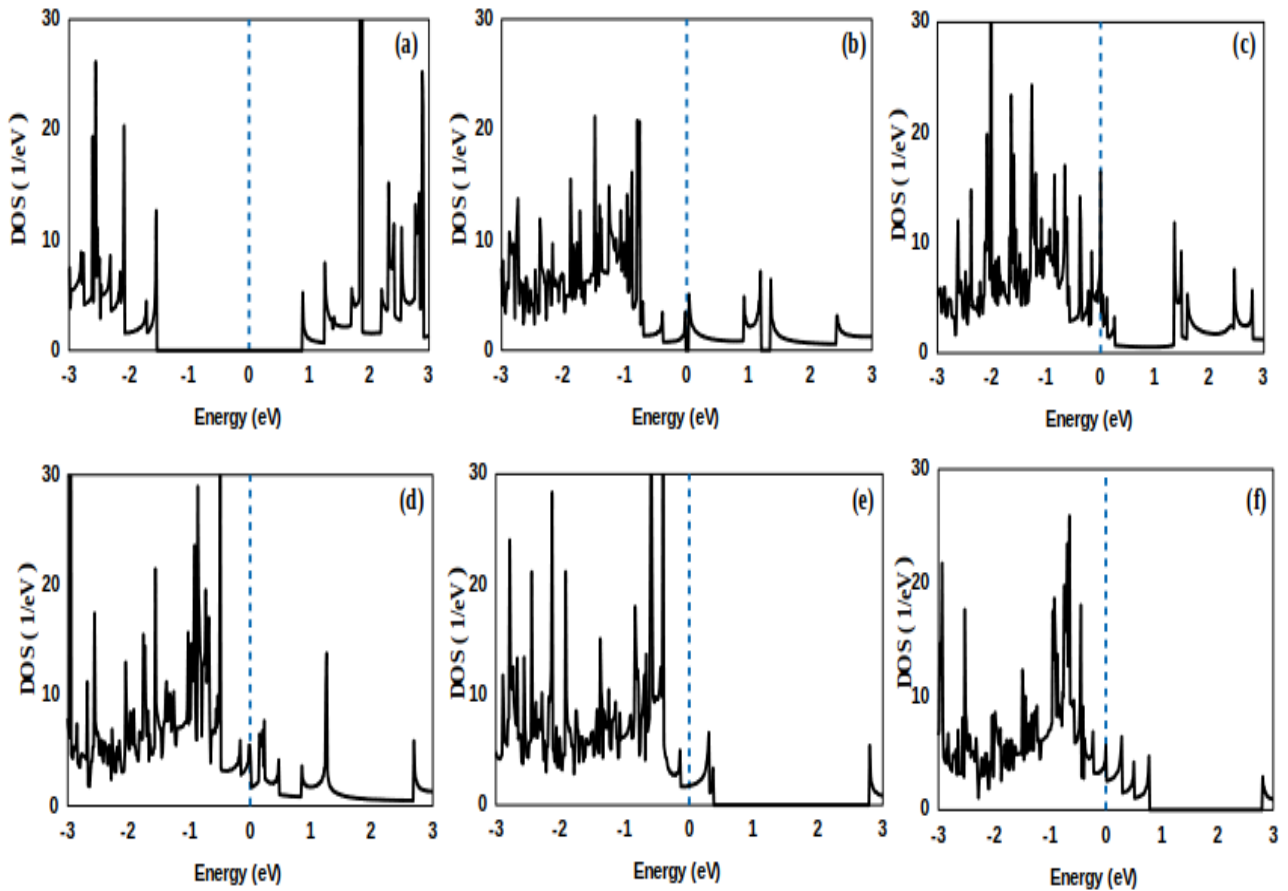
Configurations	$E_g$ (eV)				
	2	3	4	5	6
Pristine	2.92	4.67	2.64	2.62	3.48
Li-passivation @edge_Ga	M	M	M	M	0.043
Li-passivation @edges_both	M	M	M	M	M
Li-bare@edge_N	M	M	M	M	M
H-bare@edge_N	M	M	M	M	M
Bare@edges_both	M	M	M	M	M

a more negative value compared with width  $N_z=2,3,4,5$ . This negative value indicates that all unit cells are structurally feasible. Bare@edges\_both ZGaN ribbon emerged to be the most thermodynamically stable among other structures. The tabulated ZGaN ribbon binding energy ( $E_b$ ) indicates the ordering of Bare@edges\_both > Pristine > H-bare@edge\_N > Li-passivation@edge\_Ga > Li-bare@edge\_N > Li-passivation@edges\_both structures. The study shows that, the binding energies increase with increasing width of the ZGaN ribbon structures and it implies that by expanding the GaNNR width, greater stability can be achieved.

Furthermore, Li-passivation@edge\_Ga structure emerged as the most stable structure among other Li-ZGaN ribbons.

### B. ELECTRONIC PROPERTIES

The electronic properties of Li-ZGaN ribbons are analyzed using E-k diagrams and DOS profiles. As seen in Figure 2(f), the metallic behavior is observed in Bare@edges\_both ZGaN ribbon due to partially induced electronic bands and these bands are arising from dangling bonds on two side edges of the ZGaN ribbon. For Li-passivation@edges\_both ZGaN ribbon, the chemical bond is formed in between Li and the edge of N/Ga atoms that saturates one dangling bond on the edge of Ga and N atom. The residual dangling bond at the Ga and N edges are present at the Fermi level as seen in Figure 2(b). In this structure, only a single electronic band passes at the Fermi level showing metallic behaviour. Among all E-k diagrams, pristine and Li-passivation@edge\_Ga ZGaN ribbons show the semiconductor nature. The remaining structures show metallic behavior due to dangling bonds at the edge of Ga and N atoms. The pristine ZGaN ribbon possesses a direct intrinsic band gap ( $E_g$ ) that behaves as semiconducting in nature and a large bandgap behaviour (3.486 eV) is observed in band structure. An E-k diagram and their DOS profile is calculated as shown in Figure 2(a) and 3(a). The indirect band gap is observed 0.043 eV in the case of Li-passivation@edge\_Ga ZGaN ribbon



**FIGURE 3.** DOS profiles of various Li-ZGaNRR structures of (a) Pristine, (b) Li-passivation@edge\_Ga, (c) Li-passivation@edges\_both, (d) Li-bare@edge\_N, (e) H-bare@edge\_N and (f) Bare@edges\_both structures with  $N_z=6$ .

which clearly indicates the semi-conducting behaviour as shown Figure 2(b). This is due to  $\pi$  and  $\pi^*$  bonds of N and Ga atoms. The closing of the bandgap for Li-bare@edge\_N and H-bare@edge\_N structures are due to edge dangling bonds at the sites of N and Ga atoms and also noticed the presence of electronic band states around the zero Fermi level. The calculated E-k diagrams and their DOS profiles are shown in Figure 2(d-e) and Figure 3(d-e).

**C. VARIATION OF FERMİ VELOCITY AND EFFECTIVE MASS**

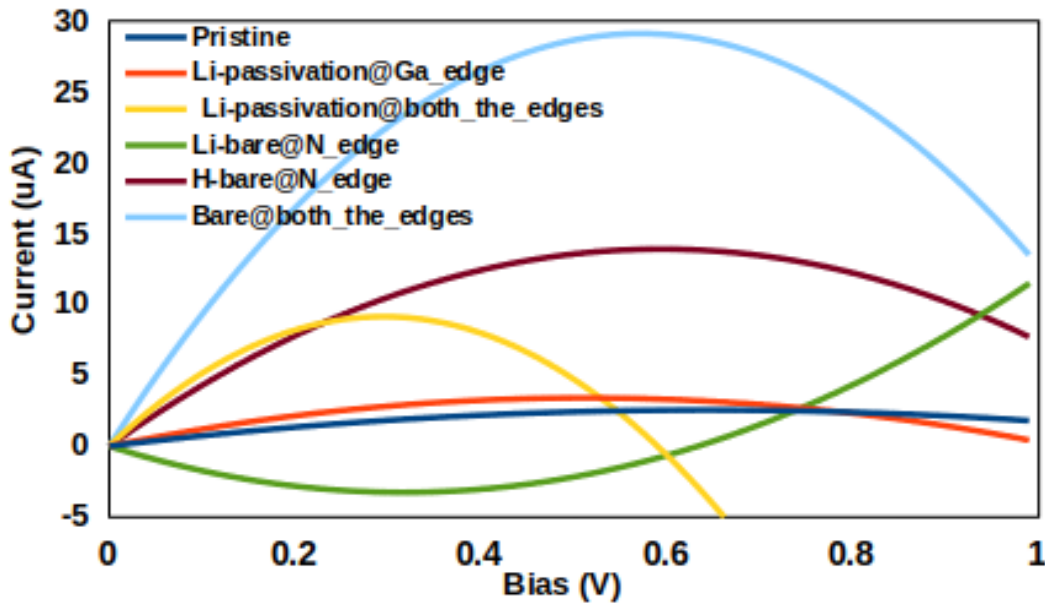
A very simple way to calculate the Fermi velocity in a Dirac material is using the energy dispersion relation. The conductivity is independent of the band slope value at the Fermi level due to linear dispersion near the origin of ribbons. Another way to detect the change in the Fermi velocity is to apply Li-passivation on ZGaNRRs and also inject carriers from Li atoms to GaNRRs. With this, the conductance should increase as a function of the applied Li passivation due to the wave vector mismatch. When the position of the Li atom changes the Fermi velocity also changes. The Fermi velocity is estimated using the standard two probe models [38].

$$V_f = \frac{1}{\hbar} \frac{dE}{dk} \tag{1}$$

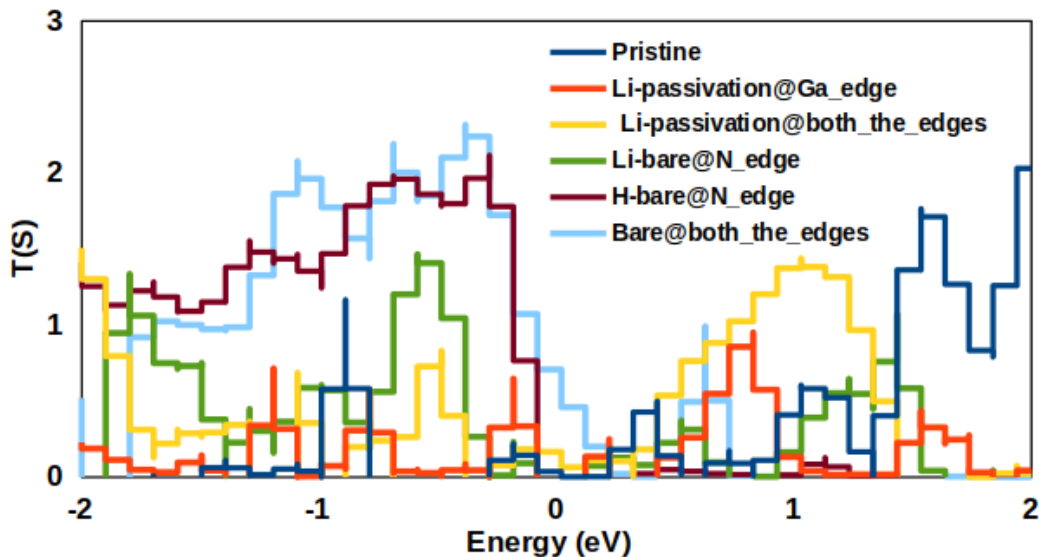
**TABLE 4.** The calculated Fermi velocity ( $V_f$ ), effective mass ( $m^*/m_0$ ) for the Li-passivated ZGaNRRs.

Configurations	$V_f$ (m/s)	$m^*/m_0$	
		electron	hole
Pristine	$6.87 \times 10^4$	0.223	0.123
Li-passivation @edge_Ga	$7.75 \times 10^4$	0.218	0.078
Li-passivation @edges_both	$7.98 \times 10^4$	0.070	0.082
Li-bare@edge_N	$8.04 \times 10^4$	0.100	0.089
H-bare@edge_N	$6.29 \times 10^4$	0.107	0.282
Bare@edges_both	$3.92 \times 10^4$	0.135	0.221

where  $\hbar = h/2\pi$ ,  $\hbar$  implies the reduced Planck’s constant and  $h$  is Planck’s constant, value of  $h$  is  $6.62 \times 10^{-34}$  Js in SI units. The center values of energy and crystal momentum are denoted by  $E$  and  $k$  respectively. The tabulated ZGaNRR Fermi velocity ( $V_f$ ) indicates the ordering of Li-bare@edge\_N>Li-passivation@edges\_both>Li-passivation@edge\_Ga>Pristine>H-bare@edge\_N>Bare@edges\_both structures. It can be seen from Table 4, the



**FIGURE 4.** I-V characteristics of Pristine, Li-passivation@ edge\_Ga, Li-passivation@edges\_both, Li-bare@edge\_N, H-bare@edge\_N and Bare@edges\_both structures of ZGaNNRs.



**FIGURE 5.** Calculated transmission spectra characteristics of Pristine, Li-passivation@edge\_Ga, Li-passivation@edges\_both, Li-bare@edge\_N, H-bare@edge\_N and Bare@edges\_both structures of ZGaNNRs.

Fermi velocity ( $V_f$ ) increases with increasing Li atoms in the ZGaNNR structures. Li-bare@edge\_N structure possesses the highest  $V_f$  ( $8.04 \times 10^4$  m/s) as compared with other structures. This is useful for designing high speed interconnect applications. Generally, Fermi velocity is inversely proportional to the quantum capacitance and kinetic inductance for designing high speed interconnect applications. In order to get the low inductive coupling and capacitive delay, the Fermi velocity should be high. It can be noticed from Table.4 that Li-bare@edge\_N structure is the strong contender for designing nanoscale interconnect circuits. The characteristic

lateral confinement of charge carriers in ZGaNNRs causes a non-zero effective charge carrier mass, which impacts on charge carrier mobility in these devices. Such a system emphasizes the importance of charge transport in order to promote future GaN based technology. Greater effective masses are observed due to the lattice Li passivation that induces quasi-particles, resulting in decreased charge mobility. In this regard, the carrier's effective mass of ZGaNNR is a main factor to design and improve GaN-based devices. Further, we investigated the effective mass ( $m^*/m_0$ ) of various Li-ZGaNNR structures using standard two probe model.

The calculated effective mass ( $m^*/m_0$ ) for the Li-passivated ZGaNRRs shown in Table 4. The effective mass of H-bare@edge\_N structure reveals the highest effective mass in both valence (0.282) and conduction (0.107) bands. This helps to design the future low bias nano-scale semiconductor devices. The  $m^*$  is calculated using a relation of energy and momentum is obtained from (2)

$$\frac{1}{m^*} = \frac{1}{\hbar^2} \frac{d^2E}{dk^2} \quad (2)$$

#### D. TRANSPORT PROPERTIES

In this section, the transport characteristics of various Li-passivated-ZGaNRRs with  $N_z = 6$  are investigated using standard two probe models. Evaluated using self-consistent methods and used bias sampling step size of 0.1V for 0 to 1V. For pristine ZGaNRR, the current-voltage curve clearly indicates that the current is rapidly enhanced during low bias, reaching a maximum value at 1.0V. The peak current around  $2.5\mu\text{A}$  reaches with a voltage of 1.0V. This linearity is due to good ohmic interaction between Ga-N bonds. Further, we found Li-bare@edge\_N structure is not conducting current at initial voltage up to 0.7V and afterwards its amplitude rises sharply. Li-passivation@edge\_Ga structure conducts less current due to less numbers of H-atom presence in the scattering region. Li-passivation@edges\_both, H-bare@edge\_N and Bare@edges\_both structures of ZGaNRRs exhibit NDR behaviour shown in Figure 4. The current-voltage characteristics of Bare@edges\_both ZGaNRR shows that the current increased rapidly during initial positive bias, reaches at a bias of 0.6V and then decreases with an increase in voltage, showing NDR behavior. The peak current of  $29.8\mu\text{A}$  was found at 0.6V and then current begins to decrease beyond 0.6V. Understanding the effect of NDR, by analyzing transmission curves, varies as a result of the voltage applied. At 0.6V, the bias window reveals a peak propagation direction with a current of approximately  $29.8\mu\text{A}$ . Decreasing in current has been observed during a reduction in electron flow in the model. The NDR effect of the proposed model could be utilized in various applications. Further, we investigated the I-V characteristics of the H-bare@edge\_N ZGaNRR. Current is enhanced at an initial positive bias, reached peak at 0.61V and then gradually decreases with a rise in voltage, shows a NDR effect. Further Li-passivation@edges\_both ZGaNRR structure increases its current at the initial stage at 0.3V and then current decreases with a rise in voltage, indicating a NDR behavior. Figure 4 and 5 depict i-v characteristics and transmission spectra of considered Li-passivated ZGaNRRs.

#### IV. CONCLUSION

In this work, investigations of various Li-passivated ZGaNRRs have been carried out by employing DFT frames. The band structures and DOS of all ZGaNRRs configurations demonstrate metallic nature and can be tuned by taking Li-passivation at different edges. Li passivation ZGaNRRs

reduces the electronic band gap as compared to pristine ones. Based on binding energy calculations, Bare@edges\_both structure to provide found the most stable configuration as compared to other ZGaNRRs. For Li-passivation@edge\_Ga structures, the minimum band gap (0.043eV) has been observed for III-V group family of nanoribbons. Further, Li-bare@edge\_N structure possesses the highest Fermi velocity ( $8.04 \times 10^4 \text{m/s}$ ). This makes it suitable for designing high speed interconnect applications. Further, we investigated the effective mass ( $m^*/m_0$ ) of various Li-ZGaNRR structures using standard two probe models with the help of E-k relationships. The effective mass of H-bare@edge\_N structure reveals the highest effective mass in both valence (0.282) and conduction (0.107) bands. This helps to design the future low bias nano-scale semiconductor devices.

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