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# Analytical Model of Torque-Prediction for a Novel Hybrid Rotor Permanent Magnet Machines

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**ABSTRACT** This paper presents an analytical method for predicting the electromagnetic performance in permanent magnet (PM) machine with the spoke-type rotor (STR) and a proposed hybrid rotor structure (HRS), respectively. The key of this method is to combine magnetic field analysis model (MFAM) with the magnetic equivalent circuit model. The influence of the irregular PM shape is considered by the segmentation calculation. To obtain the boundary condition in the MFAM, respectively, two equivalent methods on the rotor side are proposed. In the STR, the average flux density of the rotor core outer-surface is calculated to solve the Laplace's equation with considering for the rotor core outer-surface eccentric. In the HRS, based on the Thevenin's theorem, the equivalent parameters of PM remanence  $B_{reB}$  and thickness  $h_{pme}$  are obtained as a given condition, which can be utilized to compute the air-gap flux density by conventional classic magnetic field analysis model of surface-mounted PMs with air-gap region. Finally, the proposed analytical models are verified by the finite element analysis (FEA) with comparisons of the air-gap flux density, flux linkage, back-EMF and electromagnetic torque, respectively. Furthermore, the performance that the machine with the proposed hybrid structure rotor can improve the torque density as explained.

**INDEX TERMS** Analytical method, hybrid PM rotor structure, rotor core outer-surface eccentric, magnetic equivalent circuit model (MECM).

## I. INTRODUCTION

Recently, interior permanent magnet synchronous machine (IPMSM) has been widely applied to the electric vehicles (EVs) field and aerospace industry all over the world, in view of the excellent performances of high-power density, high efficiency, and high reliability [1]–[5]. To further improve the torque density, the spoke-type PM machines has been adopted due to the advantage of the flux concentrated structure [6]–[8]. In some applications, such as high dynamic response machine with low moment of inertia, rotor outer diameter is strictly restricted. Therefore, the spoke-type IPMSMs are significantly effective because it can increase magnetic flux area of IPMSMs from the larger magnet surface area compared with other types, which means that the larger

torque can be obtained [9]–[11]. Moreover, consequent-Pole and Hybrid-Pole PM Machines are adopted to improve the output torque and efficiency. Then, the flux barriers are employed to improve the magnetic symmetry and reduce the saturation of the stator yoke [12]. Two novel consequent-pole PM machines with different pole sequences are adopted. In order to provide assistant flux and suppress the subharmonics, the tangential PMs are embedded into the rotors to form the novel hybrid-pole PM machines [13].

The analytical methods can offer computation speed and flexibility at the initial design and optimization process [14], [15]. There are many related research literatures on magnetic field analysis, such as conventional scalar magnetic field analysis [16], [17] and subdomain magnetic field analysis [18]–[20]. However, in many situations, the boundary conditions are not ideal that the geometric model is irregular, which is result in inaccurate prediction of electromagnetic

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performances compared with the FEA. In order to map the spoke PMs in polar coordinates, the rectangular geometry with rectangular spoke PMs assisted by hub magnets is approximated by a combination of arc shaped subdomains [21]. To match the standard boundary conditions of the Poisson and Laplace equations, the trapezoid magnet is simplified to sector magnets, the dimension of the sector magnet is calculated with a lumped magnetic circuit model [7]. In [22], it presents an improved analytical method for the surface mounted PM machines accounting for any eccentric rotor shape by the equivalent surface current of magnet pole. Meanwhile, taking into account the saturation and nonlinearity of magnetic material, the magnetic bridges are equivalent to fan-shaped saturation regions [23], [24]. For getting standard boundary conditions, a lumped parameter magnetic circuit model and iterative method are employed to calculate the permeability [25]–[27]. However, it’s hard to generalize because of the rotor structure with different shape and various PM arrangements forms.

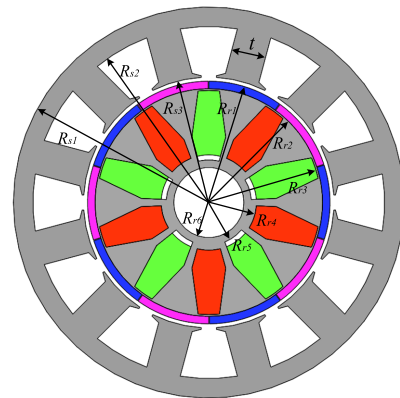
This paper presents an analytical method for the PM machine with non-ideal regular shape rotor structure, which considers the rotor core outer-surface eccentric and the proposed hybrid rotor structure, respectively. The proposed analytical model takes advantage of the characteristics of magnetic field analysis model and the magnetic equivalent circuit model. The corresponding boundary condition in the STR and the equivalent parameters in the HRS are obtained by the MECM, respectively. The predicted accuracy of the proposed two analytical models is validated by FEA. Moreover, the proposed hybrid rotor structure can improve torque density due to increasing the magnetic flux area of IPMSMs and the equivalent thickness of the PM. In Section II, the topology structure of the studied machine will be given. In Section III, the equivalent magnetic circuit model of rotor will be presented to obtain the boundary conditions and equivalent parameters. In Section IV, model of the air-gap magnetic field analysis in the STR and HRS are built. In Section V, analytical prediction accuracy will be verified by FEA. Finally, the conclusions will be derived in Section VI.

**II. TOPOLOGIES AND FEATURES**

The cross section of the hybrid PMSM with the Spoke-Type and surface-mounted PMs is shown in Fig 1. A 10-pole/12slot motor as the studied model is adopted to highlight the advantage of analytical method. For the sake of simplification, the stator is assumed to be infinitely permeable. It is assumed that the permeability of all materials in machine is isotropic and homogenous. The effects of the axial end and eddy current in the magnetic material are neglected [28].

The embedded PMs are tangentially magnetized and surface-mounted PMs are parallel magnetized with a linear demagnetization characteristic, respectively.

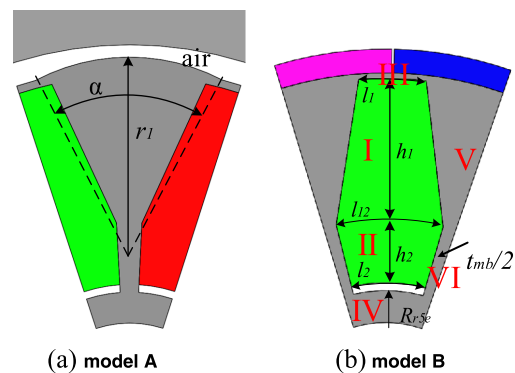
The main design parameters of the Spoke-Type rotor (model A) and hybrid rotor structure (model B) IPMSM are listed in Table 1:



**FIGURE 1. Initial machine model of 10-pole/12slot with hybrid structure rotor.**

**TABLE 1. Main PM machine parameters.**

Symbol	Quantity	Value and unit
$p$	Number of pole-pairs	5
$Q_s$	Number of stator slots	12
$\omega_r$	Rated speed	1500rpm
$I_r$	Rated current	25A
$B_r$	Magnet remanence	0.4T
$L_{ef}$	Armature effective length	100mm
$R_{s1}$	Radius of the stator outer surface	39mm
$R_{s3}$	Radius of the stator inner surface	25mm
$N$	Number of turns per phase	48
$t$	Tooth width	6.9mm



**FIGURE 2. Model of the rotor structure (a) the rotor core outer-surface eccentric (b) hybrid PM rotor structure.**

Partial enlarged views of two rotor structures are shown as Fig 2. In order to highlight the details in the STR, the core is placed in the symmetrical center of model A. It is worth noting that the model B is similar to the model A when the thickness of the surface-mounted PM in the hybrid rotor structure is close to 0 mm. In order to reduce the harmonic content of the air-gap magnetic density, the eccentric structure is often adopted on the outer surface of the rotor core.

**III. ROTOR MAGNETIC CIRCUIT MODEL**

Due to the non-idealization of the rotor structure arrangement, the EMCM is calculated to obtain the rotor flux to

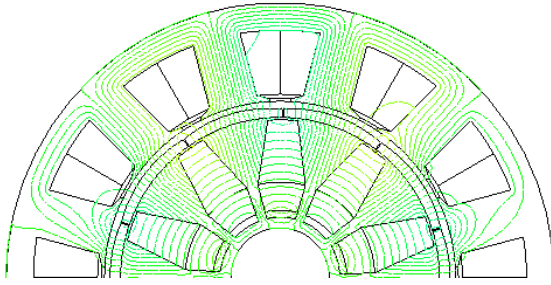


FIGURE 3. The Motor no-load magnetic line distribution.

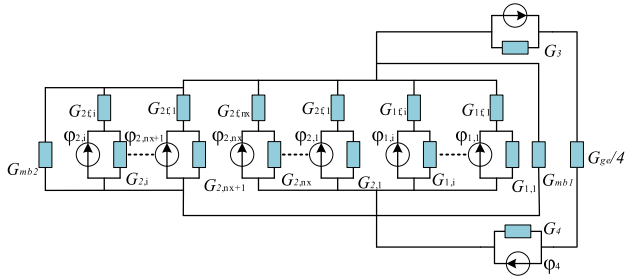


FIGURE 4. The equivalent magnetic circuit model.

air gap. For two rotor pole structures, the no-load air-gap magnetic density can be calculated by the Poisson's equation with the air-gap boundary condition in model A, and conventional classic magnetic field analysis model of equivalent surface-mounted PMs parameters with air gap in model B, respectively.

Because the magnetic circuit structures of the two models are very similar, the presented EMCM is compatible with both rotor pole structures as shown in Fig 3.

As shown in Fig 2, Fig 3 and Fig 4, it is obvious that the flux by the Spoke-Type PMs generated can be divided three parts: the main flux through the airgap, eventually, the leakage flux through the magnetic bridge of PMs top and rotor yoke, respectively. The Spoke-Type PMs is composed of two inverted trapezoids magnet, it can be replaced by  $n_1$  and  $n_2$  sector magnets in parallel.  $h_1$  and  $h_2$  are height of two inverted trapezoids magnets.  $l_1, l_{12}, l_2$  are the widths of trapezoids, respectively.  $t_{mb}$  is the magnetic bridge width between both PMs in the VI region.

The magnet flux source  $\varphi_{1,i}, \varphi_{2,i}$  and the magnetic permeance  $G_{1,i}, G_{2,i}$  of the  $i$ th magnet piece in I and II regions, respectively, can be expressed by:

$$\varphi_{1,i} = B_r \Delta h_1 L_{ef} \quad (1)$$

$$\varphi_{2,i} = B_r \Delta h_2 L_{ef} \quad (2)$$

$$G_{1,i} = \frac{\mu_0 \mu_{rm} \Delta h_1 L_{ef}}{\Delta l_{1,i}} \quad (3)$$

$$G_{2,i} = \frac{\mu_0 \mu_{rm} \Delta h_2 L_{ef}}{\Delta l_{2,i}} \quad (4)$$

where  $B_r$  is the magnet remanence,  $L_{ef}$  is the effective length of rotor and stator,  $\mu_0$  is the permeability of the vacuum,  $\mu_{rm}$

is the magnet relative permeability; and

$$\Delta h_1 = \frac{h_1}{n_1} \quad \Delta h_2 = \frac{h_2}{n_2} \quad (5)$$

$$\Delta l_{1,i} = l_1 + \frac{(l_{12} - l_1)(i - 1)}{n_1} \quad \Delta l_{2,i} = l_2 + \frac{(l_{12} - l_2)(i - 1)}{n_2} \quad (6)$$

The magnet flux source  $\varphi_3$  and the magnetic permeances  $G_3$  and  $G_4$  of the surface-mounted PMs, respectively, when considering slot less stator and the polar arc coefficient is 1, can be calculated by:

$$\varphi_3 = \frac{\pi(R_{r2} + R_{r1})}{4p} L_{ef} B_r \quad (7)$$

$$G_3 = G_4 = \mu_0 \mu_{rm} \frac{\pi(R_{r1} + R_{r2})}{4p(R_{r1} - R_{r2})} L_{ef} \quad (8)$$

where  $p$  is the number of pole pairs,  $R_{r1}$  and  $R_{r2}$  is the radius of the rotor outer surface and rotor core, respectively.

The magnetic permeance  $G_{ge}$  of air-gap can be calculated by:

$$G_{ge} = \mu_0 \frac{\pi(R_{s3} + R_{r1})}{2p(R_{s3} - R_{r1})} L_{ef} \quad (9)$$

where  $R_{s3}$  is the radius of the stator bore.

The flux density in magnetic bridge region is at a high level of saturation on the B-H curve of the silicon steel, the relative permeability  $\mu_{r1}$  and  $\mu_{r2}$  can be assumed as a small constant where the flux density is 2.1T [3], [4], [29]. The leakage magnetic permeance  $G_{mb1}$ , and  $G_{mb2}$  in III and IV regions can be calculated by:

$$G_{mb1} = \frac{\mu_0 \mu_{r1} (R_{r2} - R_{r3}) L_{ef}}{l_1} \quad (10)$$

$$G_{mb2} = \frac{\mu_0 \mu_{r2} (R_{r5e} - R_{r6}) L_{ef}}{l_2} \quad (11)$$

where  $R_{r5e}$  is the average value of  $R_{r4}$  and  $R_{r5}$ ,  $R_{r6}$  is the radius of the shaft.

The calculation method of the core magnetic permeance is the same as that of the PMs in I or II regions, it is divided into many small pieces. According to the shape of the path,  $n$  parallel computing units are divided as:

$$G_{1f,i} = \frac{\mu_0 \mu_{rf1} \Delta h_{1f} L_{ef}}{i \Delta h_1} \quad (12)$$

where

$$\Delta h_{1f} = \frac{\pi R_{r3}}{2p(n_1 + n_x)} \quad (13)$$

It is worth that the flux excited by the PMs in II region is divided into two directions: part of it is through the V region, others are through the IV region and the VI region as shown in Fig 3. The magnetic permeance  $G_{2f,i}$  of the rotor core, where the magnetic flux generated by the PMs in II region

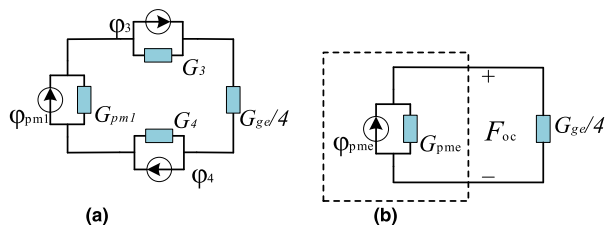


FIGURE 5. Simplified model (a) preliminary simplification (b) equivalent parameters simplification based on the Thevenin's theorem.

passes, can be calculated by:

$$\begin{cases} G_{2f,i} = \frac{\frac{\mu_0\mu_{rf1}\Delta h_2 L_{ef}}{h_1} \cdot \frac{\mu_0\mu_{rf2}\Delta h_2 L_{ef}}{n_x \Delta h_2}}{\frac{\mu_0\mu_{rf1}\Delta h_2 L_{ef}}{h_1} + \frac{\mu_0\mu_{rf2}\Delta h_2 L_{ef}}{n_x \Delta h_2}} & i \leq n_x \\ \frac{\mu_0\mu_{rf2}\Delta h_2 L_{ef}}{(n_2 - n_x - 1)\Delta h_2} & i > n_x \end{cases} \quad (14)$$

where  $h_x$  is the height of the PMs in II region that the generated flux through the main magnetic flux direction; and

$$\Delta h_{2f} = \frac{t_{mb}}{2n_2} \quad (15)$$

$$n_x = \text{round}\left(\frac{h_x n_2}{h_2}\right) \quad (16)$$

where *round* is a function for towards nearest integer.

According to the Fig.5, the magnetic flux through the rotor yoke can be expressed by:

$$\varphi_{ry} = B_{sat} A_{ry} = \sum_{n_x+1}^{n_2} \varphi_{2,i} \frac{G_{2f,i}}{G_{2f,i} + G_{2,i}} \quad (17)$$

where  $B_{sat}$  is the flux density 2.1T,  $A_{ry}$  is  $(R_{r5e}-R_{r6})L_{ef}$ . Then  $n_x$  can be solved by the equation (17). Also,  $h_x$  can be estimated by the equation (16).

The model can be simplified as shown in Fig 5 (a). Equivalent flux  $\varphi_{pm1}$  and permeance  $G_{pm1}$  are calculated by:

$$\begin{aligned} \varphi_{pm1} &= \sum_{i=1}^{n_1} \varphi_{1,i} \frac{G_{1f,i}}{G_{1f,i} + G_{1,i}} + \sum_{i=1}^{n_2} \varphi_{2,i} \frac{G_{2f,i}}{G_{2f,i} + G_{2,i}} \quad (18) \\ G_{pm1} &= \sum_{i=1}^{n_1} \frac{G_{1,i} \cdot G_{1f,i}}{G_{1,i} + G_{1f,i}} + \sum_{i=1}^{n_2} \frac{G_{2,i} \cdot G_{2f,i}}{G_{2,i} + G_{2f,i}} \\ &\quad + G_{mb1} + G_{mb2} \quad (19) \end{aligned}$$

By the Thevenin's theorem, the EMCM can be further simplified as shown in Fig 5 (b). Equivalent external magnetomotive force, internal magnetic permeance and flux source can be calculated by:

$$F_{oc} = \frac{\varphi_{pm1}}{G_{pm1}} + \frac{\varphi_3}{G_3} + \frac{\varphi_4}{G_4} \quad (20)$$

$$G_{pme} = \frac{G_{pm1} \cdot G_3}{(2G_{pm1} + G_3)} \quad (21)$$

$$\varphi_{pme} = F_{oc} \cdot G_{pme} \quad (22)$$

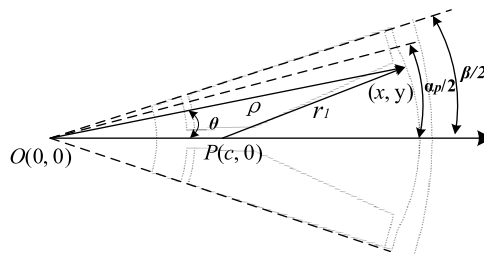


FIGURE 6. Geometric model of the rotor core outer-surface eccentric.

Therefore, average flux density on the rotor outer surface can be obtained by:

$$B_{reA} = \varphi_{pme} \frac{2G_{ge}}{(G_{ge} + 4G_{pme})A_m} \quad (23)$$

where  $A_m$  is the effectively magnetic flux area on the outer surface of the rotor.

If the EMCM is equivalent to the model that only a surface mount PM and air-gap region, then the equivalent parameters of PM remanence  $B_{reB}$  and thickness  $h_{pme}$  can be calculated by:

$$B_{reB} = \frac{\varphi_{pme}}{A_m} \quad (24)$$

$$h_{pme} = \frac{\mu_0\mu_{rm}A_m}{2G_{pme}} \quad (25)$$

#### IV. AIR-GAP MAGNETIC FIELD ANALYSIS MODEL

##### A. ROTOR CORE ECCENTRIC STRUCTURE IN THE MODEL A

Based on the simplification of the rotor magnetic circuit above all, average magnetic density on the rotor outer surface has been obtained. As Fig 6 shows, point  $O(0, 0)$ , namely origin coordinates, and  $P(c, 0)$  are the centers of the air-gap area and the surface of the rotor core, respectively.  $\alpha_p$  is polar arc coefficient.

The scalar magnetic potential distribution in the air-gap regions is governed by Laplace's equation, i.e.

$$\frac{\partial^2 \Omega}{\partial r^2} + \frac{1}{r} \frac{\partial \Omega}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \Omega}{\partial \theta^2} = 0 \quad (26)$$

According to the periodicity of the field, the general solution in the air-gap regions can be expressed by:

$$\Omega(r, \theta) = \sum_{n=1}^{\infty} (A'_n r^{np} + B'_n r^{-np}) \cos np\theta \quad (27)$$

where  $\Omega$  is related to the radial and circumferential components of  $\vec{B}$  by:

$$\begin{cases} B_r = -\mu_0 \frac{\partial \Omega}{\partial r} \\ B_\theta = -\mu_0 \frac{1}{r} \frac{\partial \Omega}{\partial \theta} \end{cases} \quad (28)$$

The boundary conditions for the interfaces between air-gaps and stator core can be expressed by:

$$B_\theta(r, \theta)|_{r=R_{si}} = 0 \quad (29)$$

In the same way, under other coordinate that the point  $P$  ( $c, 0$ ) is used as the origin coordinates, the scalar magnetic potential distribution in the air-gap regions can be built by reference to (26), (27) and (28), then, the interfaces between air-gaps and the outer surface in the rotor can be expressed by:

$$B_{r1}(r, \theta) = -\mu_0 \frac{\partial \Omega}{\partial r_1} = B_{rotor} \quad (30)$$

$$B_{rotor} = \sum_{m=1}^{\infty} \frac{4}{m\pi} B_{reA} \sin\left(\frac{m\pi\alpha}{2}\right) \quad (31)$$

where  $r_1$  is the arc radius of the surface of the rotor core,  $\alpha$  is the angle with reference to d axis.

By combining (28) and (30), the radial density can be rewritten by:

$$B_r(r, \theta) = -\mu_0 \frac{\partial \Omega}{\partial r_1} \frac{\partial r_1}{\partial r} = B_{r1}(r, \theta) \frac{\partial r_1}{\partial r} \quad (32)$$

From above boundary conditions (29), (27) can be simplified by:

$$B'_n = -A'_n R_{si}^{2np} \quad (33)$$

From above boundary condition (30), whilst combined formula (31), (27) can be rewritten by:

$$\begin{aligned} \mu_0 \sum_{n=1}^{\infty} np(A'_n r^{np-1} - B'_n r^{-np-1}) \cos np\theta \\ = \frac{\partial r_1}{\partial r} \sum_{m=1}^{\infty} \frac{4}{m\pi} B_{reA} \sin\left(\frac{m\pi\alpha}{2}\right) \end{aligned} \quad (34)$$

when  $n \neq m$ ,

$$A'_n r^{np-1} - B'_n r^{-np-1} = 0 \quad (35)$$

then  $A'_n = B'_n = 0$ , which is contradictory with (27).

Therefore, only the equation  $n = m$  is satisfied, (34) can be rewritten by:

$$-\mu_0 np(A'_n r^{np-1} - B'_n r^{-np-1}) = \frac{4}{n\pi} B_{reA} \sin\left(\frac{n\pi\alpha}{2}\right) \frac{\partial r_1}{\partial r} \quad (36)$$

Assume that  $\theta$  is the angular position referencing to in the rotor initial position;  $\theta_{mp}$  is the angle with reference to the center of a nearest magnet pole. It can be calculated by:

$$\theta_{mp} = \begin{cases} \text{mod}(\theta, \beta/2) & \text{if } k = \text{even} \\ \beta/2 - \text{mod}(\theta, \beta/2) & \text{elsewhere} \end{cases} \quad (37)$$

where  $\beta$  is the pole angle, i.e.,  $\beta = \pi/p$ ;  $k$  is the quotient that  $\theta$  and  $\beta/2$ , i.e.,  $k = \text{floor}(\theta/(\beta/2))$ ;  $\text{floor}$  represents quotient function that take an integer down;  $\text{mod}$  represents remainder function.

Assume that the coordinate  $(x, y)$  represents the position of the outer surface of rotor core, where  $\rho$  is the distance to the point  $O$ .

$$(x - c)^2 + y^2 = r_1^2 \quad (38)$$

Obviously,  $x = \rho \cos \theta_{mp}$ ,  $y = \rho \sin \theta_{mp}$ ; Substituting it to (38) can be obtained by:

$$r_1 = \sqrt{\rho^2 - 2\rho c \cos \theta_{mp} + c^2} \quad (39)$$

The partial derivative of (39) is:

$$\frac{\partial r_1}{\partial \rho} = \frac{2\rho - 2c \cos \theta_{mp}}{2\sqrt{\rho^2 - 2\rho c \cos \theta_{mp} + c^2}} = \frac{\rho - c \cos \theta_{mp}}{r_1} \quad (40)$$

Also, according to (39) and the geometric relationship as shown in Fig 6, variable  $\rho$  can be calculated by:

$$\rho(\theta) = \begin{cases} c \cos(\theta_{mp}) + \sqrt{r_1^2 - c^2 \sin^2 \theta_{mp}} & \text{if } 0 < \theta_{mp} \leq \frac{\alpha_p}{2} \\ \rho(\alpha_p/2) & \text{else } \frac{\alpha_p}{2} < \theta_{mp} \leq \frac{\beta}{2} \end{cases} \quad (41)$$

where due to  $\rho$  is close to the distance  $r$  between air-gap region and point  $O$ , the equation  $\frac{\partial r_1}{\partial r} \approx \frac{\partial r_1}{\partial \rho}$  is to be satisfied.

Therefore, according to (33) (36), coefficient  $A_n$  can be obtained by:

$$A_n = -\frac{\rho - c \cos \theta_{mp}}{\mu_0 n p r_1} \frac{4}{n\pi} \frac{B_{reA} \sin\left(\frac{n\pi\alpha}{2}\right)}{(\rho^{np-1} + R_{si}^{2np} \rho^{-np-1})} \quad (42)$$

Based on (28) (33) and (42), the radial and circumferential components of the air-gap flux density can be obtained by:

$$\begin{aligned} B_r(r, \theta) = -\frac{\rho - c \cos \theta_{mp}}{r_1} \\ \cdot \sum_{n=1}^{\infty} \frac{4}{n\pi} \frac{B_{reA} \sin\left(\frac{n\pi\alpha}{2}\right)}{\left(\left(\frac{\rho}{R_{si}}\right)^{np-1} + \left(\frac{R_{si}}{\rho}\right)^{np+1}\right)} \left(\frac{r}{R_{si}}\right)^{np-1} \\ + \left(\frac{R_{si}}{r}\right)^{np+1} \cos np(\theta - \omega_r t) \end{aligned} \quad (43)$$

$$\begin{aligned} B_\theta(r, \theta) = \frac{\rho - c \cos \theta_{mp}}{r_1} \\ \cdot \sum_{n=1}^{\infty} \frac{4}{n\pi} \frac{B_{reA} \sin\left(\frac{n\pi\alpha}{2}\right)}{\left(\left(\frac{\rho}{R_{si}}\right)^{np-1} + \left(\frac{R_{si}}{\rho}\right)^{np+1}\right)} \left(\frac{r}{R_{si}}\right)^{np-1} \\ - \left(\frac{R_{si}}{r}\right)^{np+1} \sin np(\theta - \omega_r t) \end{aligned} \quad (44)$$

## B. EQUIVALENT MODEL BASED ON ONLY A SURFACE MOUNT PM AND AIR-GAP REGION IN THE MODEL B

According to the above EMCM in the model B, the equivalent parameters of PM remanence  $B_{reB}$  and thickness  $h_{pme}$  can be obtained as (24), (25), respectively. The radial and circumferential components of the air-gap flux density can be given by (45), (46), as shown at the top of the next page [30]–[32], where  $M_{rn}$ ,  $M_{\theta n}$  is the radial and tangential magnetization, respectively. when magnetization method of surface PM is parallel magnetization, then [16]

$$\begin{cases} M_{rn} = \frac{B_{reB}}{\mu_0} \alpha_p (A_{1n} + A_{2n}) \\ M_{\theta n} = \frac{B_{reB}}{\mu_0} \alpha_p (A_{1n} - A_{2n}) \end{cases} \quad (47)$$

$$B_r(r, \theta) = \sum_{n=1}^{\infty} \frac{\mu_0}{\mu_{rm}} (M_{rn} + npM_{\theta n}) \frac{np}{(np)^2 - 1} \left[ \left( \frac{R_m}{r} \right)^{np+1} + \left( \frac{r}{R_s} \right)^{np-1} \left( \frac{R_m}{R_s} \right)^{np+1} \right] \cdot \cos np(\theta - \omega_r t)$$

$$\times \left\{ \frac{\left( np - \frac{1}{np} \right) \frac{M_{rn}}{M_{rn} + npM_{\theta n}} + \frac{1}{np} - 1 + 2 \left( \frac{R_r}{R_m} \right)^{np+1} - \left[ \left( np - \frac{1}{np} \right) \frac{M_{rn}}{M_{rn} + npM_{\theta n}} + \frac{1}{np} + 1 \right] \left( \frac{R_r}{R_m} \right)^{2np}}{\frac{\mu_r + 1}{\mu_r} \left[ 1 - \left( \frac{R_r}{R_s} \right)^{2np} \right] - \frac{\mu_r - 1}{\mu_r} \left[ \left( \frac{R_m}{R_s} \right)^{2np} - \left( \frac{R_r}{R_m} \right)^{2np}]} \right\} \quad (45)$$

$$B_{\theta}(r, \theta) = \sum_{n=1}^{\infty} \frac{\mu_0}{\mu_{rm}} (M_{rn} + npM_{\theta n}) \frac{np}{(np)^2 - 1} \left[ \left( \frac{R_m}{r} \right)^{np+1} - \left( \frac{r}{R_s} \right)^{np-1} \left( \frac{R_m}{R_s} \right)^{np+1} \right] \cdot \sin np(\theta - \omega_r t)$$

$$\times \left\{ \frac{\left( np - \frac{1}{np} \right) \frac{M_{rn}}{M_{rn} + npM_{\theta n}} + \frac{1}{np} - 1 + 2 \left( \frac{R_r}{R_m} \right)^{np+1} - \left[ \left( np - \frac{1}{np} \right) \frac{M_{rn}}{M_{rn} + npM_{\theta n}} + \frac{1}{np} + 1 \right] \left( \frac{R_r}{R_m} \right)^{2np}}{\frac{\mu_r + 1}{\mu_r} \left[ 1 - \left( \frac{R_r}{R_s} \right)^{2np} \right] - \frac{\mu_r - 1}{\mu_r} \left[ \left( \frac{R_m}{R_s} \right)^{2np} - \left( \frac{R_r}{R_m} \right)^{2np}]} \right\} \quad (46)$$

where

$$A_{1n} = \frac{\sin[(np + 1)\alpha_p \frac{\pi}{2p}]}{(np + 1)\alpha_p \frac{\pi}{2p}}$$

$$A_{1n} = 1 \quad \text{for } np = 1$$

$$A_{2n} = \frac{\sin[(np - 1)\alpha_p \frac{\pi}{2p}]}{(np - 1)\alpha_p \frac{\pi}{2p}} \quad \text{for } np \neq 1 \quad (48)$$

### C. MAGNETIC FIELD DISTRIBUTION ON ARMATURE REACTION FIELD

Under load condition, the symmetric current excitations of phase *a*, *b*, and *c* as a Fourier series are given by:

$$\begin{cases} i_a = I_u \cos(p\omega_r t + \theta_u) \\ i_b = I_u \cos(p\omega_r t - 2\pi/3 + \theta_u) \\ i_c = I_u \cos(p\omega_r t + 2\pi/3 + \theta_u) \end{cases} \quad (49)$$

where  $\theta_u$  is the harmonic phase angle,  $I_u$  is the current amplitude.

Therefore, the synthetic flux density in air-gap region excited by a symmetrical 3-phase winding current is obtained by [33]:

$$B_{rabc}(\theta, r, t) = -\mu_0 \frac{2pI_u}{\pi} \frac{1}{r} \cdot \sum_{m=1}^{Q_s} (-1)^{\text{floor}(\frac{m}{2})} \cdot \sum_{n=1}^{\infty} K_{son} K_{pn} F_{nr}(r) \times \cos(n(\theta - m\alpha_y) - p\omega_r t + \theta_u) \quad (50)$$

$$B_{\theta abc}(\theta, r, t) = \mu_0 \frac{2pI_u}{\pi} \frac{1}{r} \cdot \sum_{m=1}^{Q_s} (-1)^{\text{floor}(\frac{m}{2})} \cdot \sum_{n=1}^{\infty} K_{son} K_{pn} F_{n\theta}(r) \times \sin(n(\theta - m\alpha_y) - p\omega_r t + \theta_u) \quad (51)$$

where  $K_{son}$  is the slot-opening factor;  $K_{pn}$  is winding pitch factor;  $F_{nr}$  and  $F_{n\theta}$  are two functions depend on the radius and harmonic order [24].

### D. SLOTTED AIR-GAP FLUX DENSITY

The analytical method of the notion of complex relative air-gap permeance is adopted to calculate the magnetic field distribution in air-gap region with the effect of slotting. Refer to the documents [34], a complex number  $\lambda$  can be expressed by:

$$\lambda = \lambda_a + j\lambda_b \quad (52)$$

where

$$\begin{cases} \lambda_a = \lambda_0 + \sum_{n=1}^{\infty} \lambda_{an} \cos(nQ_s\theta) \\ \lambda_b = \sum_{n=1}^{\infty} \lambda_{bn} \sin(nQ_s\theta) \end{cases} \quad (53)$$

where the Fourier coefficients  $\lambda_{an}$  and  $\lambda_{bn}$  can be calculated by using discrete Fourier transform.

Therefore, the radial and tangential components of the flux density in the slotted air-gap can be calculated by [35]:

$$\begin{cases} B_{sr} = B_r \lambda_a + B_{\theta} \lambda_b \\ B_{s\theta} = B_{\theta} \lambda_a - B_r \lambda_b \end{cases} \quad (54)$$

$$\begin{cases} B_{srabc} = B_{rabc} \lambda_a + B_{\theta abc} \lambda_b \\ B_{s\theta abc} = B_{\theta abc} \lambda_a - B_{rabc} \lambda_b \end{cases} \quad (55)$$

### E. ELECTROMAGNETIC PERFORMANCE

The flux linkage of single tooth can be calculated by:

$$\Phi_i(t) = N_c R_g L_g \int_{\alpha_i - \tau/2}^{\alpha_i + \tau/2} (B_{sr}(\theta, t) + B_{sabcr}(\theta, t)) d\theta \quad (56)$$

where  $\alpha_i$  is the angle of stator axis,  $\tau$  is coil pitch for one coil of any phase. The total phase flux linkage can be calculated by summing the phase winding flux linkage in corresponding stator. For 10-pole/12-slot, any phase flux linkage can be expressed by:

$$\Phi_k = \Phi_i - \Phi_{i+1} - (\Phi_{i+Q_s/2} - \Phi_{i+Q_s/2+1}) \quad (57)$$

where  $k$  represents phase *a*, *b* and *c*, respectively.

**TABLE 2.** Main specifications parameters of the two structure rotors.

Symbol	Quantity	Value and unit	
		Model A	Model B
$R_{r1}$	Radius of the rotor outer surface	24.2 mm	24.2 mm
$R_{r2}$	Radius of the rotor core outer surface	23.64mm	22.7 mm
$R_{r3}$	Radius of the spoke PM outer surface	23.27mm	22.35mm
$R_{r4}$	Radius of the spoke PM inner surface	9 mm	9.5 mm
$R_{r6}$	Radius of the shaft	7 mm	7 mm
$l_{mb}$	Width of magnetic bridge	1 mm	1 mm

Thus, the back-EMF can be obtained as:

$$E_a = -\frac{d\Phi_a}{dt} \tag{58}$$

$$T_e = \frac{L_{ef} r^2}{\mu_0} \int_0^{2\pi} B_{sr}(\theta, r, t) \cdot B_{s\theta abc}(\theta, r, t) + B_{s\theta}(\theta, r, t) \cdot B_{srabc}(\theta, r, t) d\theta \tag{59}$$

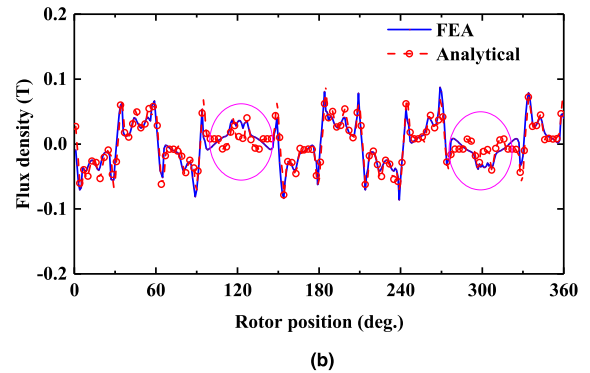
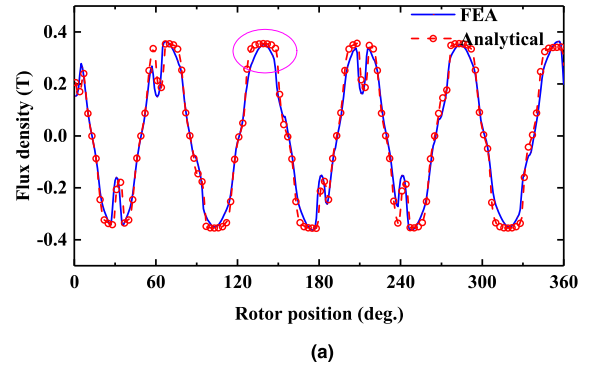
**V. ANALYTICAL PREDICTION AND VERIFICATION BY FEA**

In this section, the proposed two models will be validated by 10-pole/12-slot compared with the results of the FEA. To order to obtain an accurate electromagnetic performance, the total number of nodes in the motor with the spoke-Type rotor and hybrid rotor structure are 18429 and 23113 in 2-D FEA, respectively. The stator core B-H curve with a nonlinear change has a knee point of 1.7T. The torque will be verified when the current amplitude is 50A with the stator unsaturated. Their main parameters of the rotor are given in Table 2.

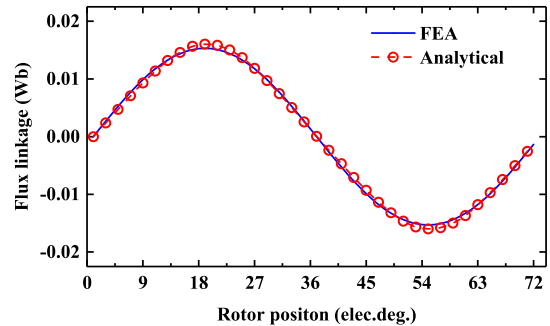
**A. THE MODEL A VALIDATION**

Based on the above parameters in Model A, the flux density components in middle air-gap region between FEA and analytical model are shown in Fig 7 (a) and (b), respectively. Fig 8 and Fig 9 are for the flux linkage and back-EMF, respectively. The electromagnetic torque is shown in Fig 10.

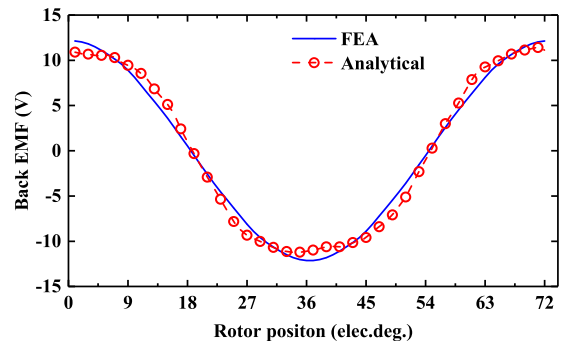
All analytical electromagnetic performances show a good agreement with FEA relatively. However, there is a difference between FEA and analytical model in circled area as shown in Fig 7. This is because that the flux density of the rotor core is not actually evenly distributed to the rotor outer surface in FEA, which is not in accordance with the assumptions in analytical predictions. Furthermore, the error of the circumferential flux density has been more impacted on the performance prediction than the radius flux density. In actual, the radius flux density is more significant influence to predict the flux linkages and back-EMF at the first stage of the design process. From the Fig 8 and Fig 9, it can be seen that the amplitude of the flux linkages and the back-EMF show good well agreement. The average torque is 5.98N.m in analytical model compared with the average torque is 6.04N.m in FEA when the peak current is 50A in the control mode of  $I_d = 0$ . The accuracy of prediction with about 1.0% is excellent although the predicted performance of torque ripples shows larger errors as shown in Table 3.



**FIGURE 7.** No-load air-gap flux density waveforms by the FEA and analytical prediction of the model A (a) radial component (b) circumferential component.



**FIGURE 8.** The flux linkage under no-load condition.



**FIGURE 9.** The back-EMF under no-load condition.

**B. THE MODEL B VALIDATION**

Based on the above parameters of the Table 2 in Model B, the air-gap flux density components between FEA and analytical

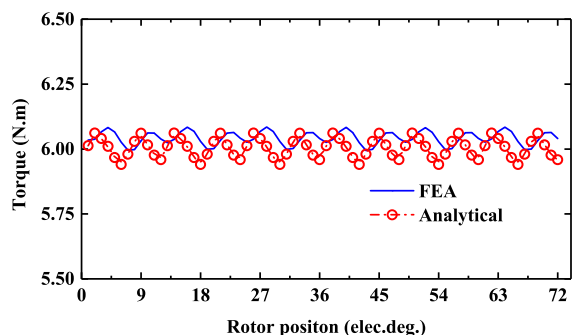


FIGURE 10. Variation of torque with electrical angle when the current  $I_u = 50A$ .

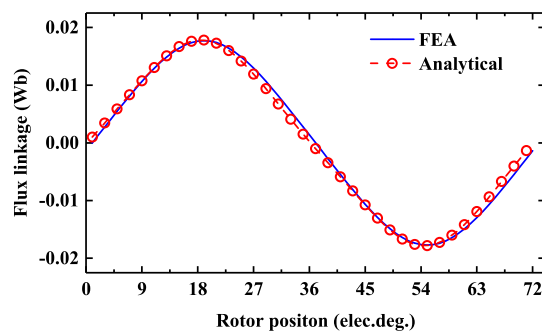
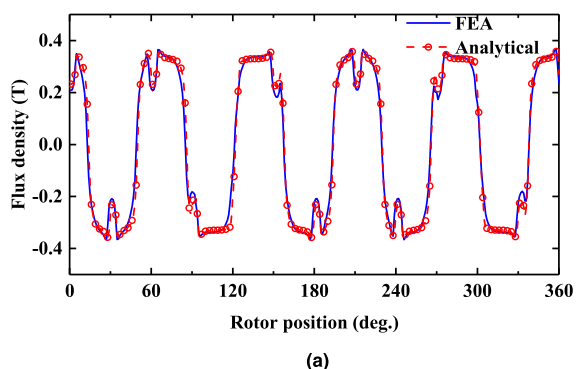
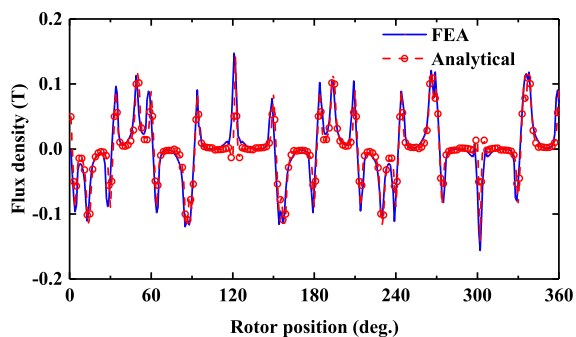


FIGURE 12. The flux linkage under no-load condition.



(a)



(b)

FIGURE 11. No-load air-gap flux density waveforms by the FEA and analytical prediction of the model B (a) radial component (b) circumferential component.

model are shown in Fig 11 (a) and (b), respectively. Fig 12 and Fig 13 are for the flux linkage and back-EMF, respectively. Fig 14 is for the electromagnetic torque. The waveforms of the electromagnetic performance show high accuracy with comparison between the FEA and analytical model.

As shown in Fig 11, the comparisons in the model B show a smaller error of the air-gap flux density between the FEA and analytical model than the errors in the model A. It is well worth to be noted that: the air-gap region is in direct contact with the surface-mount PMs in hybrid structure rotor, then the air-gap flux density can be little impacted by the rotor iron outer-surface with uneven flux density distribution as the model A. Therefore, not only the torque amplitude can be predicted very well with the error of 0.5%, excellent

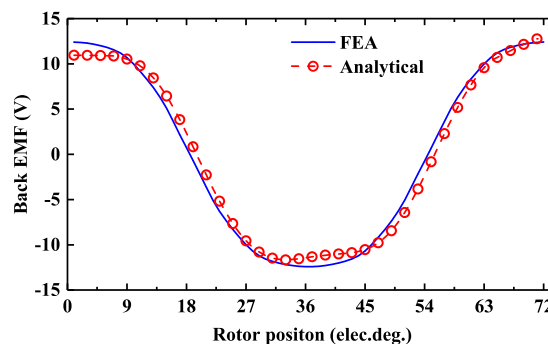


FIGURE 13. The back-EMF under no-load condition.

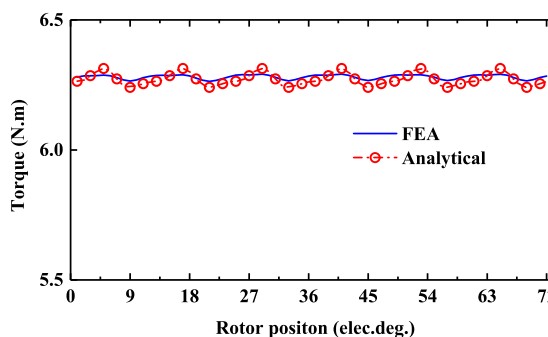


FIGURE 14. Variation of torque with electrical angle when the current  $I_u = 50A$ .

agreement between the predictions of torque ripples can be achieved as shown in Fig 14 with the same control current and control mode.

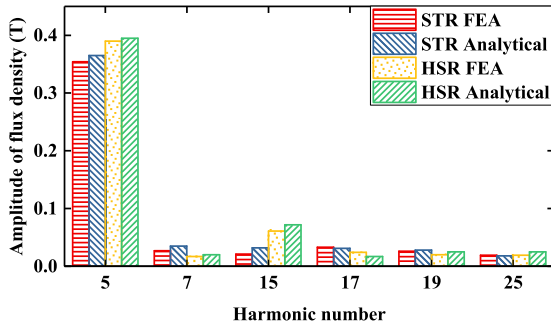
### C. COMPARISONS OF THE MAGNETIC PERFORMANCE BETWEEN TWO MODELS

The comparisons of the average torque and torque ripples are shown in Table 3. It is obvious that the torque of STR is relatively larger than that of HRS in the FEA and analytical model. The reason is that the thickness of the PMs at the main flux path is increased, even though there is a little reduction in width of the spoke-type PMs because of the restriction of the rotor size. The fundamental amplitude i.e. 5<sup>th</sup> harmonic of the air-gap flux density is higher in the HRS than in the STR as shown in Fig 15. Therefore, the torque density can be



**TABLE 3.** The comparisons of the average torque and torque ripples.

method	torque (STR)	torque (HRS)	Torque ripple (STR)	Torque ripple (HRS)
FEA	6.04	6.28	1.3%	0.9%
Analytical	5.98	6.27	3.6%	1.1%

**FIGURE 15.** Harmonic component of STR and HRS model predicted by the flux density between the FEA and analytical model.

increased by adopting the proposed model of the hybrid rotor structure (HRS).

Moreover, the torque ripples have smaller values whether it is in the FEA or analytical model, which can exclude the influence of the calculation error with uneven flux density distribution in the rotor iron outer-surface of the model A. According to the aforementioned, the surface mount PMs in HRS, which magnetic permeability is close to 1, is equivalent to increase the effective length of the air-gap. Then, the effects of the slotted opening are weakened. The ratio of the inter-harmonics contents decreases such as 7<sup>th</sup>, 17<sup>th</sup>, and 19<sup>th</sup>, as shown in Fig 15. This is not to be considered in ripples that 15<sup>th</sup> harmonic which is three times the fundamental does not react with the armature.

The FEA and analytical model has its own advantage and application. The FEA is suitable for accurate calculation of the flux density, back-EMF and torque, even considering the effects of saturation and demagnetization. However, this is the exchanged result at the expense of time compared with the analytical model. The computation times of the two methods in the model A is 143.5s and 27.7s, respectively. Correspondingly, the times in the model B is 177.6s and 32.6s, respectively. It is demonstrated that the analytical model can save nearly 80.7% and 81.7% of the simulation time compared with the FEA, respectively. It is very significant to optimize solutions at the initial design process.

## VI. CONCLUSION

In this paper, two analytical methods have been proposed to predict the magnetic field distribution and the torque performance in the machines with the irregular rotor structure. The equivalent magnetic network model has been established considering leakage flux linkage and irregular PM shape. In order to obtain boundary conditions in the model with the STR and HRS, two equivalent methods have been adopted:

one is that the average flux density in the rotor core outer surface is obtained as the boundary conditions of Laplace's equation in the air-gap regions of model A; The other is that the equivalent parameters of PM remanence  $B_{reB}$  and thickness  $h_{pme}$  can be obtained by the Thevenin's theorem in the model B.

Based on the boundary conditions, the magnetic field analysis result has been obtained with considering for the rotor core outer-surface eccentric in the model A. Meanwhile, by utilizing the equivalent parameters of  $B_{reB}$  and  $h_{pme}$ , the air-gap magnetic density has been solved for conventional classic magnetic field analysis model of surface-mounted PMs with air-gap in model B. To validate the good accuracy for the air-gap density component, flux linkage, back-EMF, and torque, the respective comparison results between the FEA and the analytical model have been shown, which revealed that results can match well with each other. Then the analytical model can save nearly 80% computation time. In addition, the machine with the proposed hybrid PM rotor structure can improve the torque density compared with that of the spoke-type PM structure rotor.

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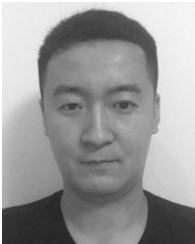
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