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# Average Symbol Error Rate Analysis of QAM Schemes Over Millimeter Wave Fluctuating Two-Ray Fading Channels

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**ABSTRACT** In this study, analytical expressions are derived by using probability density function approach over millimeter wave fluctuating two-ray (FTR) fading channels for average symbol error rate (ASER) of rectangular quadrature amplitude modulation (RQAM) and cross quadrature amplitude modulation (XQAM) schemes. First, an exact ASER expression for RQAM over millimeter wave FTR fading channels is derived. Second, two different upper bound ASER expressions of RQAM scheme under millimeter wave FTR fading conditions are obtained by using Chernoff and Chiani approximations of Gaussian *Q*-function. In addition to this, an asymptotic ASER expression for RQAM is also proposed. Then, an analytical formulation is presented for XQAM signaling over millimeter wave FTR fading channels in terms of infinite series representations which rapidly converge. Finally, the numerical results of the proposed analytical expressions are compared with the simulation results in order to validate the analytical findings in this work.

**INDEX TERMS** Rectangular QAM, Cross QAM, ASER, fluctuating two-ray fading.

#### I. INTRODUCTION

Quadrature amplitude modulation (QAM) is a bandwidth efficient modulation type and it is used in the field of digital multimedia transmission. The advantages of high power and bandwidth efficiency makes the use of QAM signaling in LTE Advanced and beyond preferable [1]. QAM, is one of many existing modulation techniques, is a promising way of enhancing the spectral efficiency [2]. Besides, QAM schemes are known as useful adaptive modulation schemes, since the constellation size of the QAM signaling can be adjusted depending on the channel quality. QAM methods have drawn considerable research attention due to the fact that it is broadly used in many wireless communication applications such as microwave, high-speed mobile communication systems, and asymmetric subscriber loop, etc [3]–[5]. There are several QAM schemes which are called as rectangular QAM (RQAM), cross QAM (XQAM) and square QAM (SQAM). RQAM is a general modulation technique since it covers a variety of well-known modulations such as SQAM, orthogonal binary frequency-shift keying, quadrature

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phase-shift keying, binary phase-shift keying [5]. Practically, RQAM has the implementation fields in microwave communications, high rate wireless communications, and telephoneline modems [6], [7]. Among the QAM schemes mentioned above, the optimal QAM signaling is XQAM scheme when the odd number of bits per symbol is transmitted due to its low average symbol energy compared to RQAM [8]. For instance, constellation size between 5-15 bits for XQAMs can be used in asymmetric and high speed subscriber lines [9], [10], while 32-XQAM and 128-XQAM are utilized in digital video broadcasting [11].

Recently, QAM methods have become popular due to its suitability in high speed communications and bandwidth efficiency. The researchers have reported a number of studies on average symbol error rate (ASER) performance of QAM schemes used in different communication systems over different fading channels [12]–[22]. In [12], the ASER of RQAM and XQAM signaling was presented in two-wave with diffuse power (TWDP) fading channels. Closed-form error probability expressions for *M*-ary XQAM signaling were proposed by using a maximal-ratio combining (MRC) technique in  $\eta$ - $\mu$  fading environments in [13]. The authors of [14] analyzed the ASER performance of QAM techniques

in two-way relaying communication networks in the presence of Nakagami-m fading. Dixit and Sahu [15] derived an expression for RQAM technique with selection combining over Nakagami-m fading conditions. The effects of imperfect channel state information on the ASER of hexagonal QAM and RQAM schemes were studied for an amplifyand-forward (AF) protocol under Nakagami-m fading conditions [16]. The impact of the channel estimation [23], [24] is important for the realization of the practical scenarios. In another study [17], the authors performed an analytical evaluation on the performance of QAM methods used in multiple AF relaying system over Rayleigh fading channels. In [18], a simple approximation was developed for generalorder RQAM with MRC receiver in Nakagami-m environments. In addition, hexagonal QAM and ROAM schemes were used in an orthogonal frequency division multiplexing with 3-hop AF relaying system and the authors presented the performance results over mixed Rician/Rayleigh fading links [19]. In [20], error probability of XQAM method was derived over Nakagami-m, Hoyt, Rice, and Rayleigh conditions by performing MRC at the receiver. Khan et al. [21] proposed two algorithms called rectangular contour algorithm (RCA) and improved RCA for blind equalization techniques of RQAM signaling with odd bit. In [22], a comprehensive study for QAM analysis over  $\kappa$ - $\mu$  shadowed fading channels is presented.

The studies mentioned above have focused on the performance analysis of QAM signaling over well-known fading conditions such as Rice, Rayleigh, Hoyt, Nakagami-m, and shadowed etc. As far as we know, there is no work that analyses and presents the ASER performance of RQAM and XQAM techniques over millimeter wave (mmWave) fluctuating two-ray (FTR) fading channels. FTR fading distribution has been recently presented as a channel model which well captures the effects of wireless medium in mmWave and device-to-device environments [25], [26]. Very recently, an outdoor measurement at 28 GHz explored that the well-known small-scale fading distributions such as Rician, Nakagami-*m* and Rayleigh cannot accurately reflect the random fluctuations in millimeter wave signals [27]. In order to overcome this, the FTR fading is proposed which is more accurate compared to Rayleigh, Nakagami-m, Rician models and to the best of our knowledge, there are only a few works over this new fading channel which can be found in [28]–[30]. Although the FTR fading distribution contains classical unimodal fading models such as Rician or Rayleigh special cases, this model is inherently bimodal. At the same time, the mmWave FTR model is a great and appropriate fading model which takes the large heterogeneity of random fluctuations into consideration that effect the mmWave radio signal when propagating in the presence of multiple scatters. Also, this model is a native generalization of the TWDP fading model that allows the fixed amplitude specular waves associated with line-of-sight propagation to randomly fluctuate. In addition to these advantages, it also contains some conventional fading distributions as special cases, i.e., Rayleigh,

Rician, Rician shadowed, one-sided Gaussian, Nakagami-*q* (Hoyt), Nakagami-*m*, and TWDP. The mmWave FTR fading model turns into these particular cases by using the parameter definitions [25], [26].

Motivated by the suitability of FTR fading for millimeter wave communications and bandwidth efficient high speed transmission capability of QAM signaling, we derive new and novel analytical expressions for RQAM and XQAM schemes over millimeter wave FTR fading channels for the first time in the literature in this work. First, we derive the exact expression for ASER of RQAM system over millimeter wave FTR fading channels in terms of bivariate Meijer's G function and then, we propose two different upper bounds for the ASER of RQAM based on well-known approximations of Gaussian O-function. Second, we also obtain an asymptotic ASER expression of RQAM system to illustrate the behaviors of the considered system at the high signal to noise ratio (SNR) regime. Then, we evaluate the performance of XQAM scheme over FTR fading channels. Finally, the obtained analytical results are compared to the corresponding simulations for validation purpose.

#### **II. SYSTEM AND CHANNEL MODELS**

Consider a single-input single-output (SISO) wireless network over FTR fading channels and let  $\gamma$  denote the instantaneous SNR.  $\overline{\gamma} = (E_b/N_0) 2\sigma^2 (1+K)$  is the average SNR [28].  $E_b$  is the energy per bit,  $N_0$  is the additive white Gaussian noise (AWGN) and  $\sigma^2$  is the variance of Gaussian noise. The probability density function (PDF) of instantaneous SNR,  $f_{\gamma}(\gamma)$ , can be given as [28]

$$f_{\gamma}(\gamma) = v \sum_{j=0}^{\infty} A_j \gamma^j \exp\left(-\frac{\gamma}{2\sigma^2}\right)$$
(1)

where  $v = \frac{m^m}{\Gamma(m)}$ ,  $A_j = \frac{K^j d_j}{j! \Gamma(j+1) (2\sigma^2)^{j+1}}$  and

$$d_{j} = \sum_{k=0}^{\infty} {\binom{j}{k}} {\binom{\Delta}{2}}^{k} \sum_{l=0}^{k} {\binom{k}{l}} \Gamma \left(j+m+2l-k\right)$$
$$\times \exp\left(\frac{\pi \left(2l-k\right)i}{2}\right) \times \left((m+K)^{2}-(K\Delta)^{2}\right)^{-\frac{(j+m)}{2}}$$
$$\times P_{j+m-1}^{k-2l} \left(\frac{m+K}{\sqrt{(m+K)^{2}-(K\Delta)^{2}}}\right)$$
(2)

where  $\Gamma(\cdot)$  is the Gamma function, *m* is a Nakagami-*m* random variable, *K* is the ratio of the average power of the dominant waves to the average power of the remaining diffuse multipath and it is defined by  $K = \frac{V_1^2 + V_2^2}{2\sigma^2}$ .  $\Delta$  denotes the relation between the powers of two dominant waves and given by  $\Delta = \frac{2V_1V_2}{V_1^2 + V_2^2}$ .  $V_1$  and  $V_2$  are two specular components,  $P_b^a(\cdot)$  is the Legendre functions of the first kind [31, eq. (8.702)]. *j* and *l* are the indexes of the summations in eqs (1) and (2). The FTR fading model includes several well-known fading models for different values of *m*, *K* and  $\Delta$  such as

Rayleigh, Nakagami-q, Rician, Rician shadowed and the onesided Gaussian as mentioned before.

## **III. PERFORMANCE ANALYSIS**

In this section, we provide an exact and two different approximate derivations of ASER for RQAM and an ASER expression for XQAM of the considered system under FTR fading conditions.

### A. EXACT ASER ANALYSIS FOR RQAM

The ASER, Pe, is generally defined as [5]

$$Pe = \int_{0}^{\infty} p(e|\gamma) f_{\gamma}(\gamma) d\gamma$$
(3)

where  $p(e|\gamma)$  is the conditional error probability for AWGN and it is given for *M*-ary RQAM as [32]

$$p(e|\gamma) = 2pQ(a\sqrt{\gamma}) + 2qQ(b\sqrt{\gamma}) - 4pqQ(a\sqrt{\gamma})Q(b\sqrt{\gamma})$$
(4)

where  $Q(\cdot)$  is the Gaussian Q-function.  $M = M_I \times M_Q$ ,  $p = 1 - (1/M_I), q = 1 - (1/M_Q), b = \beta a, \beta = d_Q/d_I$ ,  $a = \sqrt{6/(M_I^2 - 1) + (M_Q^2 - 1)\beta^2}, d_Q$  and  $d_I$  are the quadrature and in-phase decision distance, respectively.

Substituting (1) and (4) into (3), we have

$$Pe = \int_{0}^{\infty} \{2pQ(a\sqrt{\gamma}) + 2qQ(b\sqrt{\gamma}) - 4pqQ(a\sqrt{\gamma})Q(b\sqrt{\gamma})\} \times v \sum_{j=0}^{\infty} A_{j}\gamma^{j} \exp\left(-\frac{\gamma}{2\sigma^{2}}\right) d\gamma \quad (5)$$

After some manipulations and using

 $Q(x) = (1/2) \operatorname{erfc} (x/\sqrt{2})$ , the expression in (5) can be rewritten as follows

$$Pe = \int_{0}^{\infty} \left\{ p \operatorname{erfc}\left(a\sqrt{\frac{\gamma}{2}}\right) + q \operatorname{erfc}\left(b\sqrt{\frac{\gamma}{2}}\right) - pq \operatorname{erfc}\left(a\sqrt{\frac{\gamma}{2}}\right) \right\}$$
$$\times \operatorname{erfc}\left(b\sqrt{\frac{\gamma}{2}}\right) v \sum_{j=0}^{\infty} A_{j}\gamma^{j} \exp\left(-\frac{\gamma}{2\sigma^{2}}\right) d\gamma \quad (6)$$

where erfc  $(\cdot)$  is the complementary error function. By using [33, eq. (8.4.14.2)], 6 becomes as shown in (7) at the top of next page.

where  $G(\cdot | \cdot)$  is the Meijer's *G* function [31]. Then,  $F_1$  and  $F_2$  are derived with the aid of [33, eq. (2.24.3.1)] as

$$F_{1} = \left(\frac{1}{2\sigma^{2}}\right)^{-(j+1)} G_{2,2}^{2,1} \left(a^{2}\sigma^{2} \middle| \begin{array}{c} \Delta(1,-j), \Delta(1,1) \\ \Delta(1,0.5) \end{array}\right)$$
$$F_{2} = \left(\frac{1}{2\sigma^{2}}\right)^{-(j+1)} G_{2,2}^{2,1} \left(b^{2}\sigma^{2} \middle| \begin{array}{c} \Delta(1,-j), \Delta(1,1) \\ \Delta(1,0.5) \end{array}\right)$$
(8)

where  $\Delta(u, t) = \frac{t}{u}, \frac{t+1}{u}, \dots, \frac{t+u-1}{u}$ . In order to obtain a solution for  $F_3$ , we use the identity between exponential function and Meijer's *G* function in [33, eq. (8.4.3.1)] and we get

$$F_{3} = \frac{pq}{\pi} \int_{0}^{\infty} G_{1,2}^{2,0} \left( \frac{a^{2}\gamma}{2} \middle| \begin{array}{c} 1\\ 0, 0.5 \end{array} \right) G_{1,2}^{2,0} \left( \frac{b^{2}\gamma}{2} \middle| \begin{array}{c} 1\\ 0, 0.5 \end{array} \right) \\ \times G_{0,1}^{1,0} \left( \frac{\gamma}{2\sigma^{2}} \middle| \begin{array}{c} -\\ 0 \end{array} \right) \gamma^{j} d\gamma$$
(9)

With the help of [34, eq. (07.34.21.0081.01)],  $F_3$  is derived as

$$F_{3} = \left(\frac{a^{2}}{2}\right)^{-(j+1)} \times G_{2,1:1,2:0,1}^{0,2:2,0:1,0} \left(\begin{array}{c}-j, -j - 0.5 \\ -j - 1\end{array}\right) \left|\begin{array}{c}1\\0, 0.5\end{array}\right| \left|\begin{array}{c}-b^{2}\\0\end{array}\right| \left|\begin{array}{c}a^{2}}{a^{2}}, \frac{1}{\sigma^{2}a^{2}}\right)$$
(10)

where  $G_{p,q;p_1,q_1:p_2,q_2}^{m,n:m_1,n_1:m_2,n_2}$  [ $\cdot$  | $\cdot$  | $\cdot$  | $\cdot$ ,  $\cdot$ ] is the bivariate Meijer's *G* function. The exact ASER expression for RQAM is obtained by inserting (8) and (10) into (7), as shown at the top of the next page, and it is given below

$$Pe = v \sum_{j=0}^{\infty} A_j \left[ \left( \frac{1}{2\sigma^2} \right)^{-(j+1)} \times G_{2,2}^{2,1} \left( a^2 \sigma^2 \middle| \begin{array}{c} \Delta (1, -j) , \Delta (1, 1) \\ \Delta (1, 0.5) \end{array} \right) + \left( \frac{1}{2\sigma^2} \right)^{-(j+1)} \times G_{2,2}^{2,1} \left( b^2 \sigma^2 \middle| \begin{array}{c} \Delta (1, -j) , \Delta (1, 1) \\ \Delta (1, 0.5) \end{array} \right) - \left( \frac{a^2}{2} \right)^{-(j+1)} \times G_{2,1:1,2:0,1}^{0,2:2,0:1,0} \left( \begin{array}{c} -j, -j - 0.5 \\ -j - 1 \end{array} \right) \left| \begin{array}{c} 0 \\ 0 \\ 0 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ 0 \\ \end{array} \right| \left| \begin{array}{c} b^2 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\ a^2 \\ a^2 \\ a^2 \\ \end{array} \right| \left| \begin{array}{c} 0 \\ a^2 \\$$

Even though (11) is an exact ASER expression for RQAM schemes over mmWave FTR fading channels, this expression is very difficult to compute in common software packages since the bivariate Meijer's *G* function is not built-in function in Matlab, Mathematica, and Maple. In what follows, two different approximations for the ASER of RQAM in simple and easy-to-compute form will be derived.

### B. ASER ANALYSIS WITH CHERNOFF APPROXIMATION FOR RQAM

To facilitate the analysis, we require to use an approximation for  $Q(\cdot)$  functions in (4). The Chernoff approximation is a well-known approximation of  $Q(\cdot)$  and it is given by  $Q(z) \approx \frac{1}{2}e^{-z^2/2}$  [35]. By using the Chernoff approximation,  $p(e|\gamma)$ in (4) is rewritten as

$$p(e|\gamma) \approx p\left[e^{-\frac{a^2}{2}\gamma}\right] + q\left[e^{-\frac{b^2}{2}\gamma}\right] - pq\left[e^{-\frac{a^2+b^2}{2}\gamma}\right] \quad (12)$$

-

$$Pe = v \sum_{j=0}^{\infty} A_{j} \left[ \underbrace{\frac{p}{\sqrt{\pi}} \int_{0}^{\infty} G_{1,2}^{2,0} \left( \frac{a^{2}\gamma}{2} \Big|_{0,0.5}^{1} \right) \gamma^{j} \exp\left(-\frac{\gamma}{2\sigma^{2}}\right) d\gamma}_{F_{1}} + \underbrace{\frac{q}{\sqrt{\pi}} \int_{0}^{\infty} G_{1,2}^{2,0} \left( \frac{b^{2}\gamma}{2} \Big|_{0,0.5}^{1} \right) \gamma^{j} \exp\left(-\frac{\gamma}{2\sigma^{2}}\right) d\gamma}_{F_{2}} - \underbrace{\frac{pq}{\pi} \int_{0}^{\infty} G_{1,2}^{2,0} \left( \frac{a^{2}\gamma}{2} \Big|_{0,0.5}^{1} \right) G_{1,2}^{2,0} \left( \frac{b^{2}\gamma}{2} \Big|_{0,0.5}^{1} \right) \gamma^{j} \exp\left(-\frac{\gamma}{2\sigma^{2}}\right) d\gamma}_{F_{3}} \right]$$
(7)

Substituting (1) and (12) into (3), the ASER formulation with Chernoff approximation for RQAM becomes

$$Pe_{Chernoff} = \int_{0}^{\infty} \left\{ p \left[ e^{-\frac{a^2}{2}\gamma} \right] + q \left[ e^{-\frac{b^2}{2}\gamma} \right] - pq \left[ e^{-\frac{a^2+b^2}{2}\gamma} \right] \right\}$$
$$\times v \sum_{j=0}^{\infty} A_j \gamma^j \exp\left(-\frac{\gamma}{2\sigma^2}\right) d\gamma \quad (13)$$

Since the  $Pe_{Chernoff}$  expression in (13) has three similar integrals, after some mathematical manipulations and using [31, eq. (3.381.4)], we derive the ASER as follows

$$Pe_{Chernoff} = pv \sum_{j=0}^{\infty} A_j \left(\frac{a^2}{2} + \frac{1}{2\sigma^2}\right)^{-(j+1)} \Gamma(j+1) + qv \sum_{j=0}^{\infty} A_j \left(\frac{b^2}{2} + \frac{1}{2\sigma^2}\right)^{-(j+1)} \Gamma(j+1) - pqv \sum_{j=0}^{\infty} A_j \left(\frac{a^2}{2} + \frac{b^2}{2} + \frac{1}{2\sigma^2}\right)^{-(j+1)} \Gamma(j+1)$$
(14)

# C. ASER ANALYSIS WITH CHIANI APPROXIMATION FOR RQAM

To analyze the ASER for RQAM with a different way, we use the Chiani approximation of Gaussian *Q*-function which is expressed by  $Q(z) \approx \frac{1}{12}e^{-z^2/2} + \frac{1}{4}e^{-2z^2/3}$  [36]. By using the Chiani approximation,  $p(e|\gamma)$  in (4) is rewritten as

$$p(e|\gamma) \approx 2p \left[ \frac{1}{12} e^{-\frac{a^2}{2}\gamma} + \frac{1}{4} e^{-\frac{2a^2}{3}\gamma} \right] + 2q \left[ \frac{1}{12} e^{-\frac{b^2}{2}\gamma} + \frac{1}{4} e^{-\frac{2b^2}{3}\gamma} \right] - 4pq \left[ \frac{1}{144} e^{-\left(\frac{a^2+b^2}{2}\right)\gamma} + \frac{1}{48} e^{-\left(\frac{a^2}{2} + \frac{2b^2}{3}\right)\gamma} + \frac{1}{48} e^{-\left(\frac{2a^2}{3} + \frac{b^2}{2}\right)\gamma} + \frac{1}{16} e^{-\left(\frac{2a^2+2b^2}{3}\right)\gamma} \right]$$
(15)

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By inserting (15) and (1) into (3), we have

$$Pe_{Chiani} = v \sum_{j=0}^{\infty} A_j \left[ Q_1 + Q_2 - Q_3 \right]$$
(16)

where

$$Q_{1} = 2p \int_{0}^{\infty} \left[ \frac{1}{12} e^{-\frac{a^{2}}{2}\gamma} + \frac{1}{4} e^{-\frac{2a^{2}}{3}\gamma} \right] \gamma^{j} \exp\left(-\frac{\gamma}{2\sigma^{2}}\right) d\gamma$$

$$Q_{2} = 2q \int_{0}^{\infty} \left[ \frac{1}{12} e^{-\frac{b^{2}}{2}\gamma} + \frac{1}{4} e^{-\frac{2b^{2}}{3}\gamma} \right] \gamma^{j} \exp\left(-\frac{\gamma}{2\sigma^{2}}\right) d\gamma$$

$$(18)$$

$$Q_{3} = 4nq \int_{0}^{\infty} \left[ \frac{1}{-\frac{1}{2}} e^{-\left(\frac{a^{2}+b^{2}}{2}\right)\gamma} + \frac{1}{-\frac{1}{2}} e^{-\left(\frac{a^{2}}{2}+\frac{2b^{2}}{3}\right)\gamma} \right]$$

$$2_{3} = 4pq \int_{0} \left[ \frac{1}{144} e^{-\left(\frac{-2}{2}\right)\gamma} + \frac{1}{48} e^{-\left(\frac{2}{2} + \frac{-3}{3}\right)\gamma} + \frac{1}{48} e^{-\left(\frac{2a^{2}+2b^{2}}{3}\right)\gamma} \right]$$
$$+ \frac{1}{48} e^{-\left(\frac{2a^{2}}{3} + \frac{b^{2}}{2}\right)\gamma} + \frac{1}{16} e^{-\left(\frac{2a^{2}+2b^{2}}{3}\right)\gamma} \right]$$
$$\times \gamma^{j} \exp\left(-\frac{\gamma}{2\sigma^{2}}\right) d\gamma$$
(19)

Now, we have eight similar integrals in  $Q_1$ ,  $Q_2$  and  $Q_3$  to be solved for deriving an ASER expression with Chiani approximation for RQAM over FTR fading channels. To proceed, with the use of [31, eq. (3.381.4)] and after performing some analytical manipulations, we get

$$Q_{1} = 2p \left( \frac{1}{12} \left( \frac{a^{2}}{2} + \frac{1}{2\sigma^{2}} \right)^{-(j+1)} \Gamma(j+1) + \frac{1}{4} \left( \frac{2a^{2}}{3} + \frac{1}{2\sigma^{2}} \right)^{-(j+1)} \Gamma(j+1) \right)$$
(20)  
$$Q_{2} = 2q \left( \frac{1}{12} \left( \frac{b^{2}}{2} + \frac{1}{2\sigma^{2}} \right)^{-(j+1)} \Gamma(j+1) \right)$$

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$$+\frac{1}{4}\left(\frac{2b^2}{3} + \frac{1}{2\sigma^2}\right)^{-(j+1)}\Gamma(j+1)\right)$$
(21)

$$Q_{3} = 4pq \left( \frac{1}{144} \left( \frac{a^{2} + b^{2}}{2} + \frac{1}{2\sigma^{2}} \right)^{-(j+1)} \Gamma(j+1) + \frac{1}{48} \left( \frac{a^{2}}{2} + \frac{2b^{2}}{3} + \frac{1}{2\sigma^{2}} \right)^{-(j+1)} \Gamma(j+1) + \frac{1}{48} \left( \frac{2a^{2}}{3} + \frac{b^{2}}{2} + \frac{1}{2\sigma^{2}} \right)^{-(j+1)} \Gamma(j+1) + \frac{1}{16} \left( \frac{2a^{2} + 2b^{2}}{3} + \frac{1}{2\sigma^{2}} \right)^{-(j+1)} \Gamma(j+1) \right)$$
(22)

After putting  $Q_1$ ,  $Q_2$  and  $Q_3$  in (16), an ASER expression with Chiani approximation for RQAM of the considered system in the presence of FTR fading is derived as follows

$$Pe_{Chiani} = v \sum_{j=0}^{\infty} \frac{K^{j} d_{j}}{j! (2\sigma^{2})^{j+1}} \\ \times \left[ \left\{ 2p \left( \frac{1}{12} \left( \frac{a^{2}}{2} + \frac{1}{2\sigma^{2}} \right)^{-(j+1)} \right) + \frac{1}{4} \left( \frac{2a^{2}}{3} + \frac{1}{2\sigma^{2}} \right)^{-(j+1)} \right) \right\} \\ + \left\{ 2q \left( \frac{1}{12} \left( \frac{b^{2}}{2} + \frac{1}{2\sigma^{2}} \right)^{-(j+1)} \right) + \frac{1}{4} \left( \frac{2b^{2}}{3} + \frac{1}{2\sigma^{2}} \right)^{-(j+1)} \right) \right\} \\ - \left\{ 4pq \left( \frac{1}{144} \left( \frac{a^{2} + b^{2}}{2} + \frac{1}{2\sigma^{2}} \right)^{-(j+1)} + \frac{1}{48} \left( \frac{a^{2}}{2} + \frac{2b^{2}}{3} + \frac{1}{2\sigma^{2}} \right)^{-(j+1)} + \frac{1}{48} \left( \frac{2a^{2}}{3} + \frac{b^{2}}{2} + \frac{1}{2\sigma^{2}} \right)^{-(j+1)} + \frac{1}{16} \left( \frac{2a^{2} + 2b^{2}}{3} + \frac{1}{2\sigma^{2}} \right)^{-(j+1)} \right) \right\} \right]$$
(23)

So far as we know that the expressions,  $Pe_{Chiani}$  in (23) and  $Pe_{Chernoff}$  in (14), are presented for the first time in the literature. Now, we investigate the asymptotic behavior of the ASER at high SNR to specify the diversity order acquired by the considered system when Chiani approximation is used. With the assumption of  $\overline{\gamma} \rightarrow \infty$ , the  $2\sigma^2$ approaches to  $\infty$  and truncating the infinite summation at j = 0, the expression in (23) is simplified and obtained in

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asymptotic form as

 $Pe_{Chiani}^{\infty}$ 

$$= v \frac{d_0}{2\sigma^2} \times \left[ \left\{ 2p \left( \frac{1}{12} \left( \frac{a^2}{2} \right)^{-1} + \frac{1}{4} \left( \frac{2a^2}{3} \right)^{-1} \right) \right\} \\ + \left\{ 2q \left( \frac{1}{12} \left( \frac{b^2}{2} \right)^{-1} + \frac{1}{4} \left( \frac{2b^2}{3} \right)^{-1} \right) \right\} \\ - \left\{ 4pq \left( \frac{1}{144} \left( \frac{a^2 + b^2}{2} \right)^{-1} + \frac{1}{48} \left( \frac{a^2}{2} + \frac{2b^2}{3} \right)^{-1} \\ + \frac{1}{48} \left( \frac{2a^2}{3} + \frac{b^2}{2} \right)^{-1} + \frac{1}{16} \left( \frac{2a^2 + 2b^2}{3} \right)^{-1} \right) \right\} \right]$$
(24)

**D. EXACT ASER ANALYSIS FOR XQAM** For XQAM,  $p(e|\gamma)$  is given as [13]

$$p(e|\gamma)$$

$$= g_1 Q_z \left( a_0 \sqrt{\gamma}, \pi/2 \right) + \frac{4}{M} Q_z \left( a_1 \sqrt{\gamma}, \pi/2 \right)$$

$$- g_2 Q_z \left( a_0 \sqrt{\gamma}, \pi/4 \right) - \frac{8}{M} \sum_{k=1}^{w-1} Q_z \left( a_0 \sqrt{\gamma}, \alpha_k \right)$$

$$- \frac{4}{M} \sum_{k=1}^{w-1} Q_z \left( a_k \sqrt{\gamma}, \beta_k^+ \right) + \frac{4}{M} \sum_{k=2}^{w} Q_z \left( a_k \sqrt{\gamma}, \beta_k^- \right)$$
(25)

where  $M = 2^5, 2^7, ..., w = \frac{\sqrt{2M}}{8}, a_0 = \sqrt{\frac{96}{(31M-32)}}, a_k = \sqrt{2} ka_0, k = 1, 2, 3, ..., w, g_2 = 4 - \frac{12}{\sqrt{2M}} + \frac{12}{2M}, g_1 = 4 - \frac{6}{\sqrt{2M}}, \alpha_k = \arctan\left(\frac{1}{2k+1}\right), k = 1, 2, ..., (w-1), \beta_k^- = \arctan\left(\frac{k}{k-1}\right), k = 2, 3..., w, \beta_k^+ = \arctan\left(\frac{k}{k+1}\right) \text{ and, } k = 1, 2, ..., (w-1). Q_z(\cdot, \cdot) \text{ is the generalized Marcum } Q$ -function. Substituting (25) and (1) into (3), we have

Pe<sub>XQAM</sub>

$$= v \sum_{j=0}^{\infty} A_{j} \left[ \underbrace{\int_{0}^{\infty} g_{1}Q_{z} \left(a_{0}\sqrt{\gamma}, \pi/2\right) \gamma^{j} \exp\left(-\frac{\gamma}{2\sigma^{2}}\right) d\gamma}_{D_{1}} + \underbrace{\int_{0}^{\infty} \frac{4}{M}Q_{z} \left(a_{1}\sqrt{\gamma}, \pi/2\right) \gamma^{j} \exp\left(-\frac{\gamma}{2\sigma^{2}}\right) d\gamma}_{D_{2}} - \underbrace{\int_{0}^{\infty} g_{2}Q_{z} \left(a_{0}\sqrt{\gamma}, \pi/4\right) \gamma^{j} \exp\left(-\frac{\gamma}{2\sigma^{2}}\right) d\gamma}_{D_{3}} \right]$$

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$$- \underbrace{\int_{0}^{\infty} \frac{8}{M} \sum_{k=1}^{w-1} Q_{z} \left( a_{0} \sqrt{\gamma}, \alpha_{k} \right) \gamma^{j} \exp \left( -\frac{\gamma}{2\sigma^{2}} \right) d\gamma}_{D_{4}}}_{D_{4}}$$

$$- \underbrace{\int_{0}^{\infty} \frac{4}{M} \sum_{k=1}^{w-1} Q_{z} \left( a_{k} \sqrt{\gamma}, \beta_{k}^{+} \right) \gamma^{j} \exp \left( -\frac{\gamma}{2\sigma^{2}} \right) d\gamma}_{D_{5}}}_{D_{5}}$$

$$+ \underbrace{\int_{0}^{\infty} \frac{4}{M} \sum_{k=2}^{w} Q_{z} \left( a_{k} \sqrt{\gamma}, \beta_{k}^{-} \right) \gamma^{j} \exp \left( -\frac{\gamma}{2\sigma^{2}} \right) d\gamma}_{D_{6}}}_{D_{6}}$$
(26)

The expression in (26) contains six similar integrals such as  $D_1$ - $D_6$ . In order to solve  $D_1$ ,  $D_1$  is re-arranged with the use of the infinite series representation of generalized Marcum Q-function [37, eq. (4.37)] as follows

$$D_{1} = \int_{0}^{\infty} g_{1}Q_{z} \left(a_{0}\sqrt{\gamma}, \pi/2\right) \gamma^{j} \exp\left(-\frac{\gamma}{2\sigma^{2}}\right) d\gamma$$
  
$$= g_{1} \exp\left(-\frac{\left(\pi/2\right)^{2}}{2}\right) \sum_{r=1-z}^{\infty} \left(\frac{a_{0}}{\left(\pi/2\right)}\right)^{r}$$
  
$$\times \int_{0}^{\infty} \exp\left(-\frac{\left(a_{0}\right)^{2}\gamma}{2} - \frac{\gamma}{2\sigma^{2}}\right) \gamma^{r+j} I_{r} \left(\frac{a_{0}\pi}{2}\sqrt{\gamma}\right) d\gamma$$
(27)

where  $I_r(\cdot)$  is the modified Bessel function. To proceed, we need to use the series representation of  $I_r(\cdot)$  which is defined by  $I_r(x) = ((1/2)x)^r \sum_{h=0}^{\infty} \frac{((1/4)x^2)^h}{h!\Gamma(r+h+1)}$  [38, eq. (9.6.10)]. Hence,  $D_1$  becomes

$$D_{1} = g_{1} \exp\left(-\frac{(\pi/2)^{2}}{2}\right) \sum_{r=1-z}^{\infty} \left(\frac{a_{0}}{(\pi/2)}\right)^{r} \left(\frac{a_{0}\pi}{4}\right)^{r}$$
$$\times \sum_{h=0}^{\infty} \left(\frac{\left(\frac{a_{0}\pi}{2}\right)^{2} \frac{1}{4}}{h!\Gamma(r+h+1)}\right)^{h}$$
$$\times \int_{0}^{\infty} \exp\left(-\frac{(a_{0})^{2}\gamma}{2} - \frac{\gamma}{2\sigma^{2}}\right) \gamma^{r+j+h} d\gamma \qquad (28)$$

By using  $[31, eq. (3.381.4)], D_1$  is derived as

$$D_{1} = g_{1} \exp\left(-\frac{(\pi/2)^{2}}{2}\right) \sum_{r=1-z}^{\infty} \left(\frac{a_{0}}{(\pi/2)}\right)^{r} \left(\frac{a_{0}\pi}{4}\right)^{r} \times \sum_{h=0}^{\infty} \left(\frac{\left(\frac{a_{0}\pi}{2}\right)^{2} \frac{1}{4}}{h!\Gamma(r+h+1)}\right)^{h}$$

$$\times \left(\frac{(a_0)^2}{2} + \frac{1}{2\sigma^2}\right)^{-(r+j+h+1)} \Gamma(r+j+h+1)$$
(29)

Then,  $D_2$ - $D_6$  are obtained by following the same analytical steps used for  $D_1$  as

$$D_{2} = \frac{4}{M} \exp\left(-\frac{(\pi/2)^{2}}{2}\right) \sum_{r=1-z}^{\infty} \left(\frac{(a_{1})^{2}}{2}\right)^{r} \times \sum_{h=0}^{\infty} \left(\frac{\left(\frac{a_{1}\pi}{2}\right)^{2} \frac{1}{4}}{h!\Gamma(r+h+1)}\right)^{h} \times \left(\frac{(a_{1})^{2}}{2} + \frac{1}{2\sigma^{2}}\right)^{-(r+j+h+1)} \Gamma(r+j+h+1)$$
(30)

$$D_{3} = g_{2} \exp\left(-\frac{(\pi/4)^{2}}{2}\right) \sum_{r=1-z}^{\infty} \left(\frac{(a_{0})^{2}}{2}\right)^{r} \\ \times \sum_{h=0}^{\infty} \left(\frac{(\frac{a_{0}\pi}{2})^{2} \frac{1}{4}}{h!\Gamma(r+h+1)}\right)^{h} \\ \times \left(\frac{(a_{0})^{2}}{2} + \frac{1}{2\sigma^{2}}\right)^{-(r+j+h+1)} \Gamma(r+j+h+1)$$
(31)

$$D_{4} = \frac{8}{M} \sum_{k=1}^{W-1} \exp\left(-\frac{(\alpha_{k})^{2}}{2}\right) \sum_{r=1-z}^{\infty} \left(\frac{a_{0}}{\alpha_{k}}\right)^{r} \left(\frac{a_{0}\alpha_{k}}{2}\right)^{r} \\ \times \sum_{h=0}^{\infty} \left(\frac{(a_{0}\alpha_{k})^{2} \frac{1}{4}}{h!\Gamma(r+h+1)}\right)^{h} \\ \times \left(\frac{(a_{0})^{2}}{2} + \frac{1}{2\sigma^{2}}\right)^{-(r+j+h+1)} \Gamma(r+j+h+1)$$
(32)

$$D_{5} = \frac{4}{M} \sum_{k=1}^{w-1} \exp\left(-\frac{(\beta_{k}^{+})^{2}}{2}\right) \sum_{r=1-z}^{\infty} \left(\frac{a_{k}}{\beta_{k}^{+}}\right)^{r} \left(\frac{a_{k}\beta_{k}^{+}}{2}\right)^{r} \times \sum_{h=0}^{\infty} \left(\frac{(a_{k}\beta_{k}^{+})^{2}\frac{1}{4}}{h!\Gamma(r+h+1)}\right)^{h} \times \left(\frac{(a_{k})^{2}}{2} + \frac{1}{2\sigma^{2}}\right)^{-(r+j+h+1)} \Gamma(r+j+h+1)$$
(33)

$$D_{6} = \frac{4}{M} \sum_{k=2}^{w} \exp\left(-\frac{(\beta_{k}^{-})^{2}}{2}\right) \sum_{r=1-z}^{\infty} \left(\frac{a_{k}}{\beta_{k}^{-}}\right)^{r} \left(\frac{a_{k}\beta_{k}^{-}}{2}\right)^{r} \times \sum_{h=0}^{\infty} \left(\frac{(a_{k}\beta_{k}^{-})^{2}\frac{1}{4}}{h!\Gamma(r+h+1)}\right)^{h} \times \left(\frac{(a_{k})^{2}}{2} + \frac{1}{2\sigma^{2}}\right)^{-(r+j+h+1)} \Gamma(r+j+h+1)$$
(34)

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**FIGURE 1.** ASER performance of 4 × 2 RQAM with K = 30,  $\Delta = 0.5$  and different *m* parameter values.



**FIGURE 2.** ASER performance of 8 × 4 RQAM with m = 5,  $\Delta = 0.5$  and different *K* parameter values.

Then, an ASER expression for XQAM in the presence of FTR fading is obtained by inserting (29)-(34) into (26). To the best of the authors' knowledge, an ASER expression for XQAM over FTR fading has not been reported in the literature, yet.

#### **IV. NUMERICAL RESULTS**

We provide the numerical results for the proposed ASER expressions of RQAM and XQAM to show the effect of fading parameters of FTR distribution. It should be noted that infinite series are truncated after 40*th* terms to accurately calculate the derived expressions for all results in this paper. Analytical results of RQAM and XQAM are approved by the numerical solution of (5) and (26) by using Mathematica and Matlab softwares. Figs. 1 and 2 compare the *Pe<sub>Chiani</sub>* and *Pe<sub>Chernoff</sub>* performances of 4 × 2 RQAM and 8 × 4

**TABLE 1.** Absolute relative difference between exact and approximate values of the chiani and chernoff approximations for ASER of 4 × 2 RQAM with K = 30,  $\Delta = 0.5$ , m = 10.

SNR (dB)	Approximation	R:Diff. (%)	Exact
30	Chiani: 4.420452 x10 <sup>-6</sup> Chernoff: 8.993893 x10 <sup>-6</sup>	14.5 133.0	3.859998 x10 <sup>-6</sup>
35	Chiani: 8.483311 x10 <sup>-7</sup> Chernoff: 1.568500 x10 <sup>-6</sup>	11.5 106.3	7.603103 x10 <sup>-7</sup>
40	Chiani: 2.253308 x10 <sup>-7</sup> Chernoff: 4.005226 x10 <sup>-7</sup>	10.2 95.9	2.044034 x10 <sup>-7</sup>

**TABLE 2.** Absolute relative difference between exact and approximate values of the chiani and chernoff approximations for ASER of 8 × 4 RQAM with K = 15,  $\Delta = 0.5$ , m = 5.

SNR (dB)	Approximation	R:Diff. (%)	Exact
40	Chiani: 0.000094 Chernoff: 0.000164	10.58 92.94	0.000085
45	Chiani: 0.000028 Chernoff: 0.000048	12.0 92.0	0.000025
50	Chiani: 8.831291x10 <sup>-6</sup> Chernoff: 0.000015	10.08 86.97	8.022451x10 <sup>-6</sup>

RQAM over FTR fading channels for different configurations of parameters.

While Fig. 1 shows the effect of the *m* parameter for  $4 \times 2$ RQAM with K = 30,  $\Delta = 0.5$ , the impact of K parameter for 8  $\times$  4 RQAM with  $\Delta = 0.5$ , m = 5 is presented in Fig. 2. In Fig 1, it can be observed that as value of the mparameter increases, the performance of  $4 \times 2$  RQAM system over mmWave FTR fading channels improves. For example, an ASER of  $10^{-4}$  occurs at  $\overline{\gamma} \approx 30$  dB when m = 5 and the same ASER value occurs at  $\overline{\gamma} \approx 23$  dB when m = 10. On the other hand in Fig 2, an ASER level of  $10^{-3}$  takes place at  $\overline{\gamma} \approx 28$  dB when K = 30, the same ASER result with K = 15 occurs at  $\overline{\gamma} \approx 33$  dB. It can be concluded that increasing value of K provides performance improvement of  $8 \times 4$  RQAM system over mmWave FTR fading channels. Furthermore, the asymptotic behaviors of both  $4 \times 2$  RQAM and  $8 \times 4$  RQAM are given in Figs 1 and 2. It should be emphasized that the asymptotic curves are close to analytical curves and simulations at high SNR regime. Besides, when the 40 terms are used for the infinite series, Table I tabulates the absolute relative difference values for ASER of  $4 \times 2$ RQAM with  $K = 30, \Delta = 0.5, m = 10$  over mmWave FTR fading channels.

We can see from these results that the simulations of the RQAM over FTR fading become very tight with theoretical results for all SNR values. From the Fig 1, one can see that the best performance is obtained with m = 10. It is observed that  $Pe_{Chiani}$  results are closer than the  $Pe_{Chernoff}$  results to the analytical findings for all fading parameter configurations since the Chiani approximation is tighter than the Chernoff approximation. The Chiani approximation for Gaussian Q-function includes two exponential terms, so it ensures



FIGURE 3. ASER performance of 32-XQAM over FTR fading channels.

preferable agreement with the simulations when compared to the Chernoff approximation which contains one exponential term for Gaussian *Q*-function. For this reason, the ASER results of RQAM over mmWave FTR channels obtained from the Chiani approximation indicate good agreement to the simulations. This can be viewed as the reason for the deviation of the results obtained by two bounds.

The absolute relative difference percentage for ASER of  $8 \times 4$  RQAM with K = 15,  $\Delta = 0.5$ , m = 5 over mmWave FTR fading channels is given in Table II by using the first 40 terms for the infinite series in the approximate ASER expressions. In addition, Table III presents the results of evaluating the ASER expressions in (14) and (23) at a certain upper limit namely J. The table demonstrates the decimal places that have not been affected with adding more terms to the summations, in bold. From table III, it is clear that the first 40 terms for the ASER calculation of RQAM is quite sufficient.

In Fig. 3, the ASER performance of 32-XQAM under FTR fading conditions is illustrated when  $K = 30, \Delta = 0.5$ , m = 4, 5 and 10. The ASER performance of XQAM scheme with m = 10 outperforms the performance of XQAM with m = 4, 5 as expected since the higher m parameter values correspond to good fading conditions in FTR channels. In this figure, the results of the special cases are provided for mmWave FTR fading channels as Nakagamim case  $(K \rightarrow \infty, \Delta = 0, m = m)$  and Rayleigh case  $(K \rightarrow \infty, \Delta = 0, m = 1)$  are included. The case of  $\Delta = 0, K \rightarrow \infty, m = 1$  represents NLOS environments as the channel model turns into Rayleigh fading. In order to reflect the LOS effects, the fading parameters should be set as  $\Delta = 0, K = K, m \to \infty$  which corresponds to the case of Rician fading. It should be noted that the analytical results are very close to the exact simulations for 32-XQAM. By these means, the accuracy of analytical work proposed in this study is demonstrated. In Table IV, the results of evaluating the ASER expression of XQAM by inserting (29)-(34) into (26)

**TABLE 3.** Values of ASER with cherfnoff (14) and chiani (23) approximations for different levels of truncation ( $\overline{\gamma}$  = 35 dB, K = 15,  $\Delta$  = 0.5, m = 5).

J	Chernoff Approx. eq. (14)	Chiani Approx. eq. (23)
1	0.00001682165338591817	<b>9.</b> 417798128578183x10 <sup>-6</sup>
5	0.00001730452166674863	<b>9.</b> 604537283503918x10 <sup>-6</sup>
10	0.00001730454603396355	9.604543115246886x10 <sup>-6</sup>
20	0.00001730454603336928	9.604543115271153x10 <sup>-6</sup>
30	0.00001730454603398962	9.604543115251713x10 <sup>-6</sup>
40	0.00001730454603398962	<b>9.604543115251713</b> x10 <sup>-6</sup>

**TABLE 4.** Values of ASER FOR 32-XQAM with H = 100 and R = 40 for different levels of truncation ( $\overline{\gamma} = 45 \text{ dB}$ , K = 30,  $\Delta = 0.5$ , m = 4).

J	Analytical expression eq. (26)
1	0.00028722918305668550
5	<b>0.000</b> 44065811379191347
10	0.00044110649292435000
20	0.00044110650224934306
30	0.00044110650223034946
40	0.00044110650223034946

at upper limits J, H and R are illustrated. It is seen that 40 terms for the ASER computation of XQAM are enough to get a given accuracy in this paper.

#### **V. CONCLUSION**

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In this study, we have presented the ASER performance of RQAM and XQAM signaling schemes over millimeter wave FTR fading channels. New and novel expressions are derived and evaluated for various configurations of channel parameters. An exact ASER expression of RQAM over millimeter wave FTR fading channels is obtained. In addition to this, two upper bound ASER expressions are proposed for RQAM scheme based on two different approximations of Gaussian *Q*-function for comparison purpose. Then, an exact ASER expression is obtained for the ASER of XQAM technique in terms of infinite series representation. Furthermore, the analytical results are validated by exact simulations which show the accuracy of proposed analysis in this paper.

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