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A Memristive Chaotic Oscillator With Increasing Amplitude and Frequency

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ABSTRACT A chaotic oscillator utilizing a flux-controlled memristor to produce a signal that grows in amplitude and frequency over time is introduced in this paper. It was found that the initial condition can be used to change the starting oscillation as well as the amplitude and frequency. From this, a new regime of homogenous multistability was found, where various attractors with different initial conditions are of the same type but have different amplitudes and frequencies.

INDEX TERMS Homogenous multistability, increasing amplitude and frequency, memristive chaotic oscillator.

I. INTRODUCTION

Chaotic signals have great potential in engineering applications, including secure communication [1]–[3], image encryption [4]–[6], or weak signal detecting [7], [8]. Chaotic signals with a controllable amplitude [9]–[11] or frequency [12], [13] especially have promising applications since they do not need an extra peripheral to provide pre-modulation. Many chaotic circuits use memristors, a two-terminal electronic device, which was postulated in [14] and [15] and has potential applications in the next generation computers and powerful intelligent devices [16], [17]. In this paper we investigate whether a memristor can be utilized to control the amplitude or even the frequency of a chaotic signal. We further investigate the property of multistability in this system, which has attracted great interest in the field of complex problems, such as the phenomenon of conditional symmetric attractors in an asymmetric system [18]–[20] or infinitely many attractors [21], [22] in a periodically offset-boostable system that has not been investigated before. In a pinched hysteresis loop (the finger print of a memristor), it is common to see multistability in memristive chaotic systems, where different attractors [23]–[26] or even extremely multistable attractors are found [27]–[29]. Multistability with the same class of attractors can be called “homogenous

multistability” in contrast to “heterogeneous multistability”. Memristive dynamical systems have abundant multistable attractors, but the regime of homogenous multistability was rarely reported.

When a chaotic system is equipped with a memristor for amplitude-frequency control, the memristive system will have two unusual properties: a chaotic signal with growing amplitude/frequency and infinitely many homogenous attractors. In Sec. 2, a model of the new memristive chaotic system is proposed. In Sec.3, the properties of the memristive system are demonstrated. In Sec.4, the system was implemented as a chaotic circuit with the experimental results showing agreement with the predicted simulation. The conclusion and discussion is given in the last section.

II. MODEL DESCRIPTION

A new memristive chaotic system is described as,

$$\begin{cases} \dot{x} = W(u)y + yz, \\ \dot{y} = yz - axz, \\ \dot{z} = bz^2 - y^2, \\ \dot{u} = y. \end{cases} \quad (1)$$

where the flux-controlled memductance in the first dimension is $W(u) = k(|u| - c)$ is introduced in the first dimension.

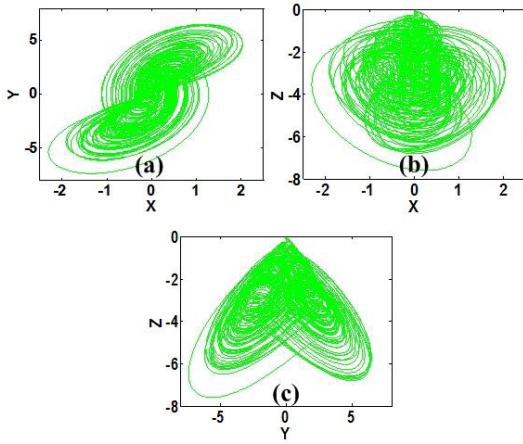


FIGURE 1. Projections of the chaotic attractor from System (1) with initial condition $[1, 0, -1, 1]$ when $a = 5.5, b = 0.55, k = 0.05, c = 1$ and the time duration is 600.

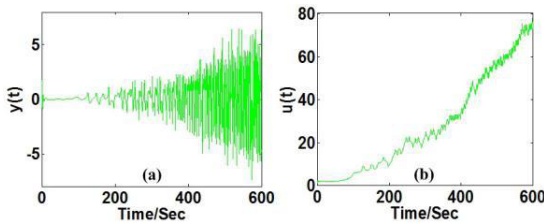


FIGURE 2. Chaotic oscillation and the internal state from System (1) with initial condition $[1, 0, -1, 1]$ when $a = 5.5, b = 0.55, k = 0.05, c = 1$ and the time duration is 600.

When $a = 5.5, b = 0.55, k = 0.05, c = 1$, System (1) has a chaotic attractor with Lyapunov exponents $(0.1375, 0.0068, -0.0086, -2.2962)$ with initial conditions $(1, 0, -1, 1)$ and increases in amplitude and frequency over a duration of 600, as shown in Fig.1. The internal state u is determined by the integration of the memristor input variable y . As shown in Fig.2, the increasing value of the control variable u increases the frequency of signals x, y and z . Here the internal state of u is continuously growing which is different from some of other memristor models because the variable y grows with the memductance $W(u)$ and give a positive feedback to the internal state.

System (1) retains the rotational symmetry with a rate of hypervolume contraction given by the Lie derivative, $\nabla V = \frac{\partial \dot{x}}{\partial x} + \frac{\partial \dot{y}}{\partial y} + \frac{\partial \dot{z}}{\partial z} + \frac{\partial \dot{u}}{\partial u} = (2b + 1)z$. In contrast to other memristive systems, the above system has an infinite plane of equilibria $(x, 0, 0, u)$ with eigenvalues $(0, 0, 0, 0)$.

Here the flux-controlled memristor is defined as,

$$\begin{cases} i = W(u)y, \\ W(u) = 0.05|u| - 0.05, \\ \dot{u} = y. \end{cases} \quad (2)$$

The flux-controlled memductance is a function of internal state of u , and is associated with the voltage y ,

$$\begin{aligned} W(u) &= 0.05|u| - 0.05 = 0.05 \left| \int_{-\infty}^t y ds \right| - 0.05 \\ &= W_0 + 0.05 \left| \int_0^t y ds \right| - 0.05 \end{aligned} \quad (3)$$

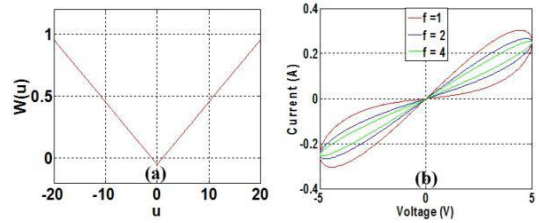


FIGURE 3. The memductance and pinched hysteresis loop.

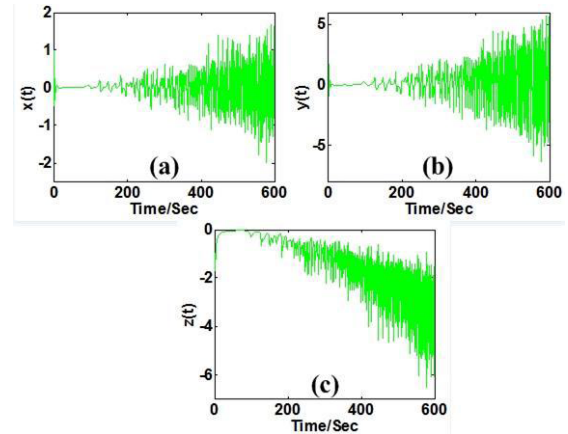


FIGURE 4. Chaotic signal of System (1) with $k = 0.05, c = 1$ and initial conditions $[1, 0, -1, 1]$ over a time duration of 600.

where $W_0 = 0.05(|\int_{-\infty}^t y ds| - |\int_0^t y ds|)$. The memductance and the theoretical pinched hysteresis loop are shown in Fig.3.

III. UNIQUE DYNAMICAL BEHAVIOR

A. AMPLITUDE AND FREQUENCY BOOSTING OVER TIME

The original system without a memristor has two coefficients to control the amplitude of the signal [12],

$$\begin{cases} \dot{x} = ny + yz, \\ \dot{y} = yz - axz, \\ \dot{z} = bz^2 - my^2. \end{cases} \quad (4)$$

where parameter n controls the amplitude and frequency of all the signals x, y and z , while m controls the amplitude of x and y . In System (1), the added memristor can control the amplitude and frequency. As the internal state u and consequently the memductance increases over time, so does the amplitude and frequency. As predicted, the chaotic signals in the time domain are modulated by the memristor with increasing amplitude and frequency, as shown in Fig.4. The phase trajectories over time are shown in Fig.5, clearly showing chaotic signals increasing in amplitude and frequency. The frequency spectra of the chaotic signals are shown in Fig. 6. As a comparison, we list the typical information including the average amplitude (AA) and Largest Lyapunov exponent (LLE) (corresponding to a boosted frequency) in Table.1. Here the increasing amplitude is indicated by the growing average value of the absolute values of corresponding chaotic signals

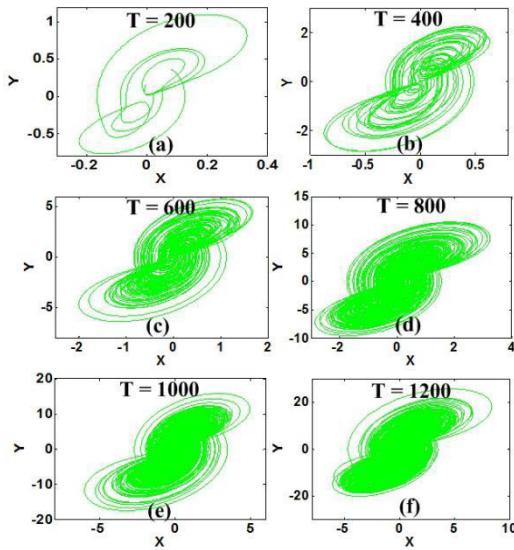


FIGURE 5. Different phase trajectories of System (1) with $k = 0.05$, $c = 1$ and various durations of time under the same initial condition $[1, 0, -1, 1]$.

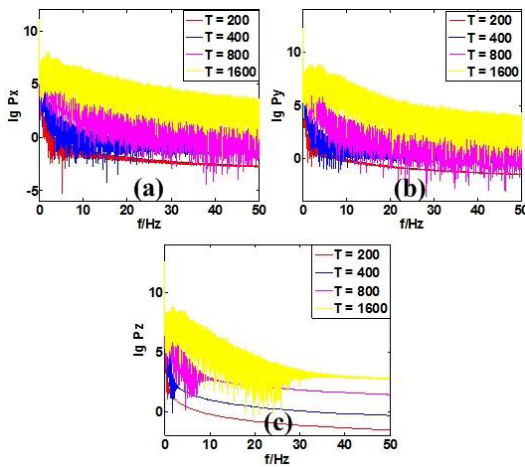


FIGURE 6. Frequency spectra of the chaotic signals from System (1) with $k = 0.05$, $c = 1$ and initial condition $[1, 0, -1, 1]$ when the run time is prolonged.

while the increasing frequency is proven by the increasing largest Lyapunov exponent based on the Wolf algorithm.

B. HOMOGENOUS MULTISTABILITY: INITIAL-CONDITION-TRIGGERED AMPLITUDE AND FREQUENCY BOOSTING

When a dynamical system has coexisting attractors of the same shape but with different positions, amplitude or even frequency, the regime of multistability can be defined as homogenous multistability. Here in System (1), we find a new regime of homogenous multistability, where different initial condition can trigger the same oscillation but with different amplitudes and frequency. In fact, the initial condition determines the start amplitude and frequency of the oscillation process. For example, fix the initial state as $(x_0, y_0, z_0) = (1, 0, -1)$ and the time duration 500, revise the

TABLE 1. Average amplitude (AA) of chaotic signal of System (1) with $k = 0.05$, $c = 1$ and the initial condition $[1, 0, -1, 1]$ and largest Lyapunov exponent (LEE).

Cases	T	AA: (mean(x , y , z , u))	LEE
A1	200	(0.0449, 0.1443, 0.2208, 5.8607)	0.0398
B1	400	(0.1212, 0.4327, 0.6578, 14.8009)	0.0701
C1	600	(0.2465, 0.8837, 1.3491, 29.2505)	0.1526
D1	800	(0.3837, 1.3842, 2.1138, 44.8955)	0.2348
E1	1000	(0.5265, 1.8868, 2.8919, 61.1531)	0.3418
F1	1200	(0.6748, 2.4159, 3.7028, 77.7349)	0.4492

TABLE 2. Average amplitude of chaotic signal of System (1) with $k = 0.05$, $c = 1$ and the initial condition $[1, 0, -1, u_0]$ and largest Lyapunov exponent when $T = 1000$.

Cases	u_0	AA: (mean(x , y , z , u))	LEE
A2	0.5	(0.3212, 1.1674, 1.7779, 37.7247)	0.2063
B2	1	(0.5265, 1.8868, 2.8919, 61.1531)	0.3202
C2	2	(0.6698, 2.4210, 3.6921, 76.8133)	0.5137
D2	4	(0.8422, 3.0538, 4.6563, 96.4674)	0.5862
E2	8	(0.7143, 2.5771, 3.9379, 82.1217)	0.6646
F2	16	(1.1997, 4.3158, 6.5991, 136.4677)	0.7262

initial value of the internal state u_0 , chaotic signals stand different stages of amplitude and frequency. As shown in Fig. 7, when the initial state u_0 is 0.5, the amplitude and frequency are both in small scale. The same trend can be seen in the average value and the largest Lyapunov exponent, as shown in Table.2. As shown in Fig.8, when the initial condition u_0 varies in $[0, 8]$, the average of the state variable u revise positively and dramatically, while the average of the variable x , y , and z will increase accordingly. The increasing frequency can be identified by the slowly climbing slope of the largest

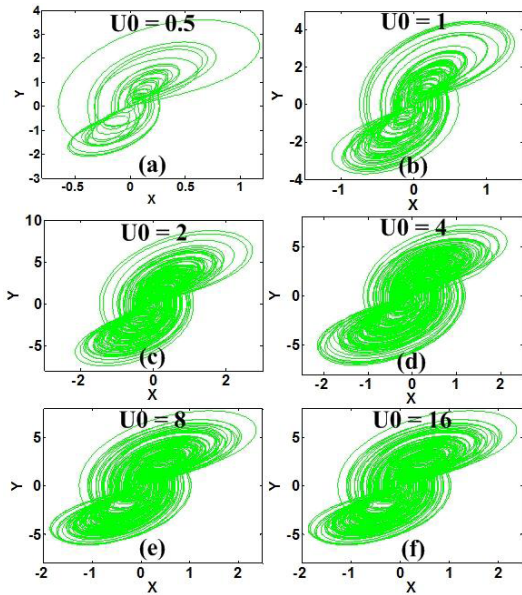


FIGURE 7. Phase trajectories of System (1) with $k = 0.05$, $c = 1$ and initial conditions $(1, 0, -1, u_0)$ under the same time ($T = 500$).

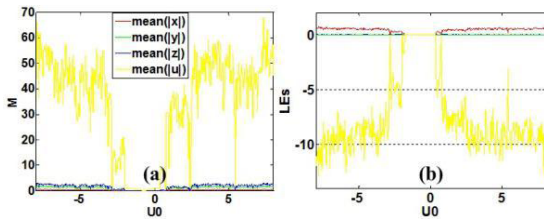


FIGURE 8. The average value of the chaotic signals and the corresponding Lyapunov exponents of System (1) with $k = 0.05$, $c = 1$ and the initial conditions $(1, 0, -1, u_0)$, u_0 varies in $[-8, 8]$. Here the run time of T is 1000.

Lyapunov exponent. In the other direction when the internal state u_0 varies in $[-8, 0]$, the amplitude and frequency evolution shows the same regularity as in the positive direction since flux-controlled memductance is an absolute function of the internal state. Frequency spectra in Fig. 9 show the same characteristic in oscillation. Typical information including the average amplitude and the Largest Lyapunov exponent (corresponds to a boosted frequency) is listed in Table.2.

Moreover, a crisis in amplitude control still exists [12]. The boosting of the amplitude and frequency may encounter a risk since some initial conditions can lead to a direct death of oscillation. For example, when $(x_0, y_0, z_0, u_0) = (-1, 0, -1, 0.5)$, $(1, 5, -1, 0.5)$, $(1, 0, 5, 0.5)$, $(1, 0, -1, -0.5)$, the oscillation is stopped by the initial condition and System (1) converges to a fixed point. Note that some initial conditions will result in a state of suspended animation, such as in $(-20, 0, -1, 0.5)$, System (1) comes back to the oscillation until the time is about over 850, as shown in Fig.10. The initial condition in the x dimension can also adjust the amplitude and frequency in a way which is not a positive correlation and different from the u dimension, as shown in Fig. 11. All the coexisting attractors as shown in Fig.7 and Fig.11 indicate the special regime of homogenous multistability.

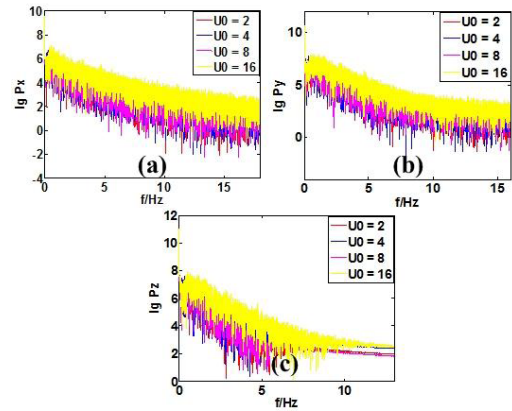


FIGURE 9. Frequency spectra of the chaotic signals of System (1) with $k = 0.05$, $c = 1$ and initial conditions $(1, 0, -1, u_0)$. Here the run time of T is 1000.

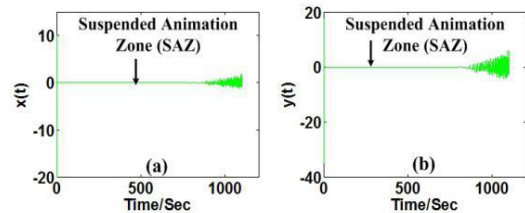


FIGURE 10. Suspended animation in System (1) with $k = 0.05$, $c = 1$ and initial conditions $[-20, 0, -1, 0.5]$ under the time duration 1100.

IV. CIRCUIT IMPLEMENTATION

From Eq. (1), we design the analog circuit shown in Fig. 12 where the circuit equations in terms of the circuit parameters are

$$\begin{cases} \dot{x} = \frac{1}{C_1}W(u)y + \frac{1}{R_1C_1}yz, \\ \dot{y} = \frac{1}{R_2C_2}yz - \frac{1}{R_3C_2}xz, \\ \dot{z} = \frac{1}{R_4C_3}z^2 - \frac{1}{R_5C_3}y^2. \end{cases} \quad (5)$$

where memristor $W(u)$ in Fig. 13 and its state is defined as,

$$\begin{cases} i = W(u)y, \\ W(u) = \frac{R_f}{R_d}|u| - \frac{R_f}{R_e}V_{dd}, \\ \dot{u} = \frac{1}{R_aC_a}y. \end{cases} \quad (6)$$

The system was rescaled by $u \rightarrow 10u$ to delay the saturation time of the circuit. The circuit consists of four channels to realize the integration, addition, and subtraction of state variables x, y, z and u . The operational amplifier TL084 performs the addition, inversion, and integration, and the analog multiplier AD633/AD performs the nonlinear product operation. The circuit is powered by $\pm 15V$. The state variables $x, y,$ and z in Eq. (1) correspond to the state voltages of the three channels, while u corresponds to the internal state of the memristor. For the system parameters $a = 5.5, b = 0.55, k = 0.05, c = 1$, the circuit element values are $C_1 = C_2 = C_3 = 1\mu F, R_1 = R_2 = R_5 = R_6 = R_7 = R_8 = R_9 = 1k\Omega,$

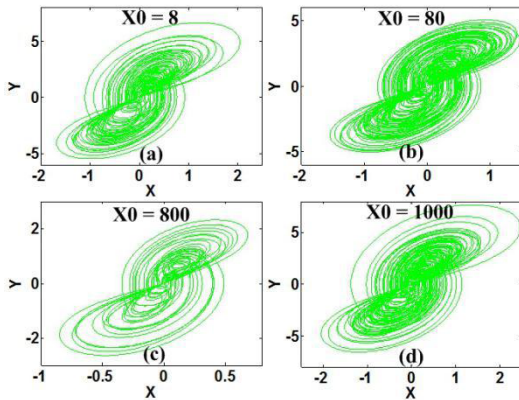


FIGURE 11. Phase trajectories of System (1) with $k = 0.05$, $c = 1$ and initial conditions $[x_0, 0, -1, 0.5]$ under the same time ($T = 500$).

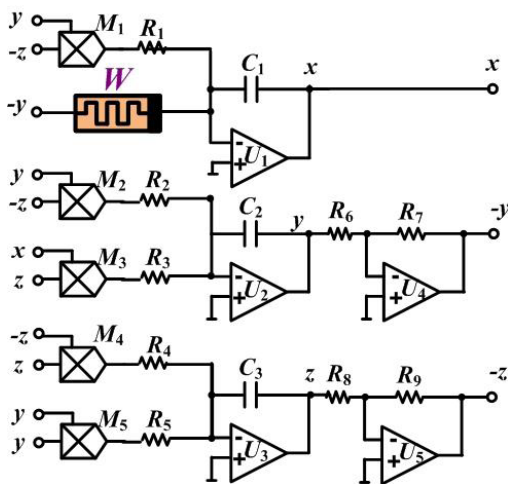


FIGURE 12. Circuit schematic of the memristive system.

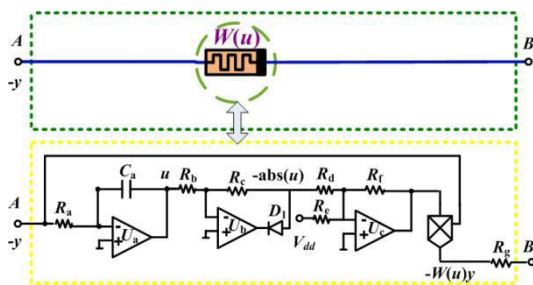


FIGURE 13. Equivalent circuit of the flux-controlled memristor.

$R_3 = 180k\Omega$ and $R_4 = 1.8k\Omega$. The memristor circuit is shown in Fig. 13 and has the following parameter values, $C_4 = 1\mu F$, $R_a = 100k\Omega$, $R_b = 470\Omega$, $R_c = 470\Omega$, $R_d = 2k\Omega$, $R_e = 300k\Omega$, $R_f = 1k\Omega$ and $R_g = 1k\Omega$. The time scale is set at 1000 to observe the increase in frequency and amplitude before saturation. Figure 14 shows a plot of the experimental memductance and pinched hysteresis loop, while Figure 15 shows the time evolution of the chaotic signal with increasing amplitude and frequency. Figure 16 shows the phase portraits observed in the oscilloscope.



FIGURE 14. Plot of experimental memductance and pinched hysteresis loop (CH1: u , CH2: $W(u)$) for the left and CH1: input voltage, CH2: the current for the right).

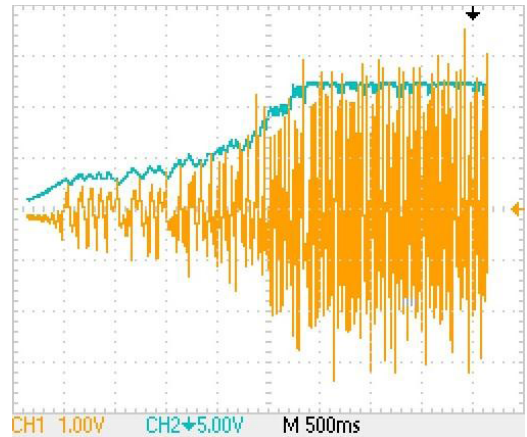


FIGURE 15. Time evolution of the chaotic signal with increasing amplitude and frequency (CH1: x , CH2: u).

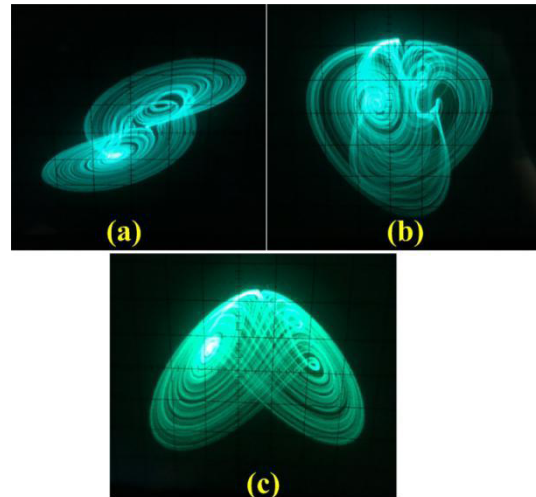


FIGURE 16. Phase trajectory of System (6) from oscilloscope (a) x - y (b) x - z (c) y - z .

V. CONCLUSIONS AND DISCUSSION

When a memristor was introduced into a chaotic system for amplitude-frequency control, a new chaotic oscillator was generated where the output chaotic signals increase in amplitude and frequency over time. The initial conditions can be used to extract chaotic signals with different amplitudes and frequencies, corresponding to a new regime of homogenous multistability which was found and demonstrated. Chaotic signals with increasing amplitude and frequency meet the broad requirement in application engineering. Regarding the

physical implementation, the memristor is utilized as a power element to control the amplitude and frequency, whereas the amplitude is typically controlled by resistor and the frequency is controlled by capacitor.

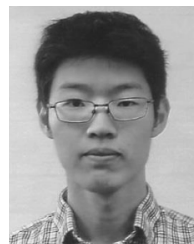
REFERENCES

- [1] S. Wang, J. Kuang, J. Li, Y. Luo, H. Lu, and G. Hu, "Chaos-based secure communications in a large community," *Phys. Rev. E, Stat. Phys. Plasmas Fluids Relat. Interdiscip. Top.*, vol. 66, no. 6, Dec. 2002, Art. no. 065202R.
- [2] C. Y. Chee and D. Xu, "Secure digital communication using controlled projective synchronisation of chaos," *Chaos, Solitons Fractals*, vol. 23, no. 3, pp. 1063–1070, 2005.
- [3] Q. Wang et al., "Theoretical design and FPGA-based implementation of higher-dimensional digital chaotic systems," *IEEE Trans. Circuits Syst. I, Reg. Papers*, vol. 63, no. 3, pp. 401–403, Mar. 2016.
- [4] Z.-H. Guan, F. Huang, and W. Guan, "Chaos-based image encryption algorithm," *Phys. Lett. A*, vol. 346, nos. 1–3, pp. 153–157, 2005.
- [5] T. Gao and Z. Chen, "A new image encryption algorithm based on hyper-chaos," *Phys. Lett. A*, vol. 372, no. 4, pp. 394–400, 2008.
- [6] E. Y. Xie, C. Li, S. Yu, and J. Lü, "On the cryptanalysis of Fridrich's chaotic image encryption scheme," *Signal Process.*, vol. 132, pp. 150–154, Mar. 2017.
- [7] G. Wang, D. Chen, J. Lin, and X. Chen, "The application of chaotic oscillators to weak signal detection," *IEEE Trans. Ind. Electron.* vol. 46, no. 2, pp. 440–444, Apr. 1999.
- [8] G. Wang and S. He, "A quantitative study on detection and estimation of weak signals by using chaotic Duffing oscillators," *IEEE Trans. Circuits Syst. I, Fundam. Theory Appl.*, vol. 50, no. 7, pp. 945–953, Jul. 2003.
- [9] C. Li, J. C. Sprott, and W. Thio, "Linearization of the Lorenz system," *Phys. Lett. A*, vol. 379, nos. 10–11, pp. 888–893, 2015.
- [10] C. Li, J. C. Sprott, Z. Yuan, and H. Li, "Constructing chaotic systems with total amplitude control," *Int. J. Bifurcation Chaos*, vol. 25, no. 10, 2015, Art. no. 1530025.
- [11] C. Li, J. C. Sprott, W. Thio, and H. Zhu, "A new piecewise linear hyper-chaotic circuit," *IEEE Trans. Circuits Syst. II, Exp. Briefs*, vol. 61, no. 12, pp. 977–981, Dec. 2014.
- [12] C. Li, J. C. Sprott, and H. Xing, "Crisis in amplitude control hides in multistability," *Int. J. Bifurcation Chaos*, vol. 26, no. 14, 2016, Art. no. 1650233.
- [13] C. Li and J. C. Sprott, "Chaotic flows with a single nonquadratic term," *Phys. Lett. A*, vol. 378, no. 3, pp. 178–183, 2014.
- [14] L. O. Chua, "Memristor—the missing circuit element," *IEEE Trans. Circuit Theory*, vol. 18, no. 5, pp. 507–519, Sep. 1971.
- [15] L. O. Chua and S. M. Kang, "Memristive devices and systems," *Proc. IEEE*, vol. 64, no. 2, pp. 209–223, Feb. 1976.
- [16] J. M. Tour and T. He, "Electronics: The fourth element," *Nature*, vol. 453, May 2008, Art. no. 42.
- [17] M. Itoh and L. O. Chua, "Memristor oscillators," *Int. J. Bifurcation Chaos*, vol. 18, no. 11, pp. 3183–3206, 2008.
- [18] C. Li, J. C. Sprott, and H. Xing, "Hypogenetic chaotic jerk flows," *Phys. Lett. A*, vol. 380, nos. 11–12, pp. 1172–1177, 2016.
- [19] C. Li, J. C. Sprott, and H. Xing, "Constructing chaotic systems with conditional symmetry," *Nonlinear Dyn.*, vol. 87, no. 2, pp. 1351–1358, 2017.
- [20] C. Li and J. C. Sprott, "Variable-boostable chaotic flows," *Optik-Int. J. Light Electron Opt.*, vol. 127, no. 22, pp. 10389–10398, 2016.
- [21] C. Li, J. C. Sprott, and Y. Mei, "An infinite 2-D lattice of strange attractors," *Nonlinear Dyn.*, vol. 89, no. 4, pp. 2629–2639, 2017.
- [22] C. Li, J. C. Sprott, W. Hu, and Y. Xu, "Infinite multistability in a self-reproducing chaotic system," *Int. J. Bifurcation Chaos*, vol. 27, no. 10, 2017, Art. no. 1750160.
- [23] G. Bao and Z. Zeng, "Multistability of periodic delayed recurrent neural network with memristors," *Neural Comput. Appl.*, vol. 23, nos. 7–8, pp. 1963–1967, 2013.
- [24] B. C. Bao, Q. D. Li, N. Wang, and Q. Xu, "Multistability in Chua's circuit with two stable node-foci," *Chaos*, vol. 26, no. 4, 2016, Art. no. 043111.
- [25] Q. Xu, Y. Lin, B. Bao, and M. Chen, "Multiple attractors in a non-ideal active voltage-controlled memristor based Chua's circuit," *Chaos, Solitons Fractals*, vol. 83, pp. 186–200, Feb. 2016.
- [26] A. Wu and J.-E. Zhang, "Multistability of memristive neural networks with time-varying delays," *Complexity*, vol. 21, no. 1, pp. 177–186, Oct. 2015.
- [27] B. C. Bao, Q. Xu, H. Bao, and M. Chen, "Extreme multistability in a memristive circuit," *Electron. Lett.*, vol. 52, no. 12, pp. 1008–1010, 2016.
- [28] F. Yuan, G. Wang, and X. Wang, "Extreme multistability in a memristor-based multi-scroll hyper-chaotic system," *Chaos*, vol. 26, no. 7, 2016, Art. no. 073107.
- [29] G. Wang, C. Shi, X. Wang, and F. Yuan, "Coexisting oscillation and extreme multistability for a memcapacitor-based circuit," *Math. Problems Eng.*, vol. 2017, 2017, Art. no. 6504969.



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