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# **RESEARCH ARTICLE**

# A Multi-Finger GHz Frequency Doubler Based on Amorphous Indium Gallium Zinc Oxide Thin Film Transistors

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**ABSTRACT** This paper presents a multi-finger doubler based on amorphous-indium gallium zinc oxide (a-IGZO) thin film transistors (TFT) operating at GHz frequency. The doubler and the TFTs have been fabricated in-house with chromium (Cr) gate electrodes. We have used the Rensselaer Polytechnic Institute amorphous silicon (RPI-a) TFT DC model and a Null-Bias method extracted AC model for simulation. The device output power, the threshold voltage mismatch, and breakdown characteristics are explained with the help of the model. The device yield reduces as the number of fingers in the TFTs increase. The model helps explain the doubler's non-idealities and compensate for these transistors' low yield. The peak second harmonic conversion gain of the doubler is measured to be -32 dB at a gate overdrive voltage of around 2.8 V for a 500 MHz input signal. This second harmonic output frequency exceeds the TFT's transit and maximum oscillation frequencies.

**INDEX TERMS** Radio frequency, TFT, a-IGZO, S-parameters, Y-parameters,  $f_T$ ,  $f_{max}$ , frequency doubler, conversion gain.

# I. INTRODUCTION

Metal-oxide thin film transistors (TFT)s have been recently popular in ongoing research for high-frequency circuit design. Compared to Si-based technologies, metal oxide semiconductors are cost-effective and can be fabricated at room temperature. As such, there have been attempts to make RF circuits from TFTs using different semiconductor devices. A frequency doubler was presented in [1] showing a conversion loss of around 44 dB for an input of 100 MHz. The same device was used as a detector to detect signals up to 1 GHz with an equivalent noise power of  $30 \text{ nW}/\sqrt{\text{Hz}}$ . Most recently, Mehlman et al. demonstrated a ZnO-based cross-coupled LC oscillator that operates at 1.25 GHz [2]. The speed of metal oxide TFTs has consistently improved over the past decade. ZnO, as well as a-IGZO TFTs, now demonstrate transit frequency  $(f_T)$  and maximum oscillation frequency  $(f_{max})$  above 1 GHz [3], [4].

However, the power capability of a-IGZO decreases as frequency increases. A frequency doubler can relax the output power versus the frequency trade-off of an a-IGZO-based signal source.

Frequency doublers operating between a few hundred MHz to GHz is nonexistent in the literature. An accurate TFT model is required to simulate amplification, mixing, and impedance matching for GHz circuit design and simulation. Although a SPICE level = 3 model is typically sufficient for fundamental AC analysis and simulation, the discontinuity in the level = 3 models leads to convergence issues for RF design involving non-linear analysis and simulation. Therefore, a model more suitable for amorphous TFTs, namely the RPI-a model, is adopted [5]. Created initially for amorphous Silicon [5], the model has been applied successfully in the MHz frequency range for a-IGZO

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transistors in [6]. Their procedure, however, has been applied to a minimum channel length of 3 µm. The RPI-a model is yet to be applied to the submicron length channel for GHz application. A combined DC and high-frequency model is necessary to simulate the power output and impedance matching for GHz frequency circuit design.

In this paper, we have used a combination of the continuous charge-based RPI-a TFT model and high-frequency elements to address the parasitic components of the TFT. The algorithm presented by [7] has been used to extract the DC model parameters. Moreover, we have used the null-bias method from [8] to extract the parasitic elements. In the same spirit, we have fabricated a four-finger common-gate (CG) doubler in-house, as described in [1] and [4]. A higher number of fingers can increase RF power handling capabilities of TFTs. We measured the CG frequency doubler and compared the measurements and the simulations to show the model's efficacy in RF circuit design. The model is derived from the DC and *S*-parameter measurements and implemented as a Verilog-A code. In that case, one can extensively reduce the cost and time required for modeling and simulation.

Moreover, the shorter length of the a-IGZO channel leads to poor device yield [9]. The yield worsens if the number of fingers increases since the probability of device failure will be higher. The model can substitute the low-yield device, helping us learn about circuit behaviors.

This paper is organized as follows: section II comments on device fabrication; in section III, we present the DC and AC models, their extraction procedure, implementation, and verification. The theoretical background of the CG doubler is presented in section IV; section V shows the measurement results and discussion for the doubler, and section VII concludes this paper.

### **II. DEVICE FABRICATION**

The fabrication steps outlined in section II of [4] are used to ensure consistency and comparability of the channel material. This work presents chromium as the gate electrode instead of the composite Ti/Au/Ti used in [4].

It was shown in the literature that water that adsorbs on the surface of the channel material can substantially affect the characteristics of metal-oxide TFTs, such as hysteresis [10]. Thus, an ALD grown  $Al_2O_3$  encapsulation layer, which is an excellent barrier against the diffusion of water molecules [11], is used to protect the TFTs in our lab environment (24°C, 60%).

As mentioned in [4], the on current and off current for a single  $100 \,\mu\text{m}/0.8 \,\mu\text{m}$  TFT are around  $5.46 \times 10^{-12}$  A and  $1.06 \times 10^{-3}$  A respectively, resulting in an on/off ratio of around  $10^8$ . For a ten-finger TFT, the on current increases to  $8.3 \times 10^{-3}$  A, while the off current stays in the same order, around  $7.36 \times 10^{-12}$  A. Thus the on/off ratio for a ten-finger TFT is around  $10^9$ .

### **III. MODELING OF TFTs**

The single-finger, ten-finger, and doubler were designed on the same substrate. The model of the TFT was developed simultaneously with the measured data. First, the singlefinger TFT was measured for both DC and S-parameters, and the model was extracted. Next, the model was used to simulate the TFT in Cadence, and the simulated data were compared with the measured data. Next, the ten-finger TFT was measured, and the model was verified for multi-finger transistors. The doubler performance was also measured, and the model was then used to simulate, verify and analyze the output results obtained.

### A. DC MODEL PARAMETER EXTRACTION

The drain current equation for a TFT of channel width W and length L according to the RPI-a TFT model [7] is given as

$$I_{d} = \frac{\frac{W}{L}C_{g}\mu_{FE}(V_{gs} - V_{th})V_{ds}}{1 + (R_{s} + R_{d})\left[\frac{W}{L}C_{g}\mu_{FE}(V_{gs} - V_{th})\right]} \times \frac{(1 + \lambda V_{ds})}{\left[1 + \left(\frac{V_{ds}}{\alpha_{sat}(V_{gs} - V_{th})}\right)^{m}\right]^{1/m}},$$
(1)

where,  $V_{ds}$  is the sum of DC and AC drain voltages,  $V_{DS}$  and  $v_{ds}$  respectively,  $V_{gs}$  is the sum of DC and AC gate voltages,  $V_{GS}$  and  $v_{gs}$ , and the other variables signify the properties as mentioned in Table 1. The mobility  $\mu_{FE}$  defines the channel mobility under the influence of gate-source voltage as

$$\mu_{FE} = \mu_0 \left( \frac{V_{gs} - V_{th}}{\text{VAA}} \right)^{\gamma}, \qquad (2)$$

where VAA is a fitting parameter describing the characteristic voltage of field effect mobility.

The two most popular methods in the literature for DC parameters extraction are the 'integral function method' [12] and the integrated non-saturating and self-heating effectbased extraction procedure in [7]. The latter is more suitable for our TFTs as the drain current does not saturate in the shortchannel devices. The modeling steps mentioned in [7] have been applied here.

The extracted parameter values are presented in Table 1. The model has been implemented as a Verilog-A code.

Fig. 1 shows the resultant simulated versus the measured output and transfer behaviors. The simulated output characteristic agrees closely with the measurement, as shown in Fig. 1a. The simulated drain current is underestimated at low  $V_{GS}$ , as shown in Fig. 1b. The model incorporates a fixed  $V_{th}$ . However, the threshold voltage decreases as  $V_{DS}$  increases, leading to this fitting error. The normalized root mean square error (NRMSE) is calculated as NRMSE = RMSE/ $(I_{D,max} - I_{D,min})$  and presented in Table 2. As gate voltage increases, the value of NRMSE decreases, leading to around 1% error at  $V_{GS} = 2$  V.

### B. RF MODEL PARAMETER EXTRACTION

One of the most critical parameters in an RF circuit is the gate resistance  $R_G$  that defines the maximum oscillation frequency  $f_{max}$ . The value of  $f_{max}$  additionally depends on the transit frequency  $f_T$  of the transistor. The transconductance and parasitic capacitances play a significant role in determining  $f_T$ .

Parameter	Unit	Description	Extracted
		-	Value
W	μm	Channel width	100
L	μm	Channel length	0.8
$\alpha_{sat}$	-	Saturation modulation parameter	0.906
$\epsilon_i$	-	Relative permitivitty of the insu-	8.5
Ť		lator	
$\gamma$	-	Power law mobility parameter	1.014
$\lambda$	1/V	Output conductance parameter	0.101
m	-	Knee shape parameter	1.3698
$\mu_0$	$m^2/Vs$	Conduction band mobility	0.001
$R_d$	m	Drain contact resistance	160.077
$R_s$	m	Source contact resistance	160.077
$T_{nom}$	°C	Nominal temperature	24
$t_{ox}$	m	Oxide $(Al_2O_3)$ thickness	5e-8
VAA	V	Characteristic voltages for field	1.902
		effect mobility	
$V_{th}$	V	Zero bias threshold voltage	-0.948
$R_{G1}$	Ω	First part of gate electrode resis-	385
0.1		tance	
$R_{G2}$	Ω	Second part of gate electrode re-	1078
		sistance	
$R_{DSE}$	Ω	Drain-source electrode	10
202		resistance	
$C_{PG}$	fF	Gate pad capacitance	36.5
$C_{PD}$	fF	Drain pad capacitance	17
$\hat{C}_{OV}$	fF	Gate-drain or gate-source over-	57
01		lap capacitance	
$r_{ad}$	Ω	Distributed gate-drain channel	0
gu		resistance	
$r_{as}$	Ω	Distributed gate-source channel	650
3~		resistance	
$\tau$	ns	Channel transit delay	150

# TABLE 1. Extracted RPI-a parameters and RF element values for a single-finger TFT.



FIGURE 1. Simulated versus measured (a) output and (b) transfer characteristics (solid lines - simulation and circular markers - measurement) for a single-finger TFT.

 TABLE 2. Normalized root mean square error in the DC output characteristics simulation.

$V_{GS}$ (V)	-1.2	-0.4	0.4	1.2	2
RMSE (mA)	0.003	0.056	0.118	0.096	0.015
$I_{D,max} - I_{D,min}$ (mA)	0.002	0.039	0.244	0.710	1.41
NRMSE	1.5	1.438	0.484	0.136	0.011

These device parameters depend on the material, geometry, and fabrication technique.

The top view of the layout of a  $100 \,\mu$ m/0.8  $\mu$ m singlefinger TFT is shown in Fig. 2a. The source terminal is large and spans both sides of the channel to facilitate



FIGURE 2. Photos (background edited) of the top views of (a) a single-finger and (b) a ten-finger TFTs with dimensions in  $\mu m$ .



**FIGURE 3.** Two lumped embedded model (red represents the intrinsic parameters).

the multi-finger architecture, as shown in the photos in Fig. 2b. Compared to the single-finger TFT, a multi-finger architecture provides a lower gate resistance, resulting in low reflected power at the input.

For high-frequency modeling of the TFT, one must separate the intrinsic and extrinsic components of the devices, as shown in Fig. 3. Here,  $C_{PG}$  and  $C_{PD}$  represent the gate and drain pad capacitance, respectively.  $C_{OV}$  is gate-drain and gate-source overlap capacitance.  $R_{G1}$  and  $R_{G2}$  are two parts of the gate resistance separated by  $C_{PG}$ .  $R_{DSE}$  represents the combined drain and source electrode resistance. The elements highlighted in red represent the intrinsic components.  $g_{m,nqs}$  and  $g_{ds,nqs}$  are the non-quasi-static (NQS) transconductance, and channel conductance, respectively [13].

The extrinsic components are derived at null-bias condition at  $V_{DS} = 0$  V and  $V_{GS} = -2$  V [8]. The gate resistance has two components, namely  $R_{G1}$  and  $R_{G2}$ . An initial calculation of the ratio  $r = R_{G2}/R_{G1}$  is first done according to formula

$$r = \frac{\#\text{squares beyond } C_{PG}}{3 \cdot \#\text{squares before } C_{PG}} = \frac{110/0.8 + 20/1.1}{3 \cdot (2 + 20/1.1)} = 2.57.$$
(3)

A factor of 3 in the denominator is used, as the gate electrode will follow the RC ladder rule beyond  $C_{PG}$  [14]. However, due to the non-uniformity of the gate electrode above the source and drain (see Fig. 2a), we must adjust the value of r slightly. First, an estimation of  $R_{G1} = 400 \Omega$  is done from the sheet resistance of the chromium electrode, which is around  $20 \Omega / \Box$  [4]. Then fitting is used until the input admittance  $Y_{11}$  is identical for both measurement and simulation. Thus, the value of r is determined to be 2.8, while the value of  $R_{G1}$  is 385  $\Omega$ .

Next,  $R_{G1}$  is subtracted from the null-biased Z-parameters and then converted to Y-parameters, from which the values of the pad capacitance  $C_{PG}$ ,  $C_{PD}$ , and the overlap capacitance



**FIGURE 4.** Extracted (a) extrinsic capacitances after de-embedding of  $R_{G1}$  and (b)  $R_{G2}$  and  $R_{DSE}$  over frequency.

 $C_{OV}$  are extracted. Fig. 4a shows the extrinsic capacitances over frequency. The values at a lower frequency give the correct approximation of the capacitive elements. At higher frequencies, the extrinsic capacitances show dispersion due to the presence of  $R_{G2}$  and  $R_{DSE}$ .

Then, the extracted pad capacitances are de-embedded, and  $R_{G2}$  and  $R_{DSE}$  are extracted from the de-embedded Z-parameters, as shown in Fig. 4b. The values at higher frequencies are taken for the resistance estimation [15]. The extracted values are not constant and vary over frequency. Hence, fitting around the average extracted values is done in Cadence Spectre to match the extrinsic Y-parameters. The resultant extrinsic resistances for the single-finger TFT are  $R_{G2} = 1078 \Omega$  and  $R_{DSE} = 10 \Omega$ .

# C. INTRINSIC EXTRACTION AND MODELING

The intrinsic model for NQS extraction has been adopted from [13], as shown in Fig. 3. Here,  $r_{gs}$  and  $r_{gd}$  represent the distributed channel resistance corresponding to capacitances  $C_{gs}$  and  $C_{gd}$ , respectively. NQS signifies the inertia of the channel to a change in gate-source voltage. This phenomenon occurs at frequencies beyond  $f_T$ . The NQS effect has been represented by a transit delay  $\tau$  in the literature. The delay appears both on the transconductance and channel conductance. The first-order representation of  $g_{m,nqs}$  and  $g_{ds,nqs}$  are given by [16],

$$g_{m,nqs} = \frac{g_m}{1+j\omega\tau},\tag{4}$$

$$g_{ds,nqs} = \frac{g_d s}{1 + j\omega\tau}.$$
(5)

The intrinsic *Y*-parameters in the strong inversion are given by the following equations [13]:

$$\begin{split} Y_{11,i} &= \omega^2 (R_{gd} C_{gd}^2) + j\omega (C_{gd} + C_{gs}), \\ Y_{12,i} &= -\omega^2 R_{gd} C_{gd}^2 - j\omega C_{gd}, \\ Y_{21,i} &= g_m - \omega^2 R_{gd} C_{gd}^2 - j\omega (C_{gd} + g_m \tau), \\ Y_{22,i} &= g_{ds} + \omega^2 R_{gd} C_{gd}^2 + j\omega (C_{gd} - \tau g_{ds}). \end{split}$$
(6)

The intrinsic capacitances, distributed resistances and transit delay are extracted from the above equations. The



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FIGURE 5. Extracted (a)  $r_{gs}$  and (b)  $\tau$  over frequency.

resistances  $r_{gs}$  and  $r_{gd}$  are extracted for different gate-source voltages while keeping the drain-source voltage zero. It is observed that applying this method results in negative  $r_{gd}$  for all bias points, which means that  $r_{gd}$  is practically zero. The maximum  $r_{gs} \approx 650 \,\Omega$  is obtained for  $V_{GS} = 0.4 \,\text{V}$ , as shown in Fig. 5a.

The channel delay  $\tau$  obtained from  $Y_{21}$  is shown in Fig. 5b for different bias points. The channel delay decreases as  $V_{DS}$  increases. The lowering of channel delay between the source and drain signifies an increased carrier speed with an increasing electric field between the source and drain.

### D. IMPLEMENTATION OF THE MODEL

The extrinsic components are implemented as constants in the Verilog-A model of the TFT.

The voltage dependency of  $C_{GD}$  and  $C_{GS}$  are implemented using the Meyer capacitance distribution of the total oxide capacitance of 90 fF. For a-IGZO TFTs, the validity of the Meyer model has already been proved for intrinsic capacitance in literature [17] and [18]. The total capacitance and the transconductance  $g_m$  determine  $f_T$  of the TFT.

The derived DC RPI-a model parameters control the low-frequency values of  $g_m$  and  $g_{ds}$ .

The NQS effect, i.e. the effect of  $\tau$  on  $g_m$ , is modeled in the charge calculation with the help of the "laplace\_nd" function in VerilogA. In the RPI-a model, there is no simple equation for  $g_m$ . So  $\tau$  is introduced in the channel charge concentration  $(n_s)$  of the RPI-a model.

$$n_s = \frac{\epsilon_i \epsilon_0}{q t_{ox} \text{VAA}^{\gamma}} \frac{(V_{gs} - V_{th})^{\gamma+1}}{1 + \omega \tau}$$
(7)

The same function is used for  $g_{ds}$ , which can be given by

$$g_{ds} = \frac{g_{chi}}{1 + g_{chi} \frac{(R_D + R_S)}{N_{\text{finger}}}} \frac{1}{1 + \omega \tau},$$
(8)

where  $g_{chi}$  is the intrinsic channel conductance of the RPIa model and  $N_{\text{finger}}$  is the number of fingers in a TFT [5]. After repeated fitting, a constant value of the channel delay  $\tau = 150$  ps is taken for our modeling purpose.



(a)  $V_{DS} = 0 \text{ V}, V_{GS} = -2 \text{ V}$  (b)  $V_{DS} = 2 \text{ V}, V_{GS} = 2 \text{ V}$ 

FIGURE 6. Simulated versus measured Y-parameters (solid lines simulation and circular markers - measurement) of a single-finger TFT for different bias conditions. Legend (11,12,21 and 22) represents the subscripts of the Y-parameters.

### E. SINGLE-FINGER TFT VERIFICATION OF THE MODEL

For the single-finger TFT, we get an almost perfect fit for all *Y*-parameters in the null-bias condition, as seen in Fig. 6a. When bias voltages are applied, as shown in Fig. 6b,  $re[Y_{21}]$  and  $re[Y_{22}]$  deviate from the measured values, especially as the frequency increases. Since the transconductance and the channel conductance are derived from the DC model, the values at higher frequencies are not estimated accurately. However, we get a good fit for  $Y_{21}$  up to around 1.0 GHz; the simulated  $re[Y_{21}]$  become higher than measured data for  $RF \geq 2$  GHz. The imaginary parts of *Y*-parameters give a good fit up to around 1.0 GHz, which also indicates the validity of the Meyer capacitance model.

### F. MULTI-FINGER TFT VERIFICATION OF THE MODEL

To verify the derived model's scalability, a ten-finger  $(N_{\text{finger}} = 10)$  was simulated for DC behavior using Cadence Spectre. It has been observed that  $V_{th}$  decreases as the number of fingers increases. For the measured multi-finger device, at  $V_{DS} = 0.4$  V, a threshold voltage  $V_{th} = -1.38$  V has been extracted using the integration method presented in [7].

The electric and magnetic field couplings between the fingers are ignored during the modeling. Due to a high separation of 20 µm between the adjacent fingers, their isolation is very high. The high isolation is verified by simulating the gate multi-finger architecture in Ansys HFSS, shown in Fig. 7a. The pad in the gate electrode is removed to simulate only the reactive coupling between the fingers. The reflection coefficient for the center finger (suffixed 'C') and its isolations with the adjacent left (suffixed 'L') and right (suffixed 'R') fingers are shown in Fig. 7b. The couplings are lower than -40 dB for frequencies up to 6 GHz. Hence the coupling can be ignored.

The simulated versus measured output and transfer characteristics are plotted in Fig. 8a and 8b, respectively. As seen in Fig. 8a, the simulated current (solid lines) matches the measurements (circular markers) better at high  $V_{GS}$ , whereas from Fig. 8b, the simulated transfer curve matches the measurement better at low  $V_{DS}$ .





**FIGURE 7.** Reactive coupling simulation (a) setup and (b) results for the ten-finger TFT. The suffixes 'C', 'L' and 'R' indicate central, left and right fingers, respectively.



FIGURE 8. Simulated versus measured (a) output and (b) transfer characteristics (solid lines - simulation and circular markers - measurement) for a ten-finger TFT.

The *Y*-parameters are also compared for the ten-finger TFT. The extrinsic drain-source resistance is kept at  $R_{DSE} = 10 \Omega$  since it gives the best fitting. The simulated *Y*-parameters agree with the measured data in the null-bias condition, as shown in Fig. 9a.

The simulated *Y*-parameters under biasing show higher deviations, as depicted in Fig. 9b. The simulated  $re[Y_{22}]$  deviates over a large frequency range.

### G. TRANSISTOR SPEED AND AMPLIFICATION

The measured versus simulated short circuit current gain  $H_{21}$ and the maximum available power gain  $G_{max}$  are extracted from the *S*-parameter. The frequencies at which  $H_{21}$  and  $G_{max}$ reach unity are  $f_T$  and  $f_{max}$ , respectively.

The relation between  $f_{max}$  and  $f_T$  is given by [4]

$$f_{max} = \frac{1}{2} \frac{f_T}{\sqrt{2\pi f_T C_{gd} R_G + g_{ds} R_G}},$$
(9)

where  $f_T = g_m / (2\pi (C_{gs} + C_{gd}))$ .

Fig. 10a and 10b show  $f_T$  and  $f_{max}$  for the single and tenfinger TFTs, respectively. The bias points in both figures correspond to equal gate overdrive voltage and equal drain voltage for both TFTs.

The top halves of both the figures show  $f_T$ . The measured  $f_T$  values for the single and ten-finger TFTs are identical. Increasing the number of fingers scales the transconductance



(a) 
$$V_{DS} = 0 \text{ V}, V_{GS} = -2 \text{ V}$$
 (b)  $V_{DS} = 2 \text{ V}, V_{GS} = 1.6 \text{ V}$ 

**FIGURE 9.** Simulated versus measured *Y*-parameters (solid lines simulation and circular markers - measurement) of a multi-finger TFT for different bias conditions. Legend (11,12,21 and 22) represents the subscripts of the *Y*-parameters.



**FIGURE 10.** Simulated (solid lines) versus measured (circular markers)  $f_T$  and  $f_{max}$  of (a) single-finger and (b) ten-finger TFTs.

and the gate capacitances proportionately. Hence,  $f_T$  remains independent of the number of fingers.

There is a good agreement between the measured and simulated  $H_{21}$  for the single-finger and ten-finger TFTs (< 5% error) at higher gate voltages. We can conclude that the Meyer capacitance model approximates the intrinsic oxide capacitances with a low error rate for higher gate voltage. However, at a low gate voltage of  $V_G = 0.4$  V, the simulated  $f_T$  is underestimated. The higher measured  $f_T$  value could be due to additional traps due to impurities or uneven electric field distribution across the oxide that reduces the gate capacitance in reality.

The bottom halves of the figures depict  $f_{max}$  of the TFTs. There is a decrease in the multi-finger  $f_{max}$  compared with the single-finger TFT. This reduction in the oscillation frequency is due to the additional gate electrode material from the probe tip to the channel (see Fig. 2b), which increases the gate resistance.

For the bias points where  $f_T$ s are in good agreement, we achieved a good agreement between the simulated and the measured  $f_{max}$  (< 20% error), thereby proving a reasonable estimation of the gate resistances.

# H. TRANSFER FUNCTION

Comparisons of the small signal transfer function between the simulation and measurement are shown in Fig. 11a and 11b for  $V_{DS}$  of 0.4 V and 2 V, respectively.



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FIGURE 11. Transfer function  $S_{21}$  of the TFT for (a)  $V_{DS} = 0.4$  V and (b)  $V_{DS} = 2$  V.

TABLE 3. Root mean square error in the simulated transfer function.

$V_{GS}\left(\mathbf{V}\right)$	-0.8	0.4	1.6	2
region	Weak	Moderate	Strong	Strong
$RMSE V_{DD} = 0.4 V$	0.08	0.067	0.086	0.088
$RMSE V_{DD} = 2.0 V$	0.129	0.264	0.360	0.365

Table 3 presents the root mean square error between the measured and simulated  $S_{21}$ . As seen here, the weak inversion presents a lower error than the strong inversion region. Moreover, the error is higher at higher  $V_{DS}$ .

The reason for this error is the fitted value of  $\tau$ , which is higher than the value extracted for high  $V_{DS}$ , as seen in Fig. 5b. The RMS error increases above 1 GHz. Thus the frequency limit of the model for different bias points is around 1 GHz, which is the frequency of operation.

# IV. THEORY OF a-IGZO TFT BASED FREQUENCY DOUBLERS

The schematic of the common-gate (CG) differential frequency doubler is shown in Fig. 12a. The device can be simulated and explained with the extracted behavioral model. The CG architecture helps avoid Miller capacitance. The differential architecture provides AC ground at the gate and the drain terminals when the doubler is symmetric [1]. If the two TFTs are symmetric, it helps eliminate the fundamental harmonic signal from the output.

### A. NUMBER OF FINGERS VERSUS OUTPUT

The extracted model simulates the doubler for output power and conversion gain. A drain voltage of  $V_{DS} = 2$  V and an input power  $P_{in}$  of 12.3 dBm are taken. The nominal  $V_{th}$  of the TFT is -0.95 V, as extracted from the single-finger TFT data (see Table 1).

First, the simulated  $P_{2f}$  is plotted over  $V_{GS}$  for different numbers of fingers  $N_{\text{fingers}}$  in Fig. 13a. The variation of peak  $P_{2f}$  from the graph over  $N_{\text{fingers}}$  is shown in Fig. 13b. As  $N_{\text{fingers}}$  increases, the peak  $P_{2f}$  also increases.

However, the fabrication-yield scales inversely with the number of gate fingers. Moreover, a higher  $N_{\text{fingers}}$  also increases the process mismatch. Hence, we chose a differential doubler with four fingers in each TFT.



FIGURE 12. (a) Schematic and (b) photo (background edited) of the top-view of a common-gate frequency doubler with dimensions in  $\mu m$ .



**FIGURE 13.** Simulated (a)  $P_{2f}$  over  $V_{GS}$  for different fingers, and (b) peak achievable  $P_{2f}$  over the number of fingers, given  $V_{DS} = 2 \text{ V}$ ,  $V_{th} = -0.95 \text{ V}$  and  $P_{in} = 12.3 \text{ dBm}$ .

### V. RESULTS, SIMULATIONS AND DISCUSSIONS

The photo of the fabricated CG doubler is shown in Fig. 12b. The TFTs are susceptible to breakdown. Hence  $V_{DS}$  is limited up to 2 V during characterization.

The measurement setup for the doubler consists of an Agilent frequency synthesizer followed by an HP 3304C 20 dB attenuator and a Mini-Circuits ZHL 3010+ power amplifier. The overall gain provided by the attenuator and amplifier chain is around 16 dB. The output from the amplifier is then fed into a Mini-Circuits ZFSCJ-2-4-S+ 180° hybrid. The differential output from the hybrid is then applied to a Picoprobe 40A-GSSG-150-P ground-signal-signal-ground probe in combination with a Mini-Circuits bias tee ZX85-12G-S+. On the output of the doubler, a Mini-Circuits bias tee ZFBT-6G+ reads the output power and provides the drain bias.

The input power is calculated after de-embedding the attenuator and power amplifier chain gain of 16 dB, the hybrid loss of 1.1 dB [19], the GSSG probe loss of 1.6 dB [20] and the total connector loss of 0.5 dB from the available power of -1 dBm of the spectrum analyzer. Moreover, the common ground paths of the doubler, composed of Molybdenum, also attenuate the input power. The resistivity of the 90 nm molybdenum source terminal is taken to be around 5.6 ×



**FIGURE 14.** Initially measured (a) output and (b) transfer I-V characteristics of the doubler.



**FIGURE 15.** Measured versus simulated (a)  $P_f$  for and (b)  $P_{2f}$  for  $V_{th} = -1.475$  V and  $\Delta V_{th} = 0$  V.

 $10^{-6} \Omega$ cm [21]. Hence, an additional source path resistance  $R_{S,add}$  of around 15  $\Omega$  is added to the inputs of each TFT.

### A. DC SIMULATION

As mentioned in subsection III-F, the TFTs'  $V_{th}$  varies as the number of fingers increases. The initial measured DC characteristics of the doubler are shown in Fig. 14a and 14b with the circular markers. The best fit for the simulated current (solid lines) is achieved by fitting  $V_{th}$  to -1.475 V. A perfect estimation of  $V_{th}$  is impossible due to device asymmetry and process variation among the fingers.

### **B. GATE VOLTAGE VERSUS OUTPUT POWER**

Since the breakdown voltage of the TFTs is unknown, a very low supply voltage  $V_{DD} = 500 \text{ mV}$  is taken first for characterization. At an input power of around 12 dBm, the gate voltage is swept from -4 V to 0.25 V. Initially, both the first harmonic power  $P_f$  and second harmonic power  $P_{2f}$ increase as  $V_{GS}$  increases, as shown in Fig. 15a and Fig. 15b, respectively. The maximum values are obtained at around  $V_{GS} = -0.725 \text{ V}$ , then both output powers decrease.

The simulated  $P_{2f}$  does not match the measurement in terms of the  $V_{GS}$  corresponding to the peak. The simulated peak  $P_{2f}$  is observed at  $V_{GS} = 0.5$  V, i.e. at a value of 1.5 V higher than the measured peak location. Moreover, for the same 12.3 dBm of input power as in the measurement, the



**FIGURE 16.** (a) Measured  $P_f$  and  $P_{2f}$  over  $V_{S1}$  and (b) measured versus simulated  $I_D$  during RF measurement.

peak simulated  $P_{2f}$  is -18 dBm. This value is 2 dB higher than the measured  $P_{2f}$ .

### C. SPECTRAL IMPURITY

The doubler also shows a high first harmonic output  $P_f$ . The spectral impurity in doublers occurs from device asymmetry. Hence, the source voltage of one of the TFTs is swept to find  $V_{th}$ -mismatch while keeping the gate voltage  $V_G = -725 \text{ mV}$  and drain voltage  $V_D = 0.5 \text{ V}$ . The lowest  $P_f$  is obtained at  $V_{S1} \approx -125 \text{ mV}$  (see Fig. 16a). The change in  $\Delta V_{th}$  does not affect  $P_{2f}$ .

### **D. BIAS STRESS**

As applied bias and  $P_{in}$  increase, the TFT  $V_{th}$  decreases. We can confirm this shift in  $V_{th}$  from the measurements and fitting of the model. The drain current is measured simultaneously as a  $V_{GS}$  sweep for RF is performed. Fig. 16b shows the measured  $I_D$  with circular markers. By keeping the same  $V_{th} = -1.475$  V (from Fig. 14), we cannot achieve the same simulated current values for the same  $V_{GS}$ . Hence, we sweep  $V_{th}$  and the fitting parameter VAA in Cadence Spectre to arrive at the same current. The best fitting is achieved at values  $V_{th1} = -3.525$  V,  $V_{th2} = -3.65$  V and VAA = 4.0. This change implies that  $V_{th}$  has moved to a more negative value under the influence of the RF signal. The adjustment in VAA indicates a change in mobility as well.

The negative shift in  $V_{th}$  is assumed to be caused by the long exposure of the TFT to bias and heat generated by the RF input, as suggested in [22]. Moreover, the presence of water molecules and hydrogen atoms in the oxide can also cause a negative  $V_{th}$  shift [22], [23].

The RF output behavior is also simulated with the final chosen values  $\gamma = 1.014$ ,  $V_{th1} = -3.525$  V,  $V_{th2} = -3.65$  V and VAA = 4.0. The input power is kept at 12.3 dBm. As shown in Fig. 17a, the simulated fundamental harmonic output power has a lower deviation ( $\approx$  3 dB) from the measured values.

The simulated second harmonic output  $P_{2f}$  agrees better with the measured result, as shown in Fig. 17b. One must note that an accurate simulation result is impossible to obtain as there are four fingers in each transistor. There may be



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FIGURE 17. Measured (circular markers) versus simulated (solid lines) (a)  $P_f$  and (b)  $P_{2f}$  for  $\gamma = 1.014$ ,  $V_{th1} = -3.525$  V,  $V_{th2} = -3.65$  V and VAA = 4.0 at  $P_{in} = 12.3$  dBm.



**FIGURE 18.** Measured versus simulated (a)  $P_f$  for and (b)  $P_{2f}$  for different values of the overlap capacitance per finger  $C_{OV1}$  of TFT 1 while keeping  $C_{OV2}$  of TFT 2 constant.

mismatches in the process parameters among all the fingers. However, the peak simulated  $P_{2f}$  is the same as the measured peak.

#### E. MISMATCH IN OVERLAP

The multi-finger layout is susceptible to alignment issues since low-resolution PET foil is used for UV exposure. This misalignment could cause a mismatch between the overlap lengths between the two TFTs. As a result, the overlap capacitances between the two TFTs can be different. The mismatch can also result in higher spectral impurity in addition to the threshold and mobility mismatch. The effect of overlap capacitance  $C_{OV}$  per finger mismatch on  $P_f$  and  $P_{2f}$ are shown in Fig. 18a and 18b, respectively. The TFT with the lower  $V_{th1} = -3.525$  V and higher  $C_{OV1} = 63$  fF per finger shows a better fit between the simulated and the measured measured  $P_f$ . The change in the overlap capacitance does not affect  $P_{2f}$ .

However, the effect of  $C_{OV}$  mismatch can only be speculated. The mismatch in intrinsic capacitance could also lead to a similar effect in  $P_f$  and spectral purity.

#### F. RF VERSUS BREAKDOWN

One of the most important characteristics of a frequency doubler is compression. It is the condition at which  $P_{2f}$ 



**FIGURE 19.** (a) Measured (circular markers) versus simulated (solid lines)  $P_{out}$  over  $P_{in}$  with the breakdown (inset - doubler after breakdown) (b) simulated  $V_{ds}$  and  $V_{gs}$  for  $\gamma = 1.014$ ,  $V_{th1} = -3.525$  V,  $V_{th2} = -3.65$  V and VAA = 4.0 at  $P_{in} = 10$  dBm; the inset shows a photo of the channels' breakdown.

becomes saturated. The doubler output power normally increases when  $V_{DS}$  increases. Here the input power is swept for a  $V_{DS} = 2$  V,  $V_{GS} = -0.775$  V, and  $V_{S1} = -125$  mV to compensate for a  $V_{th1} - V_{th2} = -125$  mV. The simulation is also done with the same biasing conditions. As shown in Fig. 19, the nominal simulation using the parameters from Table 1 agrees with the measured value in terms of the slope. The simulated  $P_{2f|nominal}$  is underestimated by around 1 dB.

However, the measured  $P_{f|nominal}$  is underestimated in the simulation by around 3 dB. The simulation shows that  $P_f$  can be matched to the measured data if the total gate capacitance  $C_{GG1}$  of TFT 1 is lowered by 3 fF. The same effect is also observed if the power law mobility parameter of the TFT 1 can be increased from its nominal value (1.0144) to 1.1. Further investigation is required to pinpoint the exact source of the spectral impurity.

The measured power starts reducing at  $P_{in} = 10$  dBm. The inset in Fig. 19b shows that the device breaks down beyond this point. Not all the fingers break down simultaneously, indicating that the heat generation is not uniform across the fingers. The simulated  $V_{ds}$  and  $I_d$  over time of the two TFTs at the input power of  $P_{in} = 10$  dBm. The root-meansquare (rms)  $V_{DS}$  is only around 0.84 V at this point, while the rms drain current per TFT is 1.35 mA. Ideally, a 47 nm thick Al<sub>2</sub>O<sub>3</sub> dielectric layer should easily sustain this value. We assume that the presence of impurities such as water or hydrogen in the oxide layer reduced the breakdown voltage of the dielectric.

### G. COMPARISON WITH THE STATE OF THE ART

The application of a-IGZO in RF is a relatively new field; hence, there are only a few publications. The presented frequency doubler is a record in terms of the frequency of operation. The previous work in [1] showed frequency doubling for an input frequency of 100 MHz. Table 4 shows some TFT technology and its performance in a frequency doubler.

#### TABLE 4. Comparison of doubler with the state of the art.

Ref	W/L (µm/µm)	RF (MHz)	Gain (dB)
a-IGZO TFT	500/3	$100 \times 2$	-44
[1]			
Graphene Am-	4/6	$11 \times 2$	-17.7
bipolar [24]			
This work	$4 \times 100/0.8$	$500 \times 2$	-32

# **VI. LIMITATION OF THE MODEL**

The model adopted suffers from the following limitations, and further investigations are required:

- The channel transit delay  $\tau$  is considered constant in the model, while this is a function of the drainsource voltage, as described in subsection III-C. The effect of velocity saturation is not included in the model.
- The simulated transfer function  $S_{21}$  deviates substantially from measured values as  $V_{DD}$  and RF increase.
- The modeling of the capacitor is done with Meyer's model, which is ideal for frequency below  $f_T$  [25]. Applying an input signal beyond  $f_T$  will cause the transistor capacitances to charge and discharge slowly compared to the input signal. The dispersive behavior of capacitance is not considered in the model.
- The error in the output admittance is high, as shown in Fig. 6b and 9b. The drain-source capacitance is included in the pad capacitance and is considered constant. Any effect of the drain-induced barrier lowering is not included.
- Temperature dependencies of model parameters are not considered.
- Hot carrier effect and trapping and de-trapping of charge carriers in the oxide cannot be simulated by the model.

### **VII. CONCLUSION**

This paper presents the efficacy of the RPI-a model in conjunction with small-signal modeling for the simulation of GHz TFT circuits based on a-IGZO. We verified the small-signal model parameters for a single-finger and a tenfinger TFT. An excellent agreement is achieved between measurement and simulation for the obtained DC and RF parameters. The agreement between most *Y*-parameters makes the model suitable for basic RF circuit simulation. However, the DC output conductance is underestimated at GHz frequencies for higher bias points.

A four-finger TFT-based doubler is characterized and also evaluated with the derived model. The doubler shows a conversion gain of -32 dB for an input signal of 500 MHz without impedance matching. It is the first doubler in literature with GHz output frequency. Owing to the non-linear relation between  $V_{gs}$  and  $I_d$ , we could achieve the frequency doubling beyond transistor  $f_T$  and  $f_{max}$ . The second harmonic frequency is a new record in terms of conversion gain of a-IGZO doubler. Our previous work presented a 100 MHz to 200 MHz doubler with only a -44 dB conversion gain [1]. The application of impedance matching can further improve the conversion gain in the future.

Applying the RF input power, a significant  $V_{th}$  shift is observed. This shift is hypothesized to be due to impurities in the dielectric. Since the TFT is handled in a laboratory atmosphere, the dielectric-semiconductor and the semiconductorencapsulation interfaces have water molecules that reduce the lifetime of the transistor.

We can feed the shift in the threshold voltage to the model to replicate its effect on the second harmonic output. The model can also simulate the spectral impurity of the doubler based on the threshold voltage mismatch.

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