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### **RESEARCH ARTICLE**

# Metaheuristic-Based Optimization and Prototype Investigation of Low Frequency Metamaterial for Wireless Power Transfer Application

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**ABSTRACT** The adoption of wireless charging technologies for consumer and industrial applications is inhibited by concerns over reliability culminating from the fluctuating output power, low transmit power, and power transfer efficiency (PTE) of these systems. Besides, the inherently high radio and microwave resonant frequency of existing Metamaterial (MTM) designs imposes high switching stress on power switches and passive components, leading to significant power loss and enormously high degradation in system performance. This manuscript presents a compact, low-frequency MTM-based Wireless Power Transfer (WPT) structure coupled with a comprehensive investigation of the effect of physical parameters on the resonant frequency where left-handed characteristics occur. Model design and simulation were conducted in ANSYS High-Frequency Structure Simulator (HFSS) to extract the transmission coefficient and reflection coefficient while realizing negative permeability at a resonant frequency ( $f_o$ ) of 1.2MHz. To further mitigate the resonant frequency of the MTM, a metaheuristic-based parameter optimization algorithm was implemented to achieve negative permeability and evanescent wave amplification at a resonant frequency of 750kHz, making it suitable for high-power WPT applications. A prototype MTM sample is fabricated for experimental measurement of power transfer efficiency and medium parameters, using the Keysight ENA5061 Vector Network Analyzer (VNA), effectively confirming the validity of the proposed design. The excellent efficiency enhancement and mutual coupling make the design an attractive solution for WPT applications. A close agreement between the experimental results and numerical simulation validates the accuracy of the optimization results.

**INDEX TERMS** Wireless power transfer, finite element analysis (FEA), metamaterial (MTM), power transfer efficiency, ANSYS, high frequency structure simulator, vector network analyzer, optimization.

#### I. INTRODUCTION

Wireless Power Transfer (WPT) has been widely studied and implemented in various applications, including mobile computing [1], wireless charging of biomedical body implants [2],

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and Electric Vehicles (EV) [3]. This technology was first discovered by Nikola Tesla in 1891 [4], and since then, significant research progress has been made. In 2007, researchers at MIT demonstrated the potential for efficient long-distance wireless power transmission using magnetic resonant coupling [5], and recently, Oak Ridge National Laboratory designed a 120-kilowatt WPT prototype for the wireless



**FIGURE 1.** Schematic showing the categorization of Metamaterial (MTM) contingent on effective permeability( $\mu_f$ ) and effective permittivity ( $\epsilon_f$ ) polarity. The electric field, magnetic field, wave number, and phase velocity, are denoted as *E*, *B*, *k*, *V*<sub>ph</sub>, respectively.

charging of EV. However, there are still challenges to be addressed to fully realize the potential of WPT, including limited efficient magnetic resonant coupling distance and the inherent dependence of power transfer efficiency (PTE) on the coupling coefficient (k). Metamaterials (MTM) have emerged as promising candidates for enhancing WPT performance over a wide transfer distance. First described by Veselago in 1968 [6], MTM represents a generic term for materials with negative refractive index and perfect lensing (negative permittivity and permeability ( $\epsilon_f < 0$ ;  $\mu_f < 0$ )) as exemplified in Eq. (1)

$$n = -\sqrt{\mu_f \epsilon_f} \tag{1}$$

The broad categorization of MTMs is shown in Fig. 1. While negative (DNG) MTM demonstrates negative effective permeability ( $\mu_f < 0$ ) and negative effective permittivity ( $\epsilon_f < 0$ ) for backward wave propagation, double positive (DPS) materials (forward wave media) are characterized by  $\mu_f > 0$  and  $\epsilon_f > 0$ . However, in extremely sub-wavelength regions, there is a separation between the electric and magnetic fields, allowing the amplification of the evanescent wave and the attainment of a negative refractive index (n < 0).

To achieve practical low-frequency MTM structures with valid medium parameters, the MTM unit cell must satisfy the homogenization criterion, which specifies a small sample size ( $l \ll \lambda$ ) compared to the wavelength at the operating frequency [7], [8], [9], [10], [11], [12], [13]. Numerous MTM designs and geometries have been reported in the literature to improve the performance of wireless power transfer (WPT). For example, in [14], split ring/spiral



FIGURE 2. Schematic summary of design implementation process.

printed resonators and rectangular spiral-based MTM structures are discussed in detail, while [15] utilizes a multi-layer array of rectangular spiral for maximizing power transfer coupled with the enhancement of the power transmission efficiency (PTE). In [16], an effective MTM design for enhancing WPT performance is analyzed. Furthermore, [17] explores a planar hybrid MTM structure that integrates zeropermeability ( $\mu_f = 0$ ) and negative permeability ( $\mu_f < 0$ ) MTMs to reduce leakage EMF and increase transmit power. Moreover, [18], [19], [20], [21] comprehensively investigates various low-frequency MTMs that rely on the tuning of physical parameters. Although these MTM-based WPT designs exhibit superior performance in comparison to conventional coil structures, nonetheless, they typically resonate at significantly high radio and microwave frequencies, which is unfavorable for high-efficiency WPT applications [22], [23], [24]. In general, high resonant frequency results in high voltage and current stresses on passive components while also subjecting power semiconductor/MOSFET to high switching stress, leading to increased power dissipation, low efficiency, and degradation in WPT performance [10], [22], [23], [24], [25], [25], [26], [27].

In this study, a low-frequency thin-PCB Metamaterial structure is studied alongside an optimization algorithm to analyze the impact of physical parameters on the resonant frequency in which evanescent wave amplification occurs. The proposed MTM structure stands out due to the unique way its individual layers are connected to achieve a constructive addition of magnetic flux. The main objective of this research is to investigate a left-handed MTM and decrease its evanescent resonant frequency to the kHz range, while simultaneously enhancing the PTE. A schematic representation of the model implementation steps is provided in Fig. 2.

The main contributions of this paper are summarized as follows:

- i To design a miniaturized, thin PCB metamaterial (MTM) that operates at low frequencies for efficient wireless power transfer.
- ii To develop an optimization algorithm to reduce the resonant frequency of the proposed MTM.
- iii To experimentally validate the accuracy of the optimization algorithm using a prototype of the fabricated MTM sample

#### II. ANSYS HFSS ELECTROMAGNETIC SIMULATION AND METAMATERIAL CHARACTERIZATION

#### A. SYSTEM DESCRIPTION AND ELECTROMAGNETIC SIMULATION

The proposed MTM design is depicted in Fig. 3(a). Its layered rectangular spirals are orientated in opposite directions (The

TABLE 1. Parameter Specification Proposed MTM Unit cell.

| Parameter                       | Unit            | Value  |
|---------------------------------|-----------------|--|
| Inter-turn spacing $(\Delta r)$ | mm              | 1.6  |
| Wire width $(t)$                | mm              | 1.9  |
| Litz wire guage                 | -               | 26 AWG, Unserved Single Build,<br>16 / 38 strand |
| Length $(l)$                    | mm              | 64   |
| Dimension                       | $\mathrm{mm}^3$ | $64 \times 32 \times 3$                          |



(a) Schematic of the Proposed MTM. The middle layer is a mirror image of both the first and third layer. The individual layers are connected in series to form a unit cell. Red arrows typify the direction of magnetic flux



(b) Fundamental Component of the proposed metamaterial



middle layer is a mirror image of the first and third layers) and also connected in series to achieve a constructive addition of the magnetic field. Apparently, the red arrow indicates the direction of travel of the magnetic flux. The series connection of the three layers constitutes an MTM unit cell. Similarly, Figure. 3(b) shows the fundamental layer of the proposed MTM in Fig. 3(a). Typically, each layer of the MTM is characterized by a wire thickness, *t*, an inter-turn spacing,  $\Delta r$ , and a length, *l*, which represents the maximum length of the unit cell. In addition, Table 1 denotes the parameter specification of the impact of design parameters on the left-handed characteristics and resonant frequency of the proposed MTM structure is carried out using the finite element solver in ANSYS HFSS.



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FIGURE 4. Configuration of unit metamaterial in a Full-wave ANSYS HFSS simulation. The setup includes the assignment of boundary conditions and excitation ports.



**FIGURE 5.** Extracted Transmission coefficient ( $S_{21}$ ) of the proposed metamaterial, indicating evanescent wave amplification at  $f_o = 1.2$  MHz.

As shown in Figure. 4, the simulated unit cell is enclosed in a Teflon substrate, having a high dielectric permittivity ( $\epsilon_r = 4.4$ ) and low loss tangent ( $tan\delta = 0.0004$ ).

In addition, perfect electric ( $E_t = 0$ ) and perfect magnetic  $(H_t = 0)$  boundary conditions were applied to terminate the MTM along the y-axes and z-axes, respectively. Moreover, excitation is assigned to the boundary along the x-axis using externally applied waveguide ports. The excitation creates a magnetic field around the MTM unit cell. Using an extraction distance equal to the length, l, of the MTM coupled with the simulated transmission coefficient  $(S_{21})$ , and reflection coefficient  $(S_{11})$  holistically referred to as scattering parameters, both the effective permeability  $(\mu_f)$  and effective permittivity ( $\epsilon_f$ ) were retrieved based on Kramar-Kronig parameter retrieval algorithm [28]. The de-embedded simulated scattering parameter,  $S_{21}$  and  $S_{11}$  are shown in Fig. 5 and Fig. 6, respectively. Ostensibly, the proposed MTM exhibits a resonance frequency,  $F_r = 1.2$  MHz, a significantly low frequency compared to the Gigahertz frequency of existing MTM structures. Additionally, the retrieved  $\mu_f$  and  $\epsilon_f$ plots depicted in Fig. 7 and Fig. 8, respectively demonstrate



**FIGURE 6.** Extracted metamaterial reflection coefficient, showing low-frequency evanescent wave amplification. The metamaterial resonates at a frequency,  $f_0 = 1.2$  MHz.



**FIGURE 7.** Retrieved permeability of proposed metamaterial, showing left-handed characteristics at a resonant frequency,  $f_0 = 1.2$  MHz.

1.2 MHz resonant frequency consistent with Fig. 5 and Fig. 6. Thus, the resonance frequency of  $\mu_f$ ,  $\epsilon_f$ ,  $S_{21}$ , and  $S_{22}$  are closely matched. In addition, the MTM demonstrates a negative real component of  $\mu_f$ , corresponding to -2.8 at the resonant operating point, effectively confirming the validity of the MTM. Moreover, at the resonant frequency, the ratio of the electromagnetic wavelength and the length of the MTM unit cell ( $\lambda_o/l$ ), which typifies a metric for appraising the compactness of MTM structures is evaluated as 3906, making it the most compact unit cell. A comparison of the proposed MTM and other existing MTM structures is shown in Table 2. Apparently, the proposed MTM demonstrates a very compact and miniaturized structure.

### B. ELECTROMAGNETIC SIMULATION AND FINITE ELEMENT ANALYSIS

In order to investigate the performance enhancing capability of the proposed MTM, an MTM-slab based on a



**FIGURE 8.** Retrieved permittivity of the proposed metamaterial, exhibiting left-handed characteristics at a resonant frequency,  $f_o = 1.2$  MHz.

**TABLE 2.** Comparison of Wavelength to Length Ratio  $(\frac{\lambda_0}{L})$  betweenexisting Works and the Proposed Metamaterial.

| Reference                                       | [14]     | [29]     | [30]     | Proposed<br>MTM |
|---|----------|----------|----------|-----------------|
| Lateral size $\left(\frac{\lambda_o}{l}\right)$ | 170      | 300      | 442.8    | 3906            |
| Resonant frequency                              | 8.69 MHz | 8.42 MHz | 7.49 MHz | 745.5 kHz       |

 $4 \times 4$  periodic array of unit MTM is integrated with a conventional two-coil WPT structure as displayed in Figure 9. Both the transmitter (Tx), and receiver (Rx) are made of three turns, 24AWG solid strand copper wire, having a coil radius of 100 mm; inductance, 10 nH; loaded chip capacitance; 100 pF, and an ohmic resistance of 0.0025  $\Omega$ . The aforementioned simulation parameters are in agreement with 5% maximum harmonic content for ripple minimization based on IEEE standard recommendation [22]. For performance comparison purposes, two full wave simulations were conducted in ANSYS electromagnetic simulator, first for the classical WPT system without MTM followed by the proposed MTM-incorporated WPT structure.

To preserve the working distance as well as prevent the intersection of the structures, a distance of 0.5 mm is maintained between transmitter (Tx) and MTM as well as receiver (Rx) and MTM while varying the working distance (MTM to MTM separation) from 50 mm to 250 mm. Moreover, it was determined based on ANSYS electromagnetic simulation that a 0.5 mm minimum distance is critical to preventing the intersection of the Tx and MTM as well as Tx and MTM. Besides, it is worth noting that attaching the MTM to the surface of either Tx or Rx effectively places the MTM on the flux traveling path, potentially enhancing leakage magnetic field and reduction in power transfer efficiency. Lastly, the power transfer efficiency of both the WPT system with and without metamaterial is analyzed for performance comparison.



(a) With Metamaterial

(b) Without Metamaterial

**FIGURE 9.** ANSYS HFSS 3D simulation schematic of the proposed MTM-based WPT system. A separation of 0.5 mm is maintained between Tx - MTM and Rx - MTM to prevent a potential intersection. Each slab comprises a  $4 \times 4$  array of unit MTM with a dimension of  $256 \times 96 \times 3 \text{ mm}^3$ .



(a) Schematic of the proposed MTM unit cell based on ANSYS HFSS



(b) Equivalent RLC lumped circuit representation of the proposed MTM

**FIGURE 10.** Frontal Section of proposed MTM unit cell as designed in ANSYS HFSS and equivalent  $R_m L_m C_m$  lump circuit.

#### III. PHYSICS-BASED MODELING AND METAHEURISTIC OPTIMIZATION OF PROPOSED METAMATERIAL FOR RESONANT FREQUENCY REDUCTION

This research combines physics-based modeling with metaheuristic-based optimization to analyze how the physical parameters of a proposed metamaterial (MTM) structure affect its resonant frequency. The main objective is to reduce the resonant frequency where evanescent wave amplification occurs. To achieve this goal, the study utilizes a Jaya-based optimization algorithm to optimize the lumped capacitance and stray inductance of the MTM, which have a comprehensive impact on the resonant frequency of the structure.

#### A. PHYSICS-BASED EXTRACTION OF MEDIUM PARAMETERS AND PERFORMANCE COMPARISON

The proposed unit cell in Figure 10(a) has been represented as an equivalent circuit model with a coupling coil composed



**FIGURE 11.** Performance comparison of effective permeability  $(\mu_f)$  and permittivity  $(\epsilon_f)$  based on full-wave simulation and analytical models. A close correlation of both  $\mu_f$  and  $\epsilon_f$  validates the accuracy of the analytical model.

of resonant parameters  $R_m$ ,  $L_m$ , and  $C_m$  [9], [10], as shown in Figure 10(b). The lumped circuit inductance, capacitance, and resistance are denoted by  $L_m$ ,  $C_m$ , and  $R_m$ , respectively. Equation (2) expresses the resistive parameter Rm of the unit cell as the sum of the lumped circuit resistance  $R_o$ , dielectric loss  $R_{dt}$  due to the copper coil, proximity loss  $R_{prox}$ , and skin effect loss  $R_{sk}$ . As exhibited in Eq. (2), the resistive parameter  $R_m$ , of the unit cell is expressed as the sum of the lumped circuit resistance ( $R_o$ ), dielectric loss ( $R_{dt}$ ) due to the copper coil, proximity loss ( $R_{prox}$ ), and skin effect loss ( $R_{sk}$ ).

$$R_m = R_{sk} + R_{prox} + R_{dt} + R_o \tag{2}$$

Moreover,  $R_{sk}$  is estimated as illustrated in Eq. 3

$$R_{sk} = R_{mc} \frac{t}{\delta \left(1 - e^{-\left(\frac{t}{\delta}\right)}\right) \left(1 + \frac{t}{\Delta r}\right)}$$
(3)

where  $R_{mc} = \frac{l}{\sigma t \Delta r}$ ;  $\sigma$  is the conductivity of the wire; t - thickness of the wire; l - length of rectangular spiral, and  $\delta$  is the skin depth as denoted in Eq. 4.

$$\delta = \sqrt{\frac{1}{\pi \mu f_r \sigma}} \tag{4}$$

where  $\mu$  is the permeability of the MTM. The proximity loss,  $R_{prox}$ , which arises from the interaction between the fields of each wire turn, can be expressed mathematically as:

$$R_{prox} = 0.1228 R_{dc} \left(\frac{\mu \Delta r \pi f_r}{R_{sh}}\right)^2 \tag{5}$$

Moreover, the resistance due to the dielectric material  $(R_{dt})$  and the length of the spiral conductor can be expressed as demonstrated in Equation (7) and Equation (6), respectively

$$l = 4(\Delta r - tN)(N - 1) - 4tN(N + 1)$$
(6)

$$R_{dt} = \frac{tan\delta}{2\pi f_r C_s} \tag{7}$$

It's worth noting that  $\tan \delta$  is the loss tangent of the Teflon substrate and N is the number of spiral rectangular turns. Subsequently, the lumped circuit capacitance ( $C_m$ ) of the metamaterial (MTM) unit cell denoted in Equation. (8) is represented by combining the compensation capacitor ( $C_{com}$ ) and the stray capacitance ( $C_s$ ) shown in Equation. (9).

$$C_{s} = \left[\frac{9}{10}\epsilon_{ar}\epsilon_{o} + \left(\frac{\epsilon_{tf}\epsilon_{o}}{10}\right)\left(\frac{t}{\Delta r}\right)l\right]$$
(8)

$$C = C_s + C_{com} \tag{9}$$

where  $\epsilon_{ar}$  and  $\epsilon_{tf}$  are the relative permittivity of air and Teflon substrate, respectively. It is worth stating that the resonant frequency of the MTM can be tuned by varying the value of  $C_{com}$ , accordingly. Further, the relationship between the lumped circuit inductance,  $L_m$ , of the unit MTM cell and the number of wire turn, N, is demonstrated in Eq. 10.

$$L_m = \frac{1.27\mu N^2 d_{av}}{2} \left[ ln \left( \frac{2.07}{\Theta} \right) + 0.18k + 0.13\Theta^2 \right]$$
(10)

where  $\Theta$  is the fill ratio and the average length,  $d_{av}$ , and k are given by Eq. 11 and Eq. 12, respectively.

$$d_{av} = \frac{\Delta r + t}{2} \tag{11}$$

$$k = \frac{\Delta r + t}{\Delta r - t} \tag{12}$$

Essentially, the holistic combination of  $L_m$  and  $C_m$  in the form of  $L_m C_m$  circuit resonance culminates in the resonant frequency,  $F_r$ , of the proposed MTM consistent with Eq. 13

$$F_r = \frac{1}{2\pi\sqrt{L_m C_m}} \tag{13}$$

The performance comparison of real effective permeability  $(\mu_f)$  and permittivity  $(\epsilon_f)$  of the proposed MTM based on full-wave HFSS simulation and derived analytical models are

presented in Fig. 11(a) and Fig. 11(b), respectively. From the results, it can be inferred that the simulated and theoretical results for both  $\mu_f$  and  $\epsilon_f$  are in close agreement, effectively validating the accuracy of the theoretical model. Moreover, the resonance frequency of the simulated and analytical real component of  $\mu_f$  and  $\epsilon_f$  are closely matched at  $\approx 1.2$  MHz. Furthermore, both the simulation and theoretical results exhibit negative real components of  $\mu_f$  and  $\epsilon_f$  at their respective resonance frequency, essentially verifying the left-handed behavior of the proposed MTM structure.

#### B. JAYA-BASED OPTIMIZATION OF METAMATERIAL DESIGN PARAMETERS FOR RESONANT FREQUENCY REDUCTION

The Jaya-based optimization algorithm [31] proposed in this study harnesses a population-based optimization algorithm to find the global minimum or maximum of the function in Eq. (14) [32]. It works by maintaining a population of candidate solutions while also iteratively updating them to find the best solution. The update rule works by moving each candidate solution towards the best solution in the population, and the movement is proportional to the distance between the two solutions. This ensures that the solutions move towards the global optimum, even when the population has converged to a local optimum.

Given the MegaHertz resonant frequency behavior of the proposed MTM (specifically, 1.2 MHz) in ANSYS EM simulations, the overarching objective of this section is to develop an optimization algorithm capable of mitigating its resonant frequency to the Kilohertz range. Due to the complex relationship between the effective permeability ( $\mu_f$ ) and physical parameters of the proposed MTM, it is formulated as an optimization problem exhibited in Eq. 14. This equation expresses the dependence of  $\mu_f$  and resonant frequency,  $F_r$ , on the physical parameters of the MTM, including inter-turn spacing ( $\Delta r$ ), number of turns of copper wire (N), and copper wire thickness (t).

$$\mu_{f}(\Delta r, N, t) = 1 - \left[\frac{\left[\frac{l}{t}\right]^{2}}{1 + \frac{\rho}{\pi \Delta r F_{r} \mu_{o}(N-1)} - \frac{c_{o}^{2}}{l^{2} \Delta r(N-1)}}\right]$$
(14)

where  $c_o$  is the speed of light,  $\rho$  is the resistivity of copper wire. The optimization objective is to reduce the resonant frequency where evanescent wave amplification (negative permeability,  $\mu_f$ ) occurs in the proposed MTM.

#### 1) DESIGN PARAMETER OPTIMIZATION CONSTRAINT

To ensure the optimized MTM can be implemented in a practical manner, three objective variables are considered for the analysis:  $\Delta r$ , N, and t, as described in Equation (15)

$$F_r = f(\Delta r, N, t) \tag{15}$$



FIGURE 12. Flowchart Architecture of Proposed Optimization Algorithm for reducing the resonant frequency of the proposed metamaterial.

where the resonant frequency,  $F_r$  represents the objective function.

In perspective, the optimization problem to be solved is a minimization of the given relation in Eq. (16). The constraints for each variable are also illustrated

$$\begin{array}{ll} \min & f(\Delta r, N, t) \\ s.t. & \Delta r_{min} < \Delta r \leq \Delta r_{max} \\ & N_{min} < N \leq N_{max} \\ & t_{min} < t \leq t_{max} \end{array}$$
(16)

It is noteworthy that the physical parameters of a metamaterial, specifically N and  $\Delta r$ , have a significant impact on the inductance  $(L_m)$  and capacitance  $(C_m)$ , respectively, which consequently affect the resonant frequency of the MTM. Therefore, optimizing these parameters can potentially lower the resonant frequency, as shown in Eq. (15). The design constraints provided below take these factors into consideration:

- 1) The inter-turn spacing or clearance between adjacent turns, denoted by  $\Delta r$ , is crucial for controlling the parasitic capacitance and leakage electromagnetic field of the resonator. The optimized value of  $\Delta r$  is chosen such that it is less than or equal to the original clearance, denoted by  $\Delta r_{max}$ , between adjacent turns before optimization.
- The number of wire turns, denoted by N, is chosen consistently with the area of the proposed MTM, denoted by A<sub>1</sub>, and the size of the substrate, denoted by A<sub>2</sub>.



FIGURE 13. Metamaterial effective permeability, showing the impact of optimizing structural parameters on the resonant frequency and evanescent wave amplification.

The optimal number of turns is determined such that N is greater than or equal to the actual number of wire turns, denoted by  $N_{max}$ , and the areas of the MTM and substrate are less than or equal to their respective maximum values, denoted by  $A_{1max}$  and  $A_{2max}$ , prior to optimization.

3) The wire thickness, denoted by t, is estimated based on the target transmit power, which consequently determines the maximum excitation current based on the wire ampacity. In perspective, the minimum thickness of wire is selected to withstand the ampacity such that tis less than or equal to the initial wire thickness prior to optimization ( $t_{max}$ ).

Figure 13 illustrates an implementation flowchart of the optimization procedure which is divided into three stages, namely Stage1, stage2 and stage3.

In stage 1, the initial values of key structure parameters/objective variables, such as wire thickness (t), number of turns (N), and inter-turn spacing  $(\Delta r)$ , are chosen based on practical applications. The optimization algorithm iteratively adjusts the wire thickness until an optimal value that can handle the ampacity is obtained. The number of turns, N, is determined using the area of the MTM and substrate size, while a moderate value of r is used to minimize parasitic effects and stray magnetic fields.

In stage 2, a moderate number of turns is determined iteratively consistent with the MTM's area and substrate size while checking the corresponding resonant frequency. Finally, stage 3 investigates the optimized design parameters while tuning other parameters to achieve optimal results. Figure 14 demonstrates the effect of optimized design parameters of the MTM on the effective permeability and resonant frequency of the structure. Figure 13.

In perspective, for a fixed wire width, t = 3 mm and interturn spacing,  $\Delta r$  iterated from 0.5 mm to 1 mm at an offset



**FIGURE 14.** Fabricated prototype of the proposed metamaterial (MTM) structure, indicating the Unit MTM and an MTM-slab based on a periodic 4 × 4 array of unit cell.

of 0.25 mm, it is observed that the MTM effectively exhibits negative real permeability at resonant frequencies 750 kHz and 850 kHz corresponding to  $\Delta r_1 = 0.5$  mm and  $\Delta r_3 =$ 1 mm, respectively, whereas  $\Delta r_2 = 0.75$  mm demonstrates positive permeability. The aforementioned apparently shows that an optimal selection of  $\frac{t}{\Delta r}$ , does effectively minimizes resonant frequency while also influencing its evanescent wave amplification property. While the optimization process yielded three resonant frequency points as shown in Figure 13, only the design parameters corresponding to the lowest optimal resonant frequency (750 kHz) was harnessed for the experimental prototype. Moreover, the selection of 750 kHz frequency is based on practical integration with wireless charging systems that operate at 85 kHz, consistent with Airfuel Alliance Standard [33], [34]. The optimal design parameters for the 750 kHz resonant frequency were used to fabricate the MTM structure, and experimental testing was conducted to verify the accuracy of the optimization process, as described in Section IV.

#### IV. EXPERIMENTAL MEASUREMENT AND PERFORMANCE VALIDATION

In order to experimentally validate the proposed MTM, a prototype sample was fabricated as shown in Figure 14, and performance measurements were conducted using the prototype testbench shown in Figure 15. The MTM-slab, which consisted of a  $4 \times 4$  periodic MTM array, was constructed using PCB technology and single-strand copper wires of 26 AWG, and positioned on a 1mm thick Teflon substrate to securely hold the copper wires. Moreover, compensation capacitors were soldered beneath the PCB board based on the specifications in Table 3.

Furthermore, the Tx-coil and Rx-coil were connected to port 1 and port 2 of the KeySight EN5061A Vector Network Analyzer (VNA) using  $50\Omega$  micro SMA connections. External holes were created in the substrate slab to provide a support framework for the experimental set-up. Measurements are taken using both port 1 and port 2 of the VNA to measure the transmission coefficient,  $S_{21}$ , reflection coefficient,  $S_{11}$ , and Z-parameters. The MTM slabs were then combined with the two coil structures shown in Figure 15 while taking measurements for sweep values of the working distance (slab1



FIGURE 15. Hardware set-up of the proposed WPT system incorporating MTM slab showing Keysight ENA 5061A Vector Network Analyzer. Tx and Rx positions are kept constant while varying the position of Slab 1 relative to slab2.

TABLE 3. Experimental design parameters.

| Parameter               | Metamaterial                         | Transmitter | Receiver |
|-------------------------|--------------------------------------|-------------|----------|
| Turns                   | 10                                   | 3           | 3        |
| Inner diameter          | N/A                                  | 20 mm       | 20 mm    |
| Outer diameter          | N/A                                  | 30 mm       | 30 mm    |
| Trace Width             | 2 mm                                 | 2 mm        | 2 mm     |
| Copper wire guage       | -                                    | 26 AWG      | 26 AWG   |
| Capacitance $(C_{com})$ | 600 pF                               | N/A         | N/A      |
| Size                    | $64 \times 32 \times 3 \text{ mm}^3$ | N/A         | N/A      |
| Substrate               | Teflon                               | PLA         | PLA      |

to slab 2 distance) ranging from 50 mm to 250 mm. To prevent the slab1 and slab2 from potentially intersecting the Rx-coil and Tx-coil, respectively, the distance between both the Tx-coil and slab1, as well as Rx-coil and slab2, was kept at 0.5mm while retaining the working distance/loading position as depicted in Figure 15. The Tx-coil was stimulated using a start frequency of 100 kHz (0.1 MHz), a stop frequency of 10 MHz, and 671 data points. The collected data including the scattering parameters ( $S_{21}$  and  $S_{11}$ ), Z-parameters, and working distances were analyzed to determine the power transfer efficiency (PTE) of the proposed system. In order to compare performance, an identical testing technique was performed on the WPT system without the MTM-slab.

#### V. PROTOTYPE MEASUREMENTS AND DISCUSSION OF RESULT

The main objective of this section is to experimentally validate the resonant operating frequency of the proposed MTM. Based on the prototype set-up in Fig 14 and Fig. 15, the experimentally retrieved transmission coefficient ( $S_{21}$ ) and reflection coefficient ( $S_{11}$ ) of the proposed MTM exhibit resonant behavior at a frequency of 743.5 kHz. It is apparent that this frequency is closely related to the resonant frequency obtained from the optimization result, effectively validating the accuracy of the optimization algorithm. In addition, it is noteworthy that the resonant frequency of the proposed MTM is significantly lower in comparison to the resonant frequencies of existing MTM designs, thus making it an attractive candidate for high-efficiency wireless charging applications.

Furthermore, when the loading position of slab 1 relative to slab 2 is varied from 50 mm to 250 mm and the



**FIGURE 16.** Experimental waveform of the proposed MTM's Transmission coefficient ( $S_{21}$ ). For all working distance ranging from 50 mm to 250 mm, the MTM resonates at a frequency,  $F_r = 745.5$  kHz.



**FIGURE 17.** Experimental waveform of the proposed MTM's reflection coefficient ( $S_{11}$ ). The MTM shows a resonant behaviour at  $F_r = 745.5$  kHz for all working distance ranging from 50 mm to 250 mm.

separation/offset between both Tx and slab1, as well as Rx and slab2, is fixed at 0.5 mm, the resonant frequency of the MTM remains fairly constant at 743.5 kHz, as indicated by the plots of  $S_{21}$  and  $S_{11}$  corresponding to Figure 16 and Figure 17, respectively. The magnitude of  $S_{21}$  and  $S_{11}$  is significantly influenced by the loading position, while the resonant operating frequency remains constant.

## A. MEASUREMENT OF TRANSFER EFFICIENCY ( $\eta$ ) OF PROPOSED METAMATERIAL BASED WPT STRUCTURE

Consistent with [35] and [10], the power transfer efficiency  $(\eta)$  was investigated using the experimentally extracted *Z*-parameters (using the prototype test-bench in Figure 15) and the expression in Eq. (17),

$$\eta = \frac{|Z_{21,\text{ef}}|^2}{\Re\{Z_{11,\text{ef}}\}R_L} \times 100 \quad [\%].$$
(17)



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**FIGURE 18.** Performance comparison of experimential and simulated WPT efficiency with and without MTM.

where  $Z_{21,ef}$  is the extracted transmission Z-parameter from port 1 to port 2, whereas  $Z_{11,ef}$  is the extracted reflection Z-parameter from port 1 to port 1. Further, experimental measurements were performed using the same working distance (MTM to MTM) distance harnessed for the ANSYS electromagnetic simulation. The simulated and experimental waveform of PTE as a function of normalized distance  $\left(\frac{X_{max}}{Y}\right)$  for the system with and without MTM is exhibited in Fig. 18. The normalized distance is evaluated as the ratio of the maximum working distance,  $X_{max}$  (slab1 to slab2 distance), and parametric working distance (X) harnessed in the electromagnetic simulation. Apparently, the observed close matching of the simulation and experimentally measured transfer efficiency confirms the accuracy of the measurement results. The slight discrepancy could be traced to the inherent approximation in the computation of the lumped element impedances. Moreover, when the WPT system is integrated with two MTM slabs (each located close to the Tx and Rx), a significant increase in PTE for all normalized distances is observed compared to the WPT system without MTM.

#### **VI. CONCLUSION**

In this study, a metamaterial structure has been analyzed for low resonant behavior and evanescent wave amplification of a near magnetic field. Full-wave electromagnetic simulation has been performed in ANSYS HFSS electromagnetic solver culminating in design characterization based on extracted scattering parameters coupled with effective permeability and permittivity of the medium. Moreover, a meta-heuristicbased optimization of MTM physical parameters has been implemented to lower the resonant frequency of the MTM to 750kHz. The optimization analyzes the impact of design parameters on the resonant frequency where evanescent wave amplification occurs. Performance validation and verification of the accuracy of the optimization results with experimental measurements have been performed using a fabricated MTM sample. Apparently, the resonant frequency resulting from measurement per fabricated prototype closely matches the optimization result, essentially validating the accuracy of the optimization algorithm Thus, as a future extension of this work, advanced optimization algorithms are currently being developed to explore lower resonant frequency regions, preferably 85 kHz.

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