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# Parameter Extraction for the PSPHV LDMOS Transistor Model

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**ABSTRACT** This paper details a robust parameter extraction flow for the PSPHV LDMOS transistor model. The procedure uses a global scaling parameter set and accounts for self-heating. We describe how to determine parameters associated with important physical effects specific to PSPHV: non-uniform lateral channel doping; the Kirk effect; internal drain voltage clamping; and the drain expansion effect. The method is verified on devices from different technologies. Verilog-A code for PSPHV is publicly available.

**INDEX TERMS** MOSFET, LDMOS, parameter extraction, semiconductor device modeling, SPICE.

## I. INTRODUCTION

Laterally diffused MOS (LDMOS) transistors are widely used in integrated circuits (ICs) for applications such as switch-mode power supplies, power amplifiers, power management, and motor drivers. Recently, we developed an accurate surface-potential based compact LDMOS transistor model PSPHV [1] to help design and optimize such ICs. PSPHV consists of an enhanced version of the PSP103.6 MOS transistor model [2] for the intrinsic channel, a modified version of the JFETIDG dual-gate JFET model [3] for the drift region, and JUNCAP2 [4], extended to improve modeling at high forward bias, for parasitic junction diodes. These components are integrated into a single Verilog-A code which is publicly available [5].

Compared to existing compact models for LDMOS transistors, such as SPHV [6], BSIM-HV [7], and the CMC standard model HiSIM\_HV [8], PSPHV:

- models nonuniform lateral doping (NULD), from out-diffused channels and source-side halo implants
- models gradual channel turn-on
- does not exhibit unphysical spikes in capacitance versus voltage [9]–[12]
- includes a significantly more accurate avalanche current model, that accounts for the Kirk effect [13]
- as an option, can include reverse recovery modeling for the junction diodes.

A model is not useful without a procedure to extract parameters of the model to fit a wide variety of device types. The parameter extraction process for PSPHV is reasonably complex because LDMOS transistors are complex devices, hence PSPHV is a fairly complex model. This paper presents a robust extraction flow for PSPHV, to help modeling engineers use PSPHV for the devices they need to model. In particular, we introduce a self-heating effect (SHE) aware extraction flow. Previously proposed LDMOS model parameter extraction procedures, e.g., [14], determine values for temperature coefficients as a final step. However, output characteristics at the reference temperature  $T_r$  cannot be fitted without good values for temperature coefficients, because behavior at high gate bias  $V_{gs}$  and drain bias  $V_{ds}$  is affected by SHE. Our procedure improves on the prior art because it self-consistently accounts for SHE.

Our procedure includes extraction of capacitance parameters, but the focus is extraction of drain current parameters for full channel width and length scaling. We show the effect of parameters on capabilities unique to PSPHV: drain voltage scaling, the Kirk effect, internal drain voltage clamping, and the drain expansion effect. Our method is verified using devices from different technologies. Extraction of parameters for other parts of the model, such as gate current, parasitic junction diodes, noise, etc. is not detailed here, and follows the procedures defined for PSP [2].

In some instances, there may be unexpected device characteristics that cause issues with using a “canned” extraction program on a new device. Our procedure recognizes this fact: below “extract” means either an initial manual tuning, or small optimization, of one or a small number of parameters to fit a limited set of data, and “optimization” means the step is more substantial, and not manual, and involves adjusting multiple parameters to fit multiple data sets, often over geometry and temperature as well as over bias.

## II. OVERVIEW AND GLOBAL SWITCHES

PSPHV inherits the local-global hierarchical parameter structure of PSP [15]. Model parameters can be specified at a local level, i.e., for one specific geometry, or at a global level. Here we use global level parameters. All PSPHV parameters are in SI units and here are typeset in constantWidth font.

PSPHV can model LDMOS transistors with different drift region structures, see Fig. 1. For devices like those, set  $swtgm\os = 1$  and  $swbgm\os = 0$ ; this defines the type of the top and bottom gates of the drift region as MOS and p-n junction, respectively.

Fig. 2 shows the equivalent circuit for PSPHV (this includes parasitics series resistance not explicitly shown in [1]). The optional second gate-drain overlap capacitance is turned on by specifying a positive value for the instance parameter  $lgdov2$ , the length of the second gate-drain overlap region, see Fig. 1, (note that the bulk connection of this component is moved from the di to the d node compared to [1]; this improves accuracy).

For LDMOS transistor structures like those of Figs. 1(a) and 1(c), the first gate-drain overlap region, whose length is specified by the instance parameter  $lgdov$ , covers only a small portion of the overall drift region length. The JFETIDG model, which is used for the drift region, assumes that depletion pinching occurs over the whole length of that region (which is true for the bottom gate of the drift region, the bulk). If parameters for the JFETIDG top gate depletion pinching effect are extracted at low  $V_{ds}$  they unphysically predict the drift region will pinchoff at high  $V_{ds}$ . For such structures, set  $dfto = 0$ , to turn off the top gate depletion pinching effect, and instead model drift region resistance modulation by the top gate via the  $mumto$  mobility modulation parameter.

For LDMOS transistor structures like Fig. 1(b),  $lgdov$  covers a substantial portion of the drift region and the JFETIDG top gate depletion pinching model can be used. Initial estimates of top gate depletion pinching factor  $dfto$  and built-in potential  $\psi_{rto}$  are

$$d_{fto} = \frac{\sqrt{\epsilon_s/(qN_c)}}{t_m + t_{ox}\epsilon_s/\epsilon_{ox}}, \quad \psi_{rto} = \frac{q\epsilon_s N_c t_{ox}^2}{\epsilon_{ox}^2} - 2V_{FB,ov} \quad (1)$$

where  $t_m$  is the drift region metallurgical thickness and  $t_{ox}$  is oxide thickness;  $\epsilon_{ox}$  and  $\epsilon_s$  are the permittivities of SiO<sub>2</sub> and Si, respectively,  $q$  is the magnitude of the electronic charge,  $N_c$  is the doping concentration in the drift region,

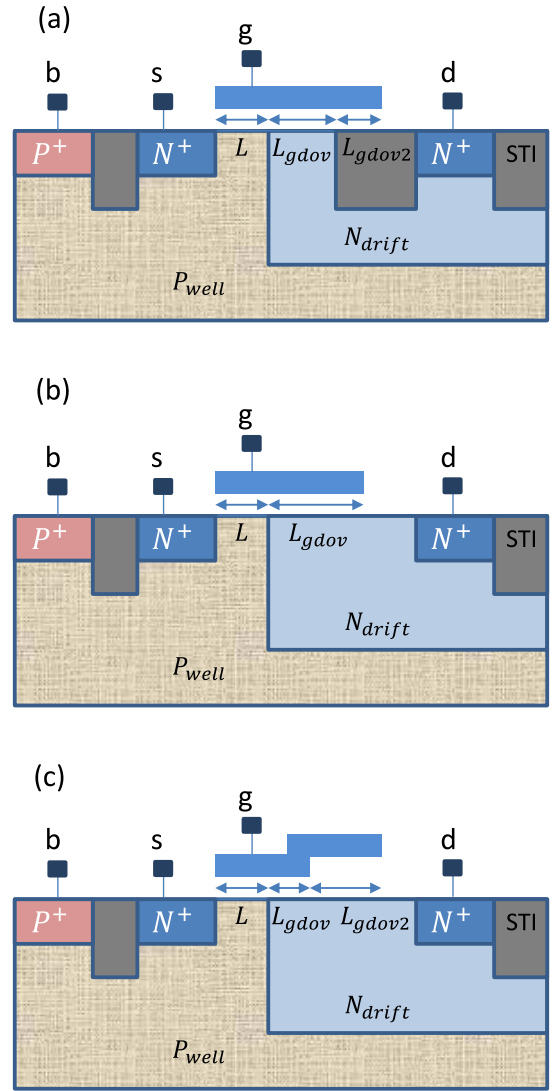


FIGURE 1. Cross-section of n-LDMOS transistors with different drift region designs.

assumed to be constant, and  $V_{FB,ov}$  is the flat-band voltage of the overlap region.

Other important global switch settings:

- $swet$ : set to 1 to turn on self-heating modeling (in which case either the thermal node must be connected to an external thermal network or the thermal resistance must be positive)
- $swvdscale$ : set to 1 to turn on  $V_{ds}$  scaling to model NULD (this triggers a separate surface potential solution for charge calculations for the intrinsic MOS transistor, as when  $swdelvtac = 1$ )
- $swvdimp$ : set to 1 to make PSP impact ionization calculations use  $V(d,s)$  instead of  $V(di,s)$  (this is necessary if  $V_{di}$  clamping is activated, when  $swvdiclmp = 1$ )
- $swdrlin$ : set to 1 to have JFETIDG treat the drift region as a linear resistor (for debugging or to speed up simulations, with associated loss of accuracy of course)

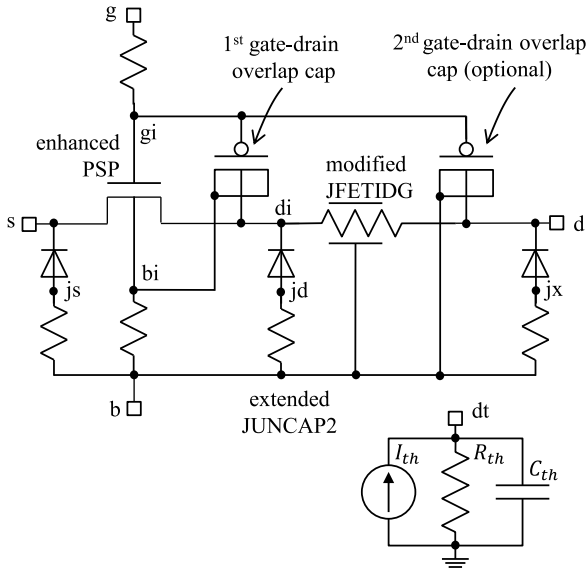


FIGURE 2. Large-signal equivalent circuit for PSPHV.

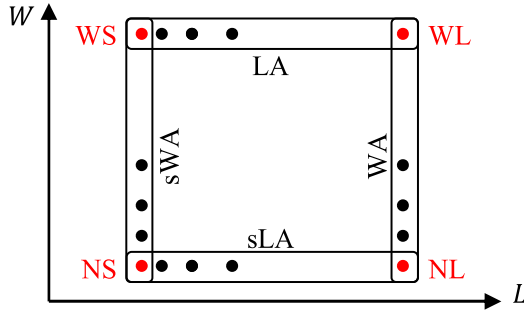


FIGURE 3. Device geometries for extraction.

- for swgeo, swedge, swigate, swimpact, swgidl, swjunc, swjunasym, swnud, and swdelvtac refer to the PSP manual [2]

### III. DEVICE GEOMETRIES AND MEASUREMENTS

Figure 3 shows the device geometries that are needed to extract a fully scalable PSPHV model. WL, WS, NL, NS stand for the wide/long, wide/short, narrow/long, narrow/short devices, respectively. Besides the four corner geometries, there are primary and secondary width  $W$  and length  $L$  arrays. Because short and narrow channel effects vary roughly as  $1/L$  and  $1/W$ , respectively, the length and width arrays should space geometries roughly equally in  $1/L$  and  $1/W$ ; this maximizes observability of parameters.

The following data are needed ( $x = d, g, s, b$ ):

- DC data (all geometries and temperatures)
  - *idvg\_lin*:  $I_d(V_{gs})$  at low  $V_{ds}$ , over  $V_{bs}$  if the bulk and source are not tied together
  - *idvg\_sat*:  $I_d(V_{gs})$  and  $I_b(V_{gs})$  at three or more high  $V_{ds}$  values and  $V_{bs} = 0$
  - *idvd data*:  $I_d(V_{ds})$  and  $I_b(V_{ds})$  for several  $V_{gs}$  in strong inversion,  $V_{bs} = 0$

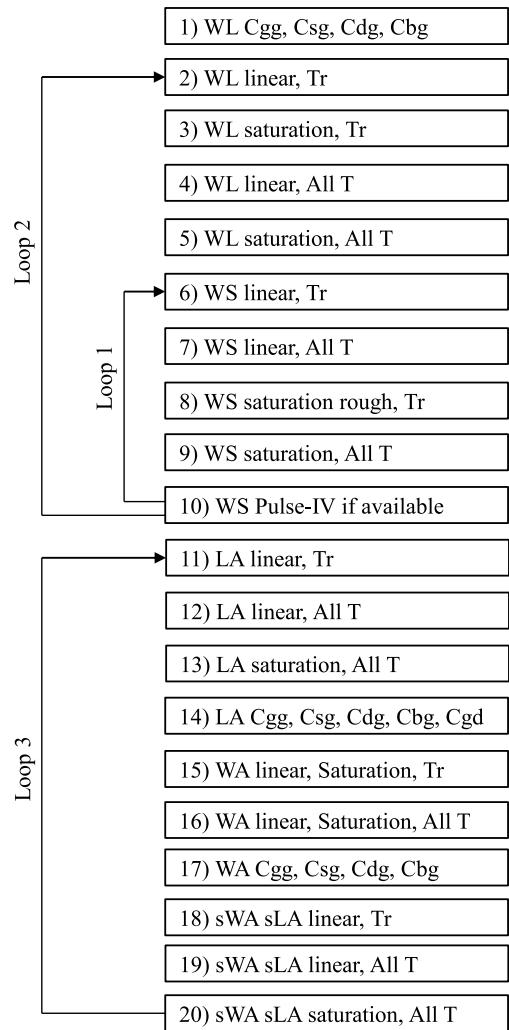


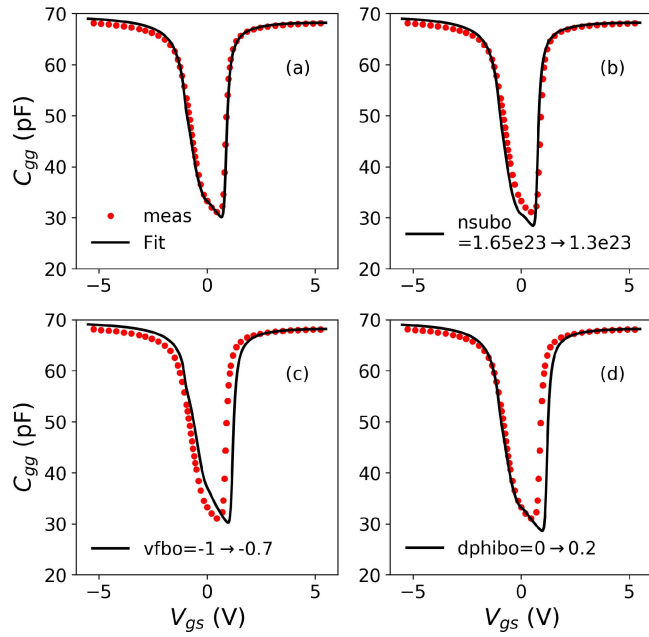
FIGURE 4. Extraction flow.

- CV data for WL, WS, NL, NS devices at  $T_r$ 
  - *cxg*:  $C_{xg}(V_{gs})$  at  $V_{ds} = V_{bs} = 0$
  - *cgd*:  $C_{gd}(V_{ds})$  at  $V_{gs} = V_{bs} = 0$
- (optional) *s*-parameter data for the WS device at  $T_r$ 
  - *y22*:  $y_{22}(f)$  from 10 kHz to 10 MHz,  $V_{gs} = V_{gs,max}/2$ ,  $V_{ds} = V_{ds,max}/2$ ,  $V_{bs} = 0$
  - *cxg\_spar*:  $C_{xg}(V_{gs})$  from *s*-parameters, for  $V_{gs}$  from  $-V_{gs,max}$  to  $+V_{gs,max}$ , several values of  $V_{ds}$  from 0 to  $V_{ds,max}/2$ ,  $V_{bs} = 0$ ,  $f = 2$  GHz
- (optional) pulsed IV data for the WS device at  $T_r$ 
  - *idvd\_pulse*:  $I_d(V_{ds})$  and  $I_b(V_{ds})$  for several  $V_{gs}$  in strong inversion,  $V_{bs} = 0$

To avoid damaging devices, the highest  $V_{ds}$  for *idvg\_sat* and *idvd* data may be limited to  $V_{ds,max}/2$ . Structures for capacitance and *s*-parameter data should have a sufficient number of gates  $N_g$  to enable reliable measurement.

### IV. EXTRACTION FLOW

Fig. 4 shows the steps in our flow, which are detailed below. Unless otherwise stated, the data are for a 45V n-LDMOS



**FIGURE 5.**  $C_{gg}$  of a WL transistor,  $W/L/N_g = 25 \mu\text{m}/10 \mu\text{m}/112$ ,  $V_{ds} = V_{bs} = 0 \text{ V}$ .

transistor in a 90nm BCD technology at  $T = T_r = 25 \text{ }^\circ\text{C}$ , measured data are plotted as symbols, and PSPHV model results are plotted as lines. If devices have the bulk and source tied, omit parameters and steps that require  $V_{sb}$  and  $I_b$  data.

Following the PSP parameter naming convention, paramo is the geometry independent part for param (WL device); paraml models the L dependence (WS device); paramw models the W dependence (NL); and paramlw models the area dependence (coupled L and W dependence, NS device). Refer to the Appendix for the meaning of each parameter.

Selection and technology parameters:

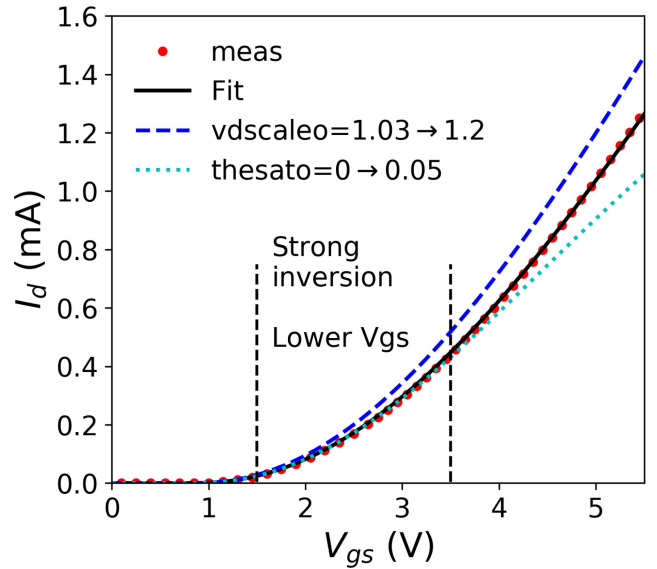
- set type to +1/-1 for n/p-type transistors, respectively
- set tr and trj to the reference temperature
- if the device has no halo implant, set lpck = 0
- set lgdov based on the device cross-section and layout
- for devices with a second gate-drain overlap region, set lgdov2 and toxovd2o from technology data

JFETIDG has several ways to define and calculate the sheet resistance of the drift region, see [3, Fig. 4]; we recommend specifying rsh0.

*Step 1:* device: WL, data: cgg, temperature:  $T_r$ .

Extract the gate oxide thickness toxo by fitting the capacitance data for  $V_{gs} = \pm V_{gs,max}$ . Set toxovo = toxovdo = toxo. The bulk doping nsubo then determines the minimum  $C_{gg}$ , see Fig. 5(b). The flat band voltage vfb0 shifts the whole  $C_{gg}$  curve horizontally, as Fig. 5(c) shows, and should be used to fit the data in depletion. The offset of the bulk potential  $\phi_B$ , dphibo, only affects  $C_{gg}$  in inversion as Fig. 5(d) shows. Fig. 5(a) shows the final fit. If the influence of poly depletion is apparent include npo.

*Step 2:* device: WL, data: idvg\_lin, temperature:  $T_r$ .



**FIGURE 6.**  $I_d(V_{gs})$  of a WL transistor,  $W/L/N_g = 25 \mu\text{m}/25 \mu\text{m}/2$ ,  $V_{ds}=28 \text{ V}$  and  $V_{bs}=0$ .

Extract nsubo, dphibo, and the interface charge factor cto to fit the data in weak inversion; if gradual channel turn on is apparent include the gate dependence of interface charge ctgo, see [1, Fig. 6]. For the data in strong inversion, extract mobility related parameters uo, mueo, themuo and, at high  $V_{bs}$ , the non-universality parameter xcoro.

*Step 3:* device: WL, data: idvg\_sat, temperature:  $T_r$ .

Extract the velocity saturation parameters thesato and thesatgo at  $V_{bs} = 0$ , thesatbo at high  $V_{bs}$ . If the model underestimates  $I_d$  even with thesato = 0 use the  $V_{di}$  scaling parameter vdscaleo. Fig. 6 shows the effect of vdscaleo and thesato. vdscaleo affects strong inversion for all  $V_{gs}$ , thesato mainly impacts the higher  $V_{gs}$  region. Therefore, use vdscaleo to fit strong inversion data at lower  $V_{gs}$ . This  $V_{ds}$  scaling technique [16] is an important feature of PSPHV, not available in PSP, and significantly improves modeling of short channel devices where NULD is significant. A similar technique is also used in the BSIM-BULK model [17].

Extract the smoothing parameter aro of the new  $C^\infty V_{dsat}$  clamping function in PSPHV [18] by fitting the nonsaturation to saturation transition region.

Extract the avalanche current parameter a1o, a2o, a3o and a3cvd1; a3cvdo to fit  $I_b$ . Fig. 11(b) shows the influence of a3cvd1; a3cvdo has similar effect.

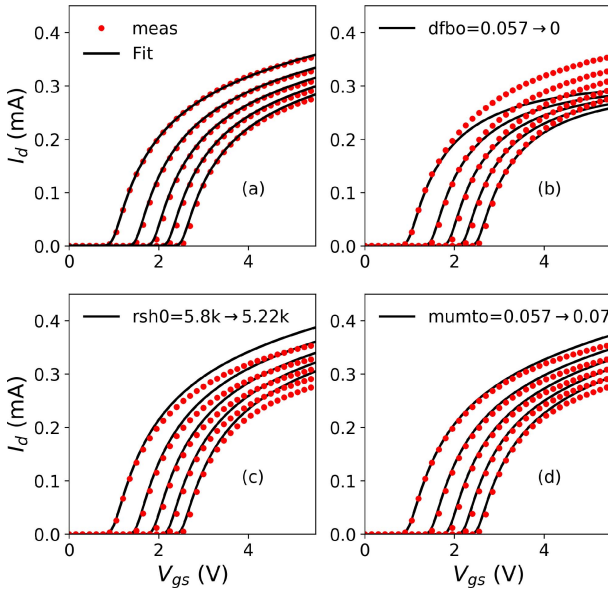
*Step 4:* device: WL, data: idvg\_lin, temperature: all.

Extract the temperature dependence parameters stvfb0, stbeto, stmueo, and stthemueo at  $V_{bs} = 0$ . From the data at high  $V_{bs}$  extract stxcoro.

*Step 5:* device: WL, data: idvg\_sat, temperature: all.

Extract stthesato and, if vdscaleo is used, stvdscaleo.

The temperature dependence parameters for the WL device, which is not significantly affected by SHE, are



**FIGURE 7.**  $I_d(V_{gs})$  of a WS transistor,  $W/L/N_g = 25 \mu\text{m}/1 \mu\text{m}/2$ ,  $V_{bs} = 0$ ,  $-1.25, -2.5, -3.75, -5 \text{ V}$ ,  $V_{ds} = 0.1 \text{ V}$ .

extracted before fitting the WS device, to ensure those parameters for the long channel device are correct.

*Step 6:* device: WS, data: idvg\_lin, temperature:  $T_r$ .

Extract  $\text{ct1}$ ,  $\text{dphib1}$ , and  $\text{foll}$  to fit  $I_d$  in weak inversion. The intrinsic series resistance  $\text{rsw1}$  can be set to zero, its impact is small and can be subsumed into the drift resistance. Extract the drift region bottom gate depletion parameter  $\text{dfbo}$  to match the spacing between the curves for different  $V_{bs}$  at  $V_{gs,max}$ , see Fig. 7(b) for the impact of  $\text{dfbo}$ . Roughly extract the effective channel length offset parameter  $\text{lvaro}$  to fit the peak  $g_m$ . Figs. 7(c) and 7(d) show how  $\text{rsh0}$  and  $\text{mumto}$  affect fitting;  $\text{rsh0}$  has an influence over a wide range of  $V_{gs}$  whereas  $\text{mumto}$  affects fitting at high  $V_{gs}$ . By tuning  $\text{lvaro}$ ,  $\text{rsh0}$ , and  $\text{mumto}$ , the nonsaturation data can be fitted well, as Fig. 7(a) shows. Fig. 8(a) shows  $g_m$  at different  $V_{bs}$ .

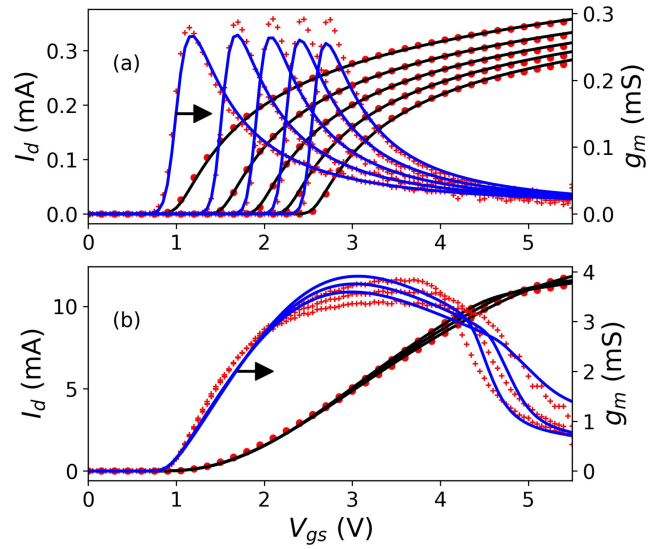
*Step 7:* device: WS, data: idvg\_lin, temperature: all.

Extract the temperature dependence of the intrinsic MOS transistor channel mobility  $\text{stbet1}$  from peak  $g_m$ . Extract the temperature dependence of the drift region mobility, and hence resistance,  $\text{tc1o}$  and  $\text{tc2o}$  from the high  $V_{gs}$  region. Coupled through SHE, these parameters affect the saturation characteristics at  $T_r$ , see Fig. 9. This is why the temperature dependence of the nonsaturation characteristics must be extracted before fitting the saturation characteristics at  $T_r$ .

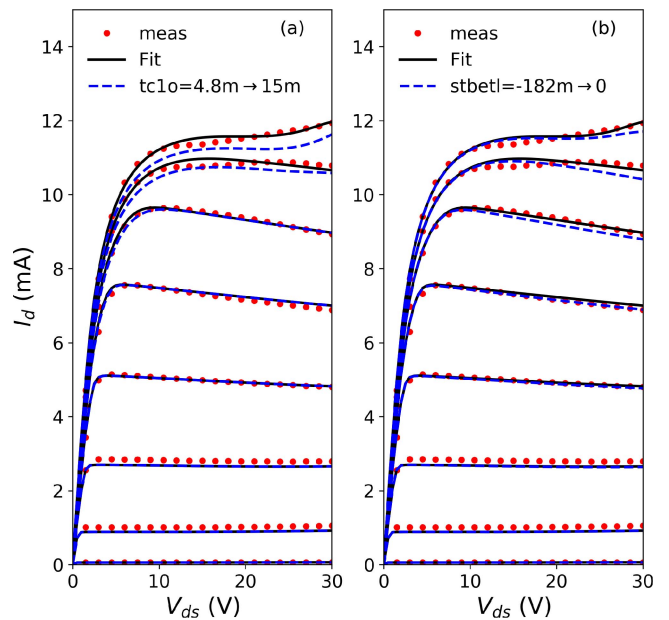
Note that  $\text{tc1o}$  and  $\text{tc2o}$  directly affect the temperature dependence of  $\text{ecrit}$  through  $E_{crit} = V_{sat}/\mu$ . This coupling complicates extraction but is crucial for self-consistent modeling of output characteristics with and without SHE, as noted in [1].

*Step 8:* device: WS, data: idvg\_sat and idvd, temperature:  $T_r$ .

Here, the device characteristics are impacted by many physical effects such as the velocity saturation of the intrinsic



**FIGURE 8.** (a)  $I_d(V_{gs})$  and  $g_m(V_{gs})$ , same transistor and biases as Fig. 7. (b)  $I_d(V_{gs})$  and  $g_m(V_{gs})$  for the same transistor,  $V_{ds} = 12, 20, 28 \text{ V}$ ,  $V_{bs} = 0 \text{ V}$ .



**FIGURE 9.**  $I_d(V_{ds})$  of the WS transistor as Fig. 7 at different  $V_{gs}$ . The temperature coefficients  $\text{tc1o}$  (a) and  $\text{stbet1}$  (b), extracted from  $\text{idvg\_lin}$  data, affect the saturation characteristics at  $T_r$ .

MOS, the quasi-saturation of the drift region, and the impact ionization in both the intrinsic MOS and the drift region. This is a complicated part of LDMOS extraction: fitting at  $T_r$ , in the presence of SHE, depends on temperature coefficients of parameters that have influence in saturation, however from the previous step we only have, so far, temperature coefficients for the parameters dominant at low  $V_{ds}$ . This step is a first rough extraction, that is refined in the next step.

8.1 Extract the DIBL parameter  $\text{cfl}$  to fit  $\text{idvg\_sat}$  data in weak inversion.

## 8.2 The thermal conductance in PSPHV scales as

$$g_{th} = g_{tho} + g_{thl}L + g_{thw}W + g_{tha}LW. \quad (2)$$

This points to “extract  $g_{tha}$  from the WL device.” However, SHE is weak for that device, to first order SHE decreases as  $1/L^2$ , so that is not possible. This leaves two practical options: extract  $g_{tha}$  from the WS device; or use TCAD to develop a scaling equation based on thermal simulation, then extract a scale factor to calibrate the TCAD-based model so the PSPHV results match silicon data. It is possible to determine  $g_{th}$  for devices that exhibit negative  $g_o$ . However, in practice  $g_{th}$  scaling can be significantly more complex than (2), especially when multiple device fingers are included. Because of the high currents involved, it can be difficult to measure  $I_d(V_{ds})$  at moderate to high  $V_{ds}$  for short multi-finger devices. For this reason we prefer to use TCAD to develop a scalable  $g_{th}$  model, including multi-finger devices.

$g_{th}$  determines the (negative) slope in output characteristics at high  $V_{gs}$ . Fig. 10(a) compares  $I_d$  data with and without SHE. The scale factor for the TCAD based  $g_{th}$  scaling model is extracted to fit the slope at moderate  $V_{gs}$  biases, see the ellipse in Fig. 10(a).

8.3 Use the  $idvg_{sat}$  data before the onset of quasi-saturation to extract  $thesat1$  and  $vdscale1$ ; their effect is the same as for the associated geometry independent parameters  $thesato$  and  $vdscaleo$ , respectively, see Fig. 6. Then use  $idvd$  data to extract the drift region velocity saturation parameters  $ecrito$  and  $ecorno$  for quasi-saturation effect; Fig. 10(b) and 10(c) show the effect of these parameters, respectively. The  $V_{dsat}$  smoothing parameter for short devices,  $ar1$ , can also be used to tune the transition region, in concert with  $ecorno$ .

At high  $V_{gs}$  and moderate  $V_{ds}$ , the parameter for the Early effect in the drift region  $clm11$  may be needed. Fig. 10(d), compared to Fig. 10(a), shows how  $clm11$  affects the output conductance. Do not make  $clm11$  too large otherwise it leads to overprediction of  $I_d$  in comparison to pulsed  $I_d(V_{ds})$  data when SHE is turned off.

At this point, do not fit the high  $V_{gs}$  and high  $V_{ds}$  region because the avalanche current parameters have not been extracted.

8.4 To extract parameters for the avalanche current, note that it has components both from the intrinsic MOS transistor, at lower  $V_{gs}$ , and from the Kirk effect in the drift region, at higher  $V_{gs}$ , see Fig. 11(a). If you have to invoke  $V_{di}$  clamping, see Section V, set  $swvdimp = 1$  to enable better avalanche current modeling.

Extract  $a11$ ,  $a31$ , and  $a3cvd1$  to fit  $I_b$  at lower  $V_{gs}$ , Fig. 11(b) shows the effect of  $a3cvd1$ . Extract Kirk effect related parameters  $alphabo$ ,  $beta$ ,  $mhc$ , and  $jhc$  to fit  $I_b$  at higher  $V_{gs}$ .  $alphabo$  controls the magnitude of the avalanche current in the drift region, Fig. 11(c) shows the effect of setting it to zero (the increase in  $I_b$  at high  $V_{gs}$  in that plot is the result of using  $V(d,s)$  instead of  $V(di,s)$  to calculate the avalanche current, from  $swvdimp = 1$ ).  $beta$  controls the spacing between the Kirk effect curves.

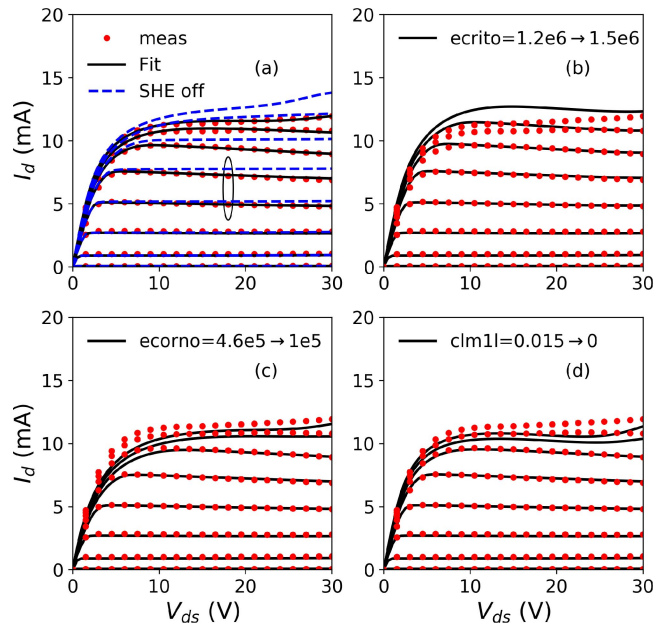


FIGURE 10.  $I_d(V_{ds})$  of the WS transistor of Fig. 7 at different  $V_{gs}$ ,  $V_{bs} = 0$ .

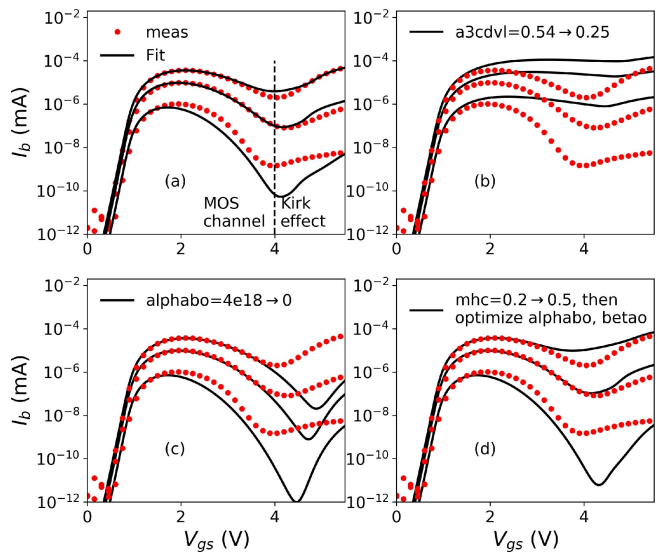


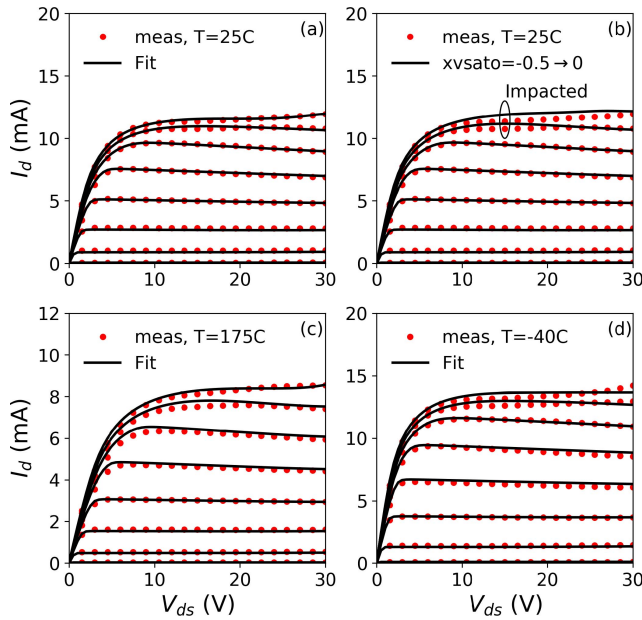
FIGURE 11.  $I_b(V_{gs})$  of the WS transistor of Fig. 7 at different  $V_{ds}$ ,  $V_{bs} = 0$ .

$mhc$  reflects the influence of the doping profile of the drift region, increasing it lowers  $I_b$  at low  $V_{ds}$ , see Fig. 11(d).

Step 9: device: WS, data:  $idvd$  and  $y22$ , temperature: all.

First, extract  $tegth$ , the temperature coefficient of  $g_{th}$ , to fit  $g_o$  at all temperatures. Then extract  $xvsato$ , the temperature dependence of velocity saturation in the drift region, see Fig. 12(b). These temperature dependence parameters also affect fitting at  $T_r$  because of SHE, which is why WS saturation parameters need to be fitted to data from all temperatures.

Next, extract  $sta2o$  and  $xbeta$  for the temperature dependence of avalanche current. Extract the drain expansion parameters  $mumi1o$  and  $mumi2o$  (Section VI gives



**FIGURE 12.**  $I_d(V_{ds})$  of the WS transistor of Fig. 7 at different  $V_{gs}$ , (a), (c), and (d) show fitting over temperature, (b) shows the impact of  $xvsato$ .

further details), and optimize  $ecrito$ ,  $gtha$ ,  $tegth$ , and  $xvsato$  to fit  $idvd$  data over all temperatures. Figs. 12(a), 12(c), and 12(d) show results after these steps at 25, 175, and  $-40^\circ\text{C}$ , respectively. Excellent fitting has been achieved. Fig. 8(b) shows that  $g_m$  at different  $V_{ds}$  is also modeled well.

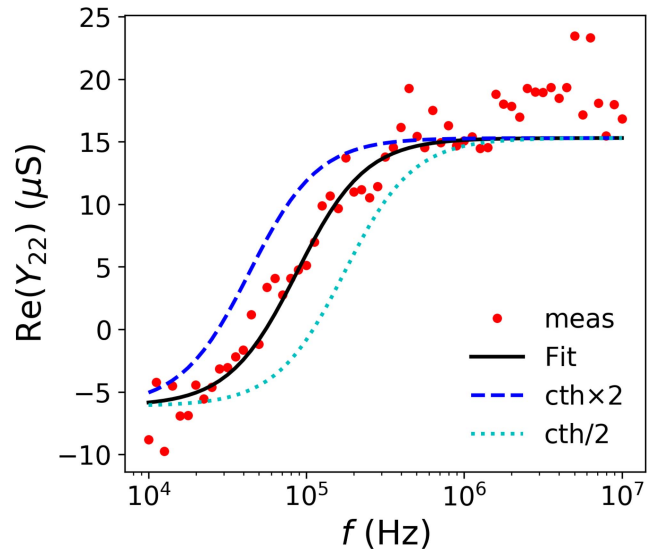
To extract thermal capacitance parameters, the best approach is to develop a scaling model using TCAD, then use a scale factor to fit that model to  $\Re(y_{22})$  versus frequency, see Fig. 13. The thermal capacitance  $c_{th}$  adjusts the frequency at which  $g_o$  transitions from being affected, to not being affected, by SHE.

If  $y_{22}$  data are not available, measurements from a variety of devices show that thermal time constants for integrated devices vary from about 0.1 to 10.0  $\mu\text{s}$ , with 1  $\mu\text{s}$  being a typical value. So, set the scaling and value such that  $c_{th}/g_{th} = 10^{-6}$  s.

*Step 10:* device: WS, data:  $idvd\_pulse$ , temperature:  $T_r$ .

Because pulsed IV data remove the SHE they are useful to verify, and refine if necessary, the accuracy of modeling of  $g_o$  of the intrinsic transistor. If  $I_d$  in saturation in the pulsed IV data are not modeled accurately, adjust the channel length modulation parameter  $clm11$  to improve the fit. See [1, Fig. 14] for an example.  $clm11$  changes the slope of output curves in the quasi-saturation region. Note that pulsed IV and DC data cannot be fitted simultaneously; they are typically measured on separate structures, so local variation between devices and differences in series resistance between the structures affect the results. Factor those differences in when adjusting to fit the pulsed IV data.

*Loop 1:* To improve fitting in both nonsaturation and saturation for the WS device, repeat steps 6 through 10, see Fig. 4.



**FIGURE 13.**  $\Re(Y_{22})$  vs.  $f$  of a WS 90V p-LDMOS transistor,  $V_{gs} = 3.5$  V,  $V_{ds} = 35$  V,  $T = 27^\circ\text{C}$ .

*Loop 2:* Repeat steps 1 through 10, see Fig. 4, to fine tune parameters for the WL device.

*Step 11:* devices: LA, data:  $idvg\_lin$ , temperature:  $T_r$ .

This step determines “shape” parameters for the length dependence of wide devices in nonsaturation. Optimize  $uo$ ,  $lvaro$ , and  $lvar1$  to fit peak  $g_m$ . (The other option is to use  $fbet1$  and  $lp1$  instead of  $lvar1$ .) Optimize  $cto$ ,  $ctgo$ ,  $ctl$ , and  $ctlexp$  to fit subthreshold slope. Optimize  $nsubo$ ,  $vfbo$ ,  $dphibo$ ,  $fol1$ ,  $fol2$ ,  $vfb1$ ,  $dphib1$ , and for devices with halo implants  $lpck$  and  $npck$ , to fit threshold voltage. Optimize  $mueo$ ,  $themuo$ ,  $xcoro$ ,  $rsh0$ ,  $mumto$ , and  $dfbo$  to fit data at high  $V_{gs}$ .

*Step 12:* devices: LA, data:  $idvg\_lin$ , temperature: all.

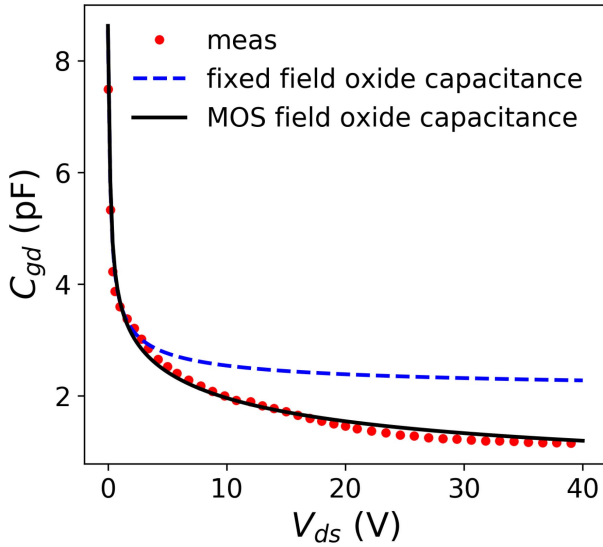
This step determines “shape” parameters for temperature dependence over length of wide devices in nonsaturation. Optimize the temperature dependence parameters  $stvfbo$ ,  $stbeto$ ,  $stmueo$ ,  $stthemueo$ ,  $stxcoro$ ,  $stvfbl$ ,  $stbet1$ ,  $tc1o$ , and  $tc2o$  to fit the  $idvg\_lin$  data over temperature.

*Step 13:* devices: LA, data:  $idvg\_sat$  and  $idvd$ , temperature: all.

This step determines “shape” parameters for the length and temperature dependence of wide devices in saturation. Optimize  $cfl$ , and  $cfllexp$  to model DIBL in the  $idvg\_sat$  data. Using both  $idvg\_sat$  and  $idvd$  data, optimize  $thesato$ ,  $thesat1$ ,  $thesatlexp$ ,  $thesatgo$ ,  $vdscaleo$ ,  $vdscale1$ , and  $vdscalele$  to fit  $I_d$  before the onset of quasi-saturation onset, optimize  $aro$ ,  $arl$ ,  $arlexp$ , and  $ecorno$  to fit the region around the transition to saturation, and optimize  $ecrito$ ,  $gtha$ ,  $tegth$ , and  $xvsato$  to fit the quasi-saturation region.

*Step 14:* devices: WS, data:  $cgg$ ,  $csg$ , and  $cdg$ , temperature:  $T_r$ .

This step fits source and drain overlap capacitances and the channel length offset for charge to fit total gate capacitance.



**FIGURE 14.**  $C_{gd}$  vs.  $V_{ds}$  for a WS transistor,  $W/L/N_g = 25 \mu\text{m}/1 \mu\text{m}/280$ ,  $V_{gs} = 0$ ,  $f = 1 \text{ MHz}$ ,  $T = 25^\circ\text{C}$ .

Extract  $lov$ ,  $vfbv$ ,  $novo$ , and  $cfrw$  to fit  $csg$ ,  $xlgdovq$ ,  $novdo$ , and  $vfbvd$  to fit  $cdg$ , and  $dlq$  to fit  $cgg$ . The field plate can be modeled as a bias independent capacitor or as a MOS capacitor through  $lgdov2$  and  $toxovd2o$ . As Fig. 14 shows, the latter greatly improves  $C_{gd}$  modeling.

*Step 15:* devices: WA, data:  $idvg\_lin$  and  $idvg\_sat$ , temperature:  $T_r$ .

This step determines “shape” parameters for width dependence of long devices in both nonsaturation and saturation. From the  $idvg\_lin$  data, extract the effective channel width offset parameters  $wvaro$  and  $wvarw$  to fit peak  $g_m$  (the other option is to use  $fbetw1$  and  $betw1$  instead of  $wvarw$ ), extract  $dphibw$  to fit threshold voltage, and extract  $muew$  to fit  $I_d$  at high  $V_{gs}$ . From the  $idvg\_sat$  data, extract  $thesatw$  and  $vdscalew$ .

*Step 16:* devices: WA, data:  $idvg\_lin$  and  $idvg\_sat$ , temperature: all.

This step determines “shape” parameters for temperature dependence over width of long devices in both nonsaturation and saturation. Extract  $stvfwb$  and  $stbetw$  to fit the  $idvg\_lin$  data over temperature. Extract  $stthesatw$  and  $stvdscalew$  to fit the  $idvg\_sat$  data over temperature.

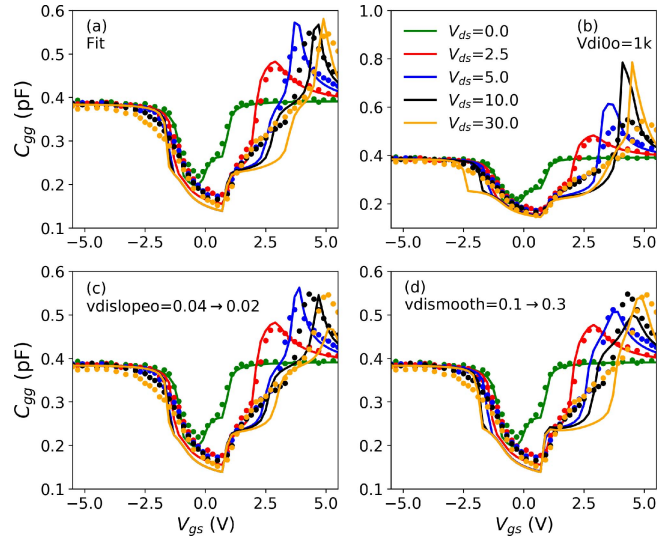
*Step 17:* devices: WA, data:  $cgg$  and  $cbg$ , temperature:  $T_r$ .

This step fits gate-bulk overlap capacitance and the channel width offset for charge to fit total gate capacitance. Extract  $cgbov1$  to fit  $cbg$ , extract  $dwq$  to fit  $cgg$ .

*Step 18:* devices: sLA and sWA, data:  $idvg\_lin$ , temperature:  $T_r$ .

This step determines “shape” parameters for coupled length and width interactions in nonsaturation. Extract  $lvarw$  to fit peak  $g_m$ . Extract  $vfb1w$  and  $dphib1w$  to fit threshold voltage. Extract  $ctlw$  to fit subthreshold slope. Extract  $xw$  to fit  $I_d$  at high  $V_{gs}$ . Extract  $dfbw$  to fit the  $V_{bs}$  dependence at high  $V_{gs}$ .

*Step 19:* devices: sLA and sWA, data:  $idvg\_lin$ , temperature: all.



**FIGURE 15.**  $C_{gg}$  vs.  $V_{gs}$  for a WS transistor,  $W/L/N_g = 10 \mu\text{m}/1 \mu\text{m}/8$ ,  $f = 2 \text{ GHz}$ ,  $T = 25^\circ\text{C}$ .

This step fits the temperature parameters for coupled length and width interactions in nonsaturation. Extract  $stvfblw$  to fit threshold voltage, and  $stbet1w$ ,  $tc1w$ , and  $tc2w$  to fit  $I_d$  at high  $V_{gs}$ .

*Step 20:* devices: sLA and sWA, data:  $idvg\_sat$  and  $idvd$ , temperature: all.

This step fits coupled length and width interactions in saturation. Extract  $thesat1w$ ,  $vdscale1w$ , and  $ecritw$  to fit both sets of data.

*Loop 3:* To improve fitting in both nonsaturation and saturation, over bias, geometry, and temperature, repeat steps 11 through 20, see Fig. 4.

## V. $V_{di}$ CLAMPING AND ITS EFFECT ON $C_{GD}$ AND $Q_G$

Figs. 15(a) and 15(b) show PSPHV modeling of gate capacitance with and without  $V_{di}$  clamping, respectively. The spikes at high  $V_{ds}$  visible in 15(b) are eliminated when  $V_{di}$  clamping is used, see 15(a). These spikes typically only occur for high voltage transistors; for low voltage transistors and transistors that do not exhibit such spikes in capacitance do not turn on  $V_{di}$  clamping (there is a computational cost associated with the clamping).

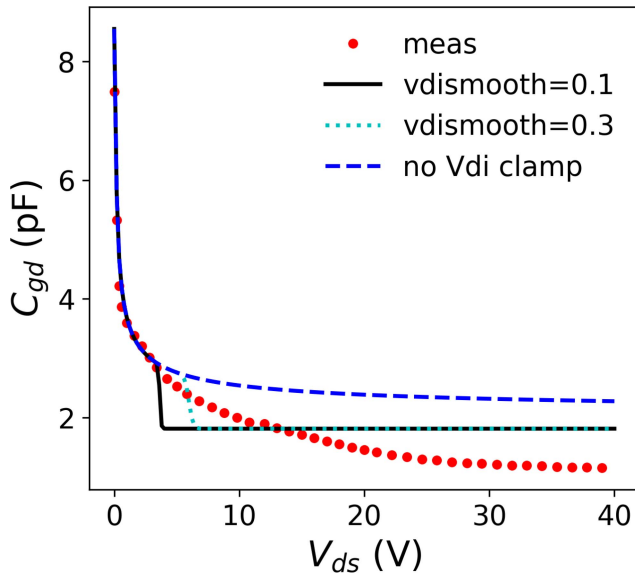
If the capacitances do exhibit spikes, set  $swvdi1c1mp = 1$  and the clamping voltage parameter  $vdi0o$  to a starting value of  $V_{gs,max}$ . Note that data from  $s$ -parameter test structures are needed to extract  $V_{di}$  clamping parameters.

The parameter  $vdislopo$  adjusts the  $V_{ds}$  dependence of the clamped  $V_{di}$ , see Fig. 15(c).  $vdismooth$  smooths the clamping, as Fig. 15(d) shows.

Adjust  $vdi0o$ ,  $vdismooth$ ,  $vdislopo$ , and  $vdislopego$  (the  $V_{gs}$  dependence of the clamped  $V_{di}$ ) to fit the measured  $C_{gg}$  data.

Fig. 16 shows the impact of  $V_{di}$  clamping on  $C_{gd}(V_{ds})$ , where the field plate is modeled as a bias independent plate capacitor. The modeled  $C_{gd}$  drops at the  $V_{di}$  clamping voltage because  $V_{ds}$  can no longer change the intrinsic transistor  $C_{gd}$

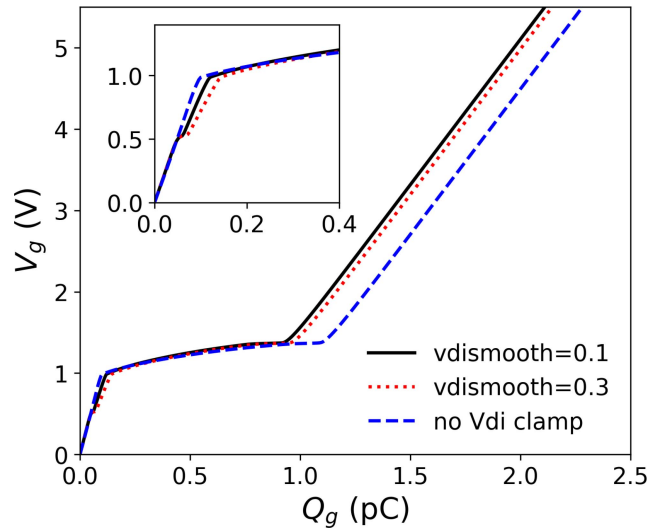




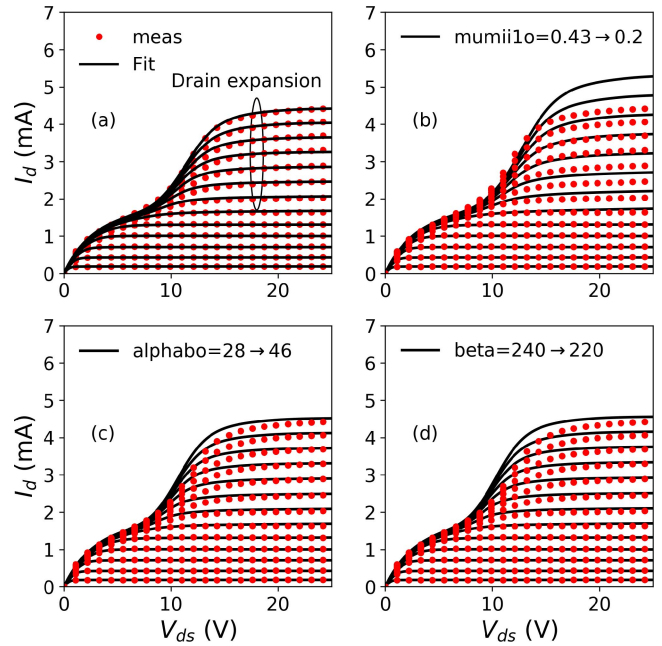
**FIGURE 16.**  $C_{gd}$  vs.  $V_{ds}$ , same transistor as Fig. 14,  $V_{gs} = 0$ ,  $f = 1$  MHz,  $T = 25$  °C.

after  $V_{di}$  is clamped. In reality, the clamping occurs gradually, through the bias dependence of the lateral depletion in the gate-drain overlap region; this phenomena is not modeled physically in PSPHV. A larger  $vdis\ smooth$  helps reduce the abruptness of the drop, see Fig. 16. For this device, the use of  $lgdov2$  readily gives a good fit, as Fig. 14 shows. However, for very high voltage devices, the drop in  $C_{gd}$  can happen even if a MOS, rather than fixed, capacitor is used for the second gate-drain overlap region. For such devices  $V_{di}$  clamping gives more accurate modeling of the total charge  $\int C_{gd} dV_d$ , and therefore more accurate modeling of switching losses.

Fig. 17 shows simulated gate charging characteristics,  $V_{gs}$  vs.  $Q_g$ .  $V_{di}$  clamping reduces the  $Q_g$  at which the transistor turns on, as expected. However, the clamping introduces a small kink at low  $V_{gs}$ , see the inset in Fig. 17. The kink becomes worse as  $vdis\ smooth$  increases, so keep  $vdis\ smooth \leq 0.3$ .



**FIGURE 17.**  $V_{gs}$  vs.  $Q_g$  for a WS transistor,  $W/L/N_g = 30 \mu\text{m}/1 \mu\text{m}/2$ ,  $T = 25$  °C.



**FIGURE 18.**  $I_d(V_{ds})$  for a 65V power n-LDMOS transistor that exhibits strong expansion effect. Bottom to top:  $V_{gs} = 2$  to 8 V by 0.5 V,  $T = 27$  °C.

**VI. EXTRACTION FOR THE DRAIN EXPANSION EFFECT**

PSPHV models the drain expansion effect [19] through modulation of the conductivity of the drift region by impact ionization. Fig. 18(a) shows  $I_d(V_{ds})$  for a transistor with pronounced drain expansion; clearly, PSPHV can model this effect well. The body and source are shorted for this device,  $I_b$  is not directly available, so we fit  $I_d$  where it has a significant component from impact ionization.

Modeling the drain expansion effect requires simultaneous fitting of avalanche current and conductivity modulation by that current.

First, extract the intrinsic transistor avalanche parameters  $a1o$  and  $a2o$  to fit  $I_d(V_{ds})$  at low  $V_{gs}$  and high  $V_{ds}$  (the avalanche current there is dominated by the intrinsic transistor), and to fit the data at moderate  $V_{gs}$ , where the Kirk effect

in the drift region becomes significant, roughly initialize beta and  $\alpha$ .

Second, the parameter  $mumiio$ , which models drift region conductivity modulation due to avalanche current, determines the point at which the expansion effect starts, see Fig. 18(b). Extract  $mumiio$  to align this onset point between PSPHV and the data.

Third, extract  $\alpha$  to fit  $I_d$  at high  $V_{ds}$ . This parameter affects the whole expansion region, see Fig. 18(c). Tune  $vdscaleo$ ,  $thesato$ , and  $thesatgo$  as well to fit  $I_d$  over all  $V_{ds}$ .

Finally, check the slope of  $I_d$  expansion just above the onset of the expansion effect. If the slope in the modeled  $I_d$

is too small, first reduce beta and re-optimize. Fig. 18(d) shows the effect of beta on slope of the expansion effect. If the fitting is still not good enough, also include the second order conductivity modulation parameter `mumii2o`.

## VII. CONCLUSION

We have described a global scaling, self-heating effect aware, large signal parameter extraction method for PSPHV. The parameters important for each extraction step are detailed, and graphical illustrations of the effect of key parameters for effects unique to PSP are provided, these include: drain voltage scaling; the Kirk effect in the drift region; internal drain voltage clamping to suppress spikes in capacitance; and the drain expansion effect. The extraction flow is verified against experimental DC and CV data from -40 °C to 175 °C. Fully scalable and accurate models have been extracted with this method.

## APPENDIX

Below is the description of the used model parameters in alphabetical order. `paramo/l/w/lw` stands for the geometry independent parameter `paramo` and the geometry dependence parameters `paraml`, `paramw`, and `paramlw`. II means impact ionization, DR means drift region.

<code>a1o/l</code>	channel II pre-factor	<code>ecrito/w</code>	velocity saturation critical field
<code>a2o</code>	channel II exponent	<code>fbet1/w</code>	relative mobility decrease due to first lateral profile
<code>a3cvdo/l</code>	drain voltage dependence of <code>a3</code>	<code>fol1</code>	first length dependence coefficient for short channel body effect
<code>a3o/l</code>	saturation-voltage dependence of II	<code>fol2</code>	second length dependence coefficient for short channel body effect
<code>alphabo</code>	DR bottom gate II magnitude	<code>gtho/l/w/a</code>	thermal conductance
<code>aro/l</code>	saturation transition smoothing	<code>jhc</code>	threshold current for DR impact ionization Kirk effect
<code>beta</code>	DR II current exponent for both gates	<code>lov</code>	gate-source overlap length for capacitance
<code>betw1</code>	first higher-order width scaling coefficient of gain factor	<code>lpck</code>	characteristics length of lateral doping profile
<code>cf1</code>	length dependence of DIBL	<code>lp1</code>	mobility-related characteristic length of first lateral profile
<code>cflexp</code>	exponent for length dependence of DIBL	<code>lvaro/l/w</code>	difference between actual and programmed gate length
<code>cgbovl</code>	gate-bulk overlap capacitance per $L_{EN}$	<code>mhc</code>	power for DR II current
<code>clm1l</code>	<code>clm1</code> length dependence coefficient	<code>mueo/w</code>	mobility reduction coefficient
<code>ctgo</code>	gate voltage dependence of interface states factor	<code>mumii1o</code>	1st order DR II modulation effect
<code>ctho/l/w/a</code>	thermal capacitance	<code>mumii2o</code>	2nd order DR II modulation effect
<code>ctlexp</code>	exponent for length dependence of interface states factor	<code>mumto</code>	DR top gate modulation effect on mobility
<code>cto/l/w/lw</code>	interface states factor	<code>novdo</code>	drain overlap region effective doping
<code>dfbo/w</code>	DR bottom gate depletion factor	<code>novo</code>	source overlap region effective doping
<code>dfto</code>	DR top gate depletion factor	<code>npck</code>	pocket doping level
<code>dlq</code>	effective channel length reduction for capacitance	<code>npo</code>	gate poly-silicon doping
<code>dphibo /l/w/lw</code>	offset of $\phi_b$	<code>nsubo</code>	bulk doping
<code>dwq</code>	effective channel width reduction for capacitance	<code>psirto</code>	twice the built-in potential of the DR top gate
<code>ecorno</code>	velocity saturation corner field	<code>rsh0</code>	DR zero-bias sheet resistance
		<code>rsw1</code>	Series resistance
		<code>sta2o</code>	temperature dependence of <code>a2</code>
		<code>stbeto/l/w</code>	temperature dependence of gain factor
		<code>stmueo</code>	temperature dependence of <code>mue</code>
		<code>stthemuo</code>	temperature dependence of <code>themu</code>
		<code>stthesato/w</code>	temperature dependence of <code>thesat</code>
		<code>stvdscaleo/w</code>	temperature dependence of <code>stvdscale</code>
		<code>stvfbo/l/w/lw</code>	temperature dependence of <code>vfb</code>
		<code>stxcoro</code>	temperature dependence of <code>xcor</code>
		<code>tc1o/w</code>	first order temperature coefficient for DR resistance
		<code>tc2o/w</code>	second order temperature coefficient for DR resistance
		<code>tegth</code>	thermal conductance temperature exponent
		<code>themuo</code>	mobility reduction exponent

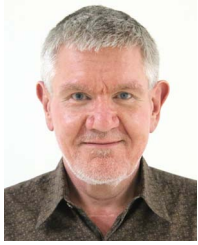
thesatbo	bulk voltage dependence of velocity saturation
thesatgo	gate voltage dependence of velocity saturation
thesatlexp	exponent for length dependence of thesat
thesato/l	temperature dependence of velocity saturation
toxo	gate oxide thickness
toxovd2o	oxide thickness for second gate-drain overlap
toxovdo	oxide thickness thickness for gate-drain overlap
toxovo	oxide thickness thickness for gate-source overlap
uo	zero-field mobility
vdi0o/l	clamping voltage for $V_{di}$
vdislopeo	drain voltage dependence of $V_{di}$ clamping
vdislopgo	gate voltage dependence of $V_{di}$ clamping
vdsmooth	smoothing factor for $V_{di}$ clamping
vdscalebo/l/w/lw	bulk voltage dependence of $V_{ds}$ scaling
vdscalele	exponent for length dependence of $V_{ds}$ scaling
vdscaleo/l/w/lw	$V_{ds}$ scaling factor for NULD
vfbo/l/w/lw	flat-band voltage
vfbov	flat-band voltage for gate-source overlap
vfbovd	flat-band voltage for gate-drain overlap
wvaro/w	difference between actual and programmed field-oxide opening
xbeta	temperature exponent of DR II parameter beta
xcoro	non-universality parameter
xlgdovq	length offset for gate-drain overlap capacitance
xvsato	temperature exponent of DR saturated velocity
xw	DR width offset.

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