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Rigidity Enhancement of GeO₂ by Y-Doping for Reliable Ge Gate Stacks

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ABSTRACT The poor reliability of the GeO₂/Ge stack is improved by appropriate cation doping [e.g., Yttrium (Y)-doping] into amorphous GeO₂ as a result of the enhancement of the GeO₂ network structure stability. In this paper, we discuss the impact of Y-doping on structural modulation of $GeO₂$ on a Ge substrate in thermal treatment. By doping a small amount of Y into amorphous $GeO₂$, the crystallization of GeO₂ to α -quartz and α -cristobalite structures is efficiently suppressed without toughening the local Ge-O bond. This is direct evidence of rigidity enhancement of the $GeO₂$ tetrahedral network structure by cation doping.

INDEX TERMS Germanium, GeO₂, reliability.

I. INTRODUCTION

Germanium has been proposed as a candidate for high performance CMOS channel material due to its high electron and hole mobility values in bulk Ge. Recently, high performance Ge MOSFETs with a $GeO₂/Ge$ gate stack formed by only oxidation have been demonstrated by many groups [\[1\]](#page-4-0), [\[2\]](#page-4-1). These gate stacks have exhibited good initial performance even without forming gas anneal (FGA) but exhibited gate stacks with poor reliability (for example, trap generation by electric stress [\[3\]](#page-4-2), [\[4\]](#page-4-3)), which are very different from $SiO₂/Si$ stacks. In $SiO₂/Si$ stacks, it is well known that FGA is needed to achieve excellent initial properties, and the long-term reliability is excellent. On the other hand, we have demonstrated that the appropriate cation doping into $GeO₂$ substantially improves the long-term reliability in terms of the D_{it} generation at no cost to other initial performance degradations [\[5\]](#page-4-4). Here, the appropriate cation are, for example, Y, Sc or Al which do not form direct bond with substrate Ge. Furthermore, Y is one of the best dopants in terms of leakage current density at metal/doped-GeO₂/Ge stack [\[6\]](#page-4-5).

We have proposed a rigidity model [\[7\]](#page-4-6) to reasonably understand the difference between $SiO₂$ and $GeO₂$ and the impact of Y-doping on $GeO₂$ based on the stability of the amorphous network structure as described in the next section. However, such a doping effect on the structure of thin amorphous $GeO₂$ film, for example, has not been investigated in detail from such a view point yet. In this work, we discuss the reasonability of the rigidity model through physical analyses of local and network structures of amorphous pure $GeO₂$, Y-doped $GeO₂$, and $SiO₂$ films in thermal annealing.

II. RIGIDITY MODEL FOR INITIAL PERFORMANCE AND LONG-TERM RELIABILITY OF GATE STACK

To describe the trade-off relationship between initial performance and long-term reliability based on both the single bond nature and network structure, we have proposed a rigidity model [\[7\]](#page-4-6), which was inspired by the anticipation of insulator/Si interface initial defects based on the constraint theory of network structure [\[8\]](#page-4-7). We simply defined the rigidity as a product of the single bond strength γ and average coordination number N_{av} as the simplest model. As shown in Fig. [1,](#page-1-0) good initial performance and long-term reliability favor low and high γN_{av} , respectively.

In the model, the rigidity of $SiO₂$ is well-balanced and moderate considering the good initial performance and longterm reliability of the $SiO₂/Si$ stack. Hence, the rigidity of $GeO₂$ is lower than that of $SiO₂$ because the Ge-O bond is weaker than the Si-O one [\[9\]](#page-4-8), and rigidity enhancement is needed to improve the long-term reliability of the $GeO₂/Ge$ stack. On the other hand, the doped-Y in the $GeO₂$ network

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FIGURE 1. Initial performance and long-term reliability of MOS gate stack as function of rigidity *γ* **Nav, which is product of single bond strength** *γ* **and average coordination number Nav in rigidity model [\[6\]](#page-4-5). Initial performance and long-term reliability are in trade-off relationship. Comparing moderate (well-balanced) rigidity of SiO2, rigidity of GeO2 is relatively low due to weaker Ge-O bond. In this model, rigidity enhancement is required to improve long-term reliability of GeO2 gate stack.**

is expected to change the network structure with additional Y-O bonds due to a very high N_{av} of 5.6 in crystallized Y2O3. Namely, the improvement of long-term reliability in the Y-doped $GeO₂/Ge$ stack might be mainly due to an increase of N_{av}. The detail of γN_{av} increase by doping and appropriate γN_{av} value for gate stack is semi-quantitatively discussed in ref. 7. From the next section, we experimentally analyze the actual Y-doping impact on local and network structures of $GeO₂$.

III. STABILITY OF AMORPHOUS GEO₂ STRUCTURE

Firstly, we investigated the structural stability of amorphous $GeO₂$ in the thermal annealing process. 270-nm-thick pure $GeO₂$ or Y-doped $GeO₂$ (YGO) films were stacked by RFsputtering of $GeO₂$ or RF-co-sputtering of $GeO₂$ and $Y₂O₃$ on a chemically cleaned Ge(100) substrate, respectively. We also prepared a thermally grown $SiO₂/Si$ stack as a reference. The pure GeO_2 and YGO/Ge and SiO_2/Si stacks were thermally annealed in an ultra-high-vacuum chamber (base pressure: $\sim 10^{-8}$ Pa) at 660°C and 1200°C, respectively. The crystal structure was analyzed by XRD measurement. A Cu Kα X-ray source was used for the XRD measurement.

Figure [2](#page-1-1) shows the typical XRD profiles of thermally annealed pure $GeO₂$ and YGO stacks. The pure $GeO₂$ film clearly exhibited diffraction peaks assigned to the α -quartz and α -cristobalite structures in 5 min and longer annealing, which is in good agreement with previous work [\[10\]](#page-4-9). On the other hand, only the 1 % (cation ratio) doped-YGO film does not exhibit any clear diffraction peaks caused by crystallization even in 15 min annealing. The small diffraction peaks appear again for the 24 % YGO film, but this is due to another instability involved in phase separation,

FIGURE 2. Typical out-of-plane XRD profiles of 270-nm-thick GeO2/ and YGO/Ge stacks annealed at 660◦C. Pure GeO2 is crystallized into *α***-quartz and** *α***-cristobalite structures in annealing, while 1, 6, and 12 % cation ratio YGO do not exhibit clear diffraction peaks. Only 1 % Y-doping is significantly effective to suppress crystallization of GeO2 film. On other hand, 24 % YGO film slightly crystallized into** *α***-quartz structure with highly visible surface roughening. This might be due to another heavy doping instability resulting in spinodal decomposition.**

resulting in decomposition with terrible surface roughening. In addition, the $SiO₂$ film does not exhibit any diffraction peaks caused by crystallization in the 1200◦C annealing (not shown), though the annealing temperature of $1200\degree$ C (1473 K) for $SiO₂$ corresponds to that of 760 $°C$ (1033 K) for $GeO₂$ when the annealing temperature is normalized by the melting point of $SiO₂$ (1983 K) and $GeO₂$ (1388 K) in the absolute temperature scale. It is inferred that the amorphous network structure of $GeO₂$ has been improved a great deal and is like that of $SiO₂$ by the slight amount of Y-doping.

IV. LOCAL BOND STRUCTURE OF GEO2

Secondly, we focused on the local $GeO₂$ structure. We analyzed the 270 -nm-thick amorphous pure $GeO₂$ and YGO (1, 6, and 12 %) films that had been annealed for 1 min at 660° C and the crystallized pure GeO₂ film that had been annealed for 15 min at 660◦C by IR absorption and Raman spectroscopy measurements. The IR measurement was carried out in the vertical incident mode. The absorption peak around 870 cm^{-1} observed in both the amorphous and crystallized $GeO₂$, as shown in Fig. [3\(](#page-2-0)a), is assigned to the asymmetric stretching TO mode of Ge-O-Ge between neighboring GeO⁴ tetrahedral units [\[11\]](#page-4-10). For Raman spectroscopy measurement, an Ar laser ($\lambda = 488$ nm) was used as an optical source. The Raman spectra were observed from 400 to

FIGURE 3. (a) IR absorption spectra and (b) Raman shift spectra of 270-nm-thick amorphous pure GeO2 and YGO in annealing at 660◦C for 1 min and those of crystallized GeO2 in annealing at 600◦C for 15 min. IR absorption peaks around 870 cm−¹ in Fig. [3\(](#page-2-0)a) (blue arrow) and strongest Raman peaks around 430 cm−¹ in Fig. [3\(](#page-2-0)b) (red arrow) are assigned to asymmetric and symmetric stretching mode of Ge-O-Ge bond between neighboring GeO4 tetrahedral units, respectively. Obvious low energy shift of IR peak and high energy shift of Raman peak by crystallization of pure GeO2 are reasonably understandable by narrowing of average Ge-O-Ge bond angle.

 600 cm^{-1} in both the amorphous and crystallized films, as shown in Fig. [3\(](#page-2-0)b). The strongest Raman peak around 430 cm^{-1} was assigned to the symmetric stretching mode of the 6-membered ring [\[12\]](#page-4-11), [\[13\]](#page-4-12) and low frequency edge of the band which originates in the breathing mode in the isolated molecule [\[14\]](#page-4-13).

In pure $GeO₂$, peak energies and peak widths of the IR absorption and the Raman shift are characterized by the most likely Ge-O-Ge bond angle and its distribution. The peak energy of asymmetric stretching (AS) mode (IR absorption peak) and that of symmetric stretching (SS) mode (Raman peak) is clearly shifted down and up by crystallization, respectively. According to the well discussed phonon spectrum in AX_2 tetrahedral amorphous glass [\[14\]](#page-4-13), [\[15\]](#page-4-14),

FIGURE 4. (a) Peak energy and (b) FWHM of IR absorption and Raman peaks as a function of doped Y ratio. The peak energy and FWHM were evaluated by fitting with a single peak. The AS and SS mode peaks do not exhibit obvious modulation by 1 % doping. By increasing Y-doping up to 6 % or higher, the peak energies are shifted in different direction and peak width is increased. The former effect is not explainable by Ge-O bond strength modulation but by slight narrowing of Ge-O-Ge bond angle. The later one is by increasing of bond angle distribution.

those peak shifts are reasonably understandable by the narrowing of the Ge-O-Ge bond angle in neighboring GeO⁴ units [\[14\]](#page-4-13). The most likely Ge-O-Ge bond angle in crystallized $GeO₂$ is speculated to be lower than that in amorphous GeO₂ at 133 \degree [\[16\]](#page-4-15). The narrowing of the peak width in IR absorption and Raman shift corresponds to the tightening of the Ge-O-Ge bond angle distribution, which is reasonably understandable based on the ordering of the local bond structure due to crystallization. From the above results, the crystallization of pure $GeO₂$ is described as the rearrangement and ordering of the local GeO₄ tetrahedral structure.

Next, we analyze the impact of Y-doping on local bond structure in amorphous $GeO₂$. The AS and SS mode peaks do not exhibit clear modulation with 1% Y-doping in terms of the peak energy and width, as shown in Fig. [4](#page-2-1) (a) and (b), which suggests that the local structure of $GeO₄$ units is very little affected by 1 % Y-doping. By increasing the Y-doping ratio to 6 % or higher, the AS and SS mode peak energies are slightly shifted in the opposite direction, and the peak width is increased. According to the $AX₂$ tetrahedral amorphous glass model based on the central force between the A and X atoms [\[14\]](#page-4-13), [\[15\]](#page-4-14), the modulation of the A-X bond stretching constant induces the peak energy shifts of AS and SS mode in the same direction. Therefore, the slight peak shift is not mainly due to Ge-O bond strength modulation but to Ge-O-Ge bond angle narrowing, which is observed in pure $GeO₂$ crystallization. On the other hand, the broadening of the peak width corresponds to the increase of the Ge-O-Ge bond angle distribution. Thus, it is indicated that the significant stabilization of the amorphous $GeO₂$ structure by Y-doping shown in Fig. [1](#page-1-0) is not mainly caused by toughening of the local Ge-O bonds.

FIGURE 5. (a) TDS spectra of ⁷⁴Ge¹⁶O (M/Z = 90) desorption in 270-nm-thick pure GeO₂/ and YGO/Ge structure in annealing at 660°C which **corresponds to GeO desorption rate. By increasing Y-doping density, initial GeO desorption is delayed. (b) Relationship between annealing time at 660◦C and total amount of GeO desorption which is estimated from an integration of TDS signal shown in (a).**

FIGURE 6. The XRD peak area intensity assigned to (101) *α***-quartz GeO2 (observed around 2***θ* **∼ 26◦) as a function of total amount of desorbed GeO shown in Fig. [4\(](#page-2-1)b) It is obvious that YGO does not crystallize in highly GeO desorbed region, where pure GeO2 is already crystallized.**

V. CORRELATION BETWEEN AMORPHOUS STRUCTURAL STABILITY AND GEO DESORPTION

It is known that crystallization of $GeO₂$ is accompanied by GeO desorption [\[10\]](#page-4-9), and Y-doping into $GeO₂$ effectively suppresses the GeO desorption [\[6\]](#page-4-5). Since GeO desorption from $GeO₂/Ge$ stack is triggered by the O vacancy diffusion caused by $GeO₂/Ge$ interface reaction [\[17\]](#page-4-16), it is inferred that the destabilized $GeO₄$ tetrahedron network due to O vacancy diffusion might drive a further structural stabilization leading to the crystallized form and that the doped-Y in $GeO₂$ might stabilize the amorphous structure through perfect suppression of O vacancy diffusion in the thermal annealing. Actually, desorption of sub-oxide is not observed in $SiO₂/Si$ stack in thermal annealing up to 1200◦C at the lowest. Therefore, we also investigate the direct contribution of Y-doping to stabilization of an amorphous structure based on a correlation among Y-doping, O vacancies, and crystallization.

As an indicator of the O vacancy amount in $GeO₂$ and YGO films, we estimate the total amount of desorbed GeO in the annealing process. Fig. [5\(](#page-3-0)a) shows the GeO desorption spectra measured by thermal desorption spectrometer (TDS). Thus, the total amount of desorbed GeO can be evaluated as a function of annealing time by integrating the TDS spectrum. As shown in Fig. [5\(](#page-3-0)b), Y-doping does not perfectly suppress the GeO desorption though the GeO desorption from the YGO/Ge stack starts with a time lag. Fig. [6](#page-3-1) shows the XRD peak intensity assigned to the α -quartz structure of $GeO₂$ as a function of the total amount of desorbed GeO. It is clearly observed that YGO is hard to crystallize compared with pure $GeO₂$ even in the case that the total amount of desorbed GeO is the same. Thus, the result clearly indicates that the doped-Y stabilizes the amorphous structure in addition to delaying the GeO desorption.

On the other hand, the increasing of GeO desorption rate in initial stage in pure $GeO₂/Ge$ is comparable to that in 1%-YGO/Ge, and then the GeO desorption rate from pure $GeO₂/Ge$ obviously decreases with increasing annealing time as shown in Fig. $5(a)$. It suggests that the GeO₂/Ge interface reaction and/or O vacancy diffusion is drastically suppressed with progressing crystallization of $GeO₂$.

VI. DISCUSSION

These experimental results are reconsidered based on the rigidity model. As denoted in Section III, Y-doping is quite effective to suppress crystallization, which is described as a rearrangement of the $GeO₄$ tetrahedral network. This cannot be explained by perfect suppression of GeO desorption but by stabilization of the amorphous network structure, as shown in Section V. Therefore, it directly corresponds to the enhancement of the rigidity by Y-doping.

Furthermore, slight Y-doping into amorphous $GeO₂$ does not strengthen the local bond structure in $GeO₄$ tetrahedral units, as shown in Section IV. Thus, it is concluded that the rigidity enhancement of $GeO₂$ by Y-doping is mainly due to an increase of N_{av} by anchoring the amorphous GeO₄ tetrahedron network frame around the doped-Y as summarized in Fig. [7.](#page-4-17)

It is also reported that Y-doping in bulk germanate glass increases the glass transition temperature and bend strength [\[18\]](#page-4-18), which is expected to be another aspect of the rigidity enhancement. Therefore, understandings of

FIGURE 7. Schematic summary of impact of Y-doping on GeO2 structure. Local bond strength of GeO4 tetrahedral units in amorphous YGO is not different from that in amorphous pure GeO2. However, amorphous network structure stability is different. It clearly appears when local bonds are partially broken by O vacancy diffusion as rearrangement and ordering of GeO4 tetrahedral network. On other hand, Y-doping delays O vacancy diffusion but does not completely suppress it. Therefore, it is speculated that doped-Y may reinforce amorphous cation frame in GeO2.

amorphous network structures control by doping in bulk glass seem to be effectively applicable to thin film in gate stacks, and it is concluded that the control of network rigidity of an insulator by appropriate doping is useful for designing gate stacks of new semiconductor materials as well as those of Ge.

VII. CONCLUSION

Crystallization of GeO₂ on Ge to α -quartz and α -cristobalite structures in thermal treatment is significantly suppressed by a slight amount of Y-doping, while the local bond structure is not significantly modulated by the doping. Therefore, it is understandable that toughening of the amorphous $GeO₂$ network structure occurs as a result of Y-doping. To realize a Ge gate stack with both good initial performance and long-term reliability, the design of a network structure based on the rigidity model will be a key to achieving high performance gate stacks.

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