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# A Continuous Compact DC Model for Dual-Independent-Gate FinFETs

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**ABSTRACT** Multiple-independent-gate (MIG) silicon FinFETs were recently shown capable of enabling: 1) device-level polarity control; 2) dynamic threshold modulation; and 3) subthreshold slope tuning down to ultra-steep-slope operation. These operation mechanisms can unlock a myriad of opportunities in digital as well as analog design. Here we discuss a continuous compact direct-current (dc) model, capable of describing the current–voltage characteristics of a class of MIG FinFETS, namely dual-independent-gate (DIG) FinFETs, over all its biasing design space. This model captures some of the unique features of DIG FinFETs including the dependence of its super-steep subthreshold swing on drain bias and polarity gate bias. An excellent agreement is shown between the model and measured experimental current–voltage characteristics in these devices. Moreover, the predictive nature of the model is evaluated by foreseeing the perspectives of DIG FinFETs as efficient RF detectors at very high frequencies.

**INDEX TERMS** Compact model, FinFET, Schottky barrier, steep subthreshold slope, impact ionization, feedback.

## I. INTRODUCTION

In order to reduce the power consumption in digital integrated circuits, increase the trans-conductance generation efficiency  $(g_m/I_D)$  of transistors in analog circuits, and for attaining a very sensitive nonlinear response to radio-frequency, transistors attaining very steep subthreshold slope (SS) at room-temperature are required. The subthreshold slope in conventional MOSFETs is limited to 60 mV/dec due to the turn-on mechanism for current in such devices being thermionic emission. To overcome this fundamental limit, different types of field-effect transistors, based on alternative current mechanisms, have been proposed during the last decade. In this regard, Tunnel FETs (TFETs) [1]–[5], in which the current turn-on mechanism is band-to-band tunneling rather than thermionic emission, have emerged as an attractive alternative to traditional MOSFETs. Experimental demonstration of such devices showing SS below 30 mV/dec over four decades of current has been recently reported [4]. In addition to TFETs, Dual-Independent-Gate (DIG) FinFETs [6], [7] have been also recently demonstrated capable of achieving a very steep

subthreshold slope (SS << 60 mV/dec at room-temperature). The reason behind this super steep slope is a positive feedback induced by carrier multiplication due to weak impact ionization under proper biasing conditions [6]–[10]. In this regard, experimental demonstrations of DIG FinFETs have shown SS as small as 3.4 mV/dec at room-temperature and the possibility of attaining such small slopes over five decades of current swing [6], [10].

In order to design circuits based on DIG FinFETs and to evaluate and project the performance of these devices for analog and RF applications, it is of interest to develop a simple, closed form, continuous analytical model for its current-voltage characteristics. *This model must: (i) contain the basic physics of weak impact ionization during the ON-transition, and, also, (ii) be capable of capturing the dependence of the current on the voltages applied at all fourterminals of the device.* Here, we report for the first time on a simple and continuous model capable of describing the current-voltage characteristics of DIG FinFETs in all its operation regions. Moreover, the proposed model is capable of capturing all the essential features observed in this device.



**FIGURE 1.** (a) Structure sketch and SEM images of fabricated DIG FinFET, and (b) conceptual band diagram for an *n*-type device illustrating the mechanism of impact-ionization during the ON-transition. Figure adapted from [6].

From the measured *I-V* characteristics of fabricated devices previously reported in [6], we find a set of parameters that enable an excellent fit between the experimental data and the proposed continuous compact model. Moreover, besides enabling circuit design, we discuss how the proposed model can also predict potential future applications of these devices, for instance as efficient RF detectors.

The remainder of this paper is organized as follows. Section II describes the operation of DIG FinFETS. Section III introduces a compact, continuous, model, and describes its peculiarities. Experimental data is fitted to this model showing an excellent agreement. Finally, Section IV discusses the predictive nature and applications of the proposed model.

## **II. OPERATION OF DIG FINFETS**

A sketch of the structure and a SEM image of a DIG FinFET are depicted in Fig. 1a. The device consists of a finshaped silicon channel, metallic source and drain contacts, and two independent gate electrodes. In general, in multipleindependent-gate (MIG) FinFETs, there are two polarity gates, a source-polarity gate (PG<sub>S</sub>) and a drain-polarity gate (PG<sub>D</sub>), which modulate the Schottky barriers at the source and drain ends of the device. Moreover, a control gate (CG) controls the potential barrier in the channel and thus can govern whether the device is in ON or OFF state. In DIG FinFETs, the PG<sub>D</sub> and PG<sub>S</sub> terminals are electrically connected, thus leading to a single polarity gate (PG), as depicted in Fig. 1(a) [6], [7]. The device depicted in Fig. 1(a), which will be discussed herein, has a channel length of  $\sim 600 \, \text{nm}$ (CG length 200 nm); the fin height and width are 340 nm and  $\sim$ 60 nm, respectively [6]. Although fit of experimental data to the proposed model will be performed at the 200 nm CG node, results from numerical simulations employing TCAD Sentaurus predict the validity of the model for shorter channel devices (e.g., 50 nm). Control of carrier injection at the Schottky barrier by the polarity gate, offers the possibility of arbitrarily setting the operation mode of the device as either *n*-type or *p*-type [6], [7]. The polarity gate tunes the Schottky barriers at the source/drain junctions, therefore

selecting the type of majority charge carriers injected in the channel. Although the discussion in this manuscript will be focus on the quadrant where *n*-type operation is achieved, the same functional model, although employing a different set of parameters, could be employed in the quadrant where the device exhibits *p*-type operation.

Depicted in Fig. 1b is the operation mechanism of the device. Application of a positive voltage on the polarity gate creates a potential well under this terminal. When a voltage is applied on the control gate, at the onset of conduction, conducting electrons can acquire enough energy so that electron-hole pairs are generated by impact ionization as indicated by step #1 in Fig. 1b. The electrons drift towards the drain and the generated holes accumulate in the potential well under the control gate as indicted by step #2 in Fig. 1b. This lowers the potential barrier and provides for further electrons for impact ionization. In this process, the more electron-hole pairs that are generated, the lower the potential barrier becomes, thus setting a positive feedback. Another factor that further enhances this carrier multiplication process is modulation of the Schottky barrier by the holes that are swept towards the source, as depicted in step #3 in Fig. 1b. Because of this positive feedback, very steep subthreshold slope ( $SS \ll 60 \text{ mV/dec}$ ) is possible in DIG FinFETs.

## III. MODEL

To physically understand the steep subthreshold in the DIG FinFET I-V characteristics, it is worth noticing that the impact ionization current can be mathematically expressed by [7] and [8]:

$$I(V_{GS}) = I_0 e^{\frac{q_V_{GS}}{nkT}},\tag{1}$$

where:

$$n = (1+r) / \left(1 + r \frac{dV_{BS}}{dV_{GS}}\right),\tag{2}$$

and,  $r = \frac{2\epsilon_{Si} t_{ox}}{\epsilon_{ox} t_{Si}}$ . Here, q is the electron charge, k is Boltzmann's constant, T is temperature;  $t_{ox}$  is the thickness of the oxide layer and  $t_{si}$  is the width of the fin,  $\varepsilon_{ox}$  and  $\varepsilon_{si}$ , are the permittivities of the oxide layer and Si, respectively. Moreover,  $V_{BS}$  is the difference between  $V_{B,CG}$  and the voltage at the source terminal ( $V_S$ ), where the potential  $V_{B,CG}$  is the potential at the middle point between the two sidewalls of the fin in the gate-controlled region. When the carriers generated by impact ionization charge the body, it is possible to have  $dV_{BS}/dV_{GS} > 1$ , and thus n smaller than unity; which leads to SS < 60 mV/dec.

As discussed in [7] and [8],  $dV_{BS}/dV_{GS}$  is given by the following (approximate) expression:

$$\frac{dV_{BS}}{dV_{GS}} \approx \frac{m}{n} \frac{1}{1 + I_{gt}/I_{gi}} \\ \propto \frac{m}{n} e^{\left[\frac{q}{kT}\left(\frac{V_{GS}}{n} - \frac{V_{B,CG} - V_{B,PGS}}{m}\right)\right]} e^{-\frac{\beta_i l}{V_{DS}}}, \quad (3)$$

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where  $V_{B,PGS}$  is the center potential in the source polarity gate region, *m* is an ideality factor in the junction between the region controlled by the control gate and the region controlled by the source-end of the polarity gate,  $I_{gt}$  is the *off*-state leakage current, and  $I_{gi}$  is the impact ionization current [7]. Furthermore, in Eqn. (3),  $\beta_i$  is a constant and *l* is a structural parameter.

In our model, so to properly represent the device characteristics across all bias operation regions, the current voltage characteristics of the DIG FinFET will be expressed by:

$$I = C_T f_{IIF} V_{IF}, \tag{4}$$

where: (i)  $f_{IIF}$  is a dimensionless, bias dependent, factor capable of capturing the steep subthreshold slope in the device, and also capable of smoothly connecting the subthreshold and the above-threshold regions of operation. And, (ii)  $V_{LF}$  is a bias-dependent term that mainly captures the device behavior above threshold. Moreover, (iii)  $C_T$  is a factor that sets the current levels thus is dependent on the Drain-Source (DS) bias as well as the polarity gate (PG) bias. Equation (4) is employed in analogy to the compact analytical model developed for steep-slope TFETs in [11]-[16]. Although the physical mechanisms behind the steep current turn-on mechanism in TFETs and DIG FinFETs are intrinsically different, the functional dependences for current on  $V_{GS}$  during the onset of conduction are rather similar. In both devices the subthreshold current depends exponentially on  $V_{GS}$  and the associated SS can be << 60 mV/dec. In this regard, it is worth mentioning that in [11] and [12] the authors developed a similar compact model for the I-V characteristics of TFETs based on analyzing the device operation in each region of operation and extending the equations ruling the device response in each region to be continuous across all regions of operation.



**FIGURE 2.** Impact ionization factor  $(f_{IIF})$  as a function of  $V_{GS}$  for: (a)  $c_0$  ranging from 10mV to 70mV in 20mV steps (corresponds to variations in SS of approx. 23mV/dec per curve), and  $b_0 = 10^{-7}$ , thus 7 decades variation in drain current; and (b)  $b_0$  ranging from  $10^{-4}$  to  $10^{-7}$ , and  $c_0 = 10$ mV. In both plots  $V_{OFF} = 0.2$ V is employed.

#### A. SUBTHRESHOLD REGION

A characteristic signature of DIG FinFETs is the possibility of attaining (at proper  $V_{DS}$  and  $V_{PG}$  biases) a very steep subthreshold slope, which, as previously discussed, arises from impact ionization processes during the ON-transition. When the bias voltage becomes much smaller than the threshold voltage, experimental results evidence the presence of a constant (i.e.,  $V_{GS}$  independent) leakage current. In order to accurately depict the operation of the device in the subthreshold regime, a compact model should be able to capture the following two effects, namely: (a) leakage current (current level at OFF-state), and (b) steep subthreshold slope at the onset of ON-state, which is dictated by an exponential dependence on  $V_{GS}$  as per Eqn. (1). These two features are simultaneously captured in our model by considering an impact ionization factor, which is given by:

$$f_{IIF} = \left\{ \frac{1}{2} + \left(\frac{1}{2} - b_0\right) tanh\left(\frac{V_{GS} + V_{OFF}}{c_0}\right) \right\}, \quad (5)$$

where: (i)  $b_0$  is a factor setting the ON/OFF current ratio, (*ii*)  $c_0$  is a fitting parameter, which sets the subthreshold slope, and (iii) V<sub>OFF</sub> is an offset voltage. Depicted in Fig. 2a is the dependence of  $f_{IIF}$  on  $c_0$  (for a fixed  $V_{OFF}$  and  $b_0$ ). It is important to mention, that besides being capable of capturing the OFF-state behavior and subthreshold slope, because of the asymptotic limit of hyperbolic tangent to one,  $f_{IIF}$  also works as a smoothing function capable of efficiently connecting the subthreshold and above-threshold regions. The dependence of  $f_{IIF}$  on  $b_0$  is analyzed in Fig. 2b, where  $c_0$  was fixed to  $c_0 = 10$  mV, which corresponds to a  $\sim$ 12 mV/dec subthreshold slope. Since experimental data on DIG FinFETs shows constant, steep, slopes over multiple decades of current swing [6], and due to the fact that in Eqn. (5)  $b_0$  constitutes a factor related to the minimum attainable value by  $f_{IIF}$  (i.e., it can be noticed that  $f_{IIF}$ ranges between 1 -  $b_0$  and  $b_0$  at its asymptotic limits on  $V_{GS}$ ),  $b_0$  will usually be a "small number," i.e.,  $b_0 << 1$ thus the upper bound of  $f_{IIF}$  approaches one. Typical values of  $b_0$ , ranging from  $10^{-4}$  to  $10^{-7}$ , are used in order to show the dependence of  $f_{UF}$  on  $b_0$ , which is depicted in Fig. 2(b). As can be noticed in Fig. 2(b), besides being related to the ON/OFF current ratio,  $b_0$  is also related to the current swing over which maximum SS is observed. These observations are useful for understanding how these parameters are related to the current-voltage characteristics of the device, and for later exploring functional forms for their dependences on  $V_{DS}$  and  $V_{PG}$ . Finally, it is also worth mentioning that this model can be adapted to describe the *p*-type operation of DIG FinFETs simply by just changing the sign of the parameter  $c_0$ .

#### **B. ABOVE-THRESHOLD REGION**

The region above threshold is modeled employing a functional dependence of the form:

$$V_{LF} = \begin{cases} b_1 (V_{GS} + V_{OFF})^{c_1} + 1; & V_{GS} + V_{OFF} > 0\\ 1; & otherwise \end{cases} , \quad (6)$$

where:  $c_1$  and  $b_1$  are fitting parameters setting the order and strength of the current voltage dependence, and  $V_{OFF}$  is an offset voltage. Based on experimental data (see Section III-D), we observe that in general  $1 < c_1 < 2$ , thus although piecewise defined,  $V_{LF}$  results continuous and its derivative is also continuous. The dependence of  $V_{LF}$  on  $b_1$  and  $c_1$  is shown in Fig. 3a and Fig. 3b, respectively. By tuning these two parameters ( $b_1$  and  $c_1$ ) simultaneously, the properties of the device in the region above threshold can be well represented, as will be discussed in Section III-D, by means of fitting the model to experimental data.



**FIGURE 3.**  $V_{LF}$  function plotted (a) for varying  $b_1$  with  $c_1 = 1.6$ ; and (b) for varying  $c_1$  from 1 to 1.9 in steps of 0.3 with  $b_1 = 50$ ; according to Eqn. (6), for both plots  $V_{OFF} = 0.45$  is employed.

## C. IDS-VGS MODEL

After applying the definitions for  $f_{IIF}$  and  $V_{LF}$  given in Eqn. (5) and Eqn. (6), the complete current equation for the DIG FinFET becomes of the form:

$$I(V_{GS}) = C_T \left[ b_1 (V_{GS} + V_{OFF})^{c_1} H(V_{GS}, V_{OFF}) + 1 \right].$$

$$\times \left[ \frac{1}{2} + \left( \frac{1}{2} - b_0 \right) tanh \left( \frac{V_{GS} + V_{OFF}}{c_0} \right) \right].$$
(7)

where  $H(V_{GS}, V_{OFF})$  is the Heaviside step function evaluated at  $V_{GS} - V_{OFF}$ . A typical plot of current as a function of  $V_{GS}$  (assuming  $C_T = 1$ ) is shown in Fig. 4 together with the correspondent plots of  $f_{IIF}$  and  $V_{LF}$  leading to it. We observe, in general, that choosing the same values for the offset voltages in  $f_{IIF}$  and  $V_{LF}$  qualitatively produces a very smooth transition betwen both bias regimes. In the following discussion we will assume this condition.



**FIGURE 4.**  $f_{IIF}$ ,  $V_{LF}$ , and their product, as a function of the control gate voltage ( $V_{GS}$ ). The parameters employed in this calculation are:  $b_0 = 10^{-7}$ ,  $c_0 = 10$  mV,  $V_{OFF} = 0.4$ V,  $b_1 = 2$ ,  $c_1 = 1.1$ .

## D. DEPENDENCE ON V<sub>DS</sub> AND V<sub>PG</sub>

In order to provide for a complete compact model, the current dependences on  $V_{DS}$  and  $V_{PG}$  should be analyzed from this point of view, and, in accordance with Eqn. (7), we will model for the dependences of the unknown parameters, i.e.,  $b_0$ ,  $c_0$ ,  $b_1$ ,  $c_1$ ,  $V_{OFF}$ , and  $C_T$ , on these voltages. Of these parameters, the parameters that have a more clear physical meaning, and which will be discussed in most depth are  $b_0$  and  $c_0$ , which are directly related to the ON/OFF current ratio and SS.



FIGURE 5. Flow chart of the procedure employed to determine the model parameters.

In accordance with Eqns. (2) and (3) it is observed that *n* decreases exponentially with  $V_{DS}$  and  $V_{PG}$ . From Eqn. (2) and (3), the following expression can be derived for the mathematical dependence of  $c_0$  on  $V_{DS}$  and  $V_{PG}$ :

$$c_0 = \frac{\alpha_0}{\left(1 + \beta_0 e^{\gamma_0 V_{PG}} e^{\delta_0 V_{DS}}\right)},\tag{8}$$

where  $\alpha_0$ ,  $\beta_0$ ,  $\gamma_0$ , and  $\delta_0$  are constants. By means of determining  $c_0$  at different values of  $V_{DS}$  and  $V_{PG}$  via fitting the measured  $I_{DS} - V_{GS}$  to the model in Eqn. (7), the values of these unknown constants can be extracted. Similar procedures will be employed to find the voltage dependences of the other five parameters in Eqn. (7). Depicted in Fig. 5 is a sketch of the procedure employed so to extract the complete voltage dependences for all the model parameters. More details about the parameter extraction will be discussed in the next subsection.

## E. PARAMETER EXTRACTION

We extracted the  $b_0$ ,  $c_0$ ,  $b_1$ ,  $c_1$ ,  $C_T$ ,  $V_{OFF}$  from the measured  $I_{DS} - V_{GS}$  characteristics at different values of  $V_{PG}$  and  $V_{DS}$  by means of fitting experimental data to the model in Eqn. (4). The extracted parameters were found by using the Nonlinear Least Squares Error method (NLSE) [17] to minimize the following error function:

$$e_0(b_0, c_0, b_1, c_1, C_T, V_{OFF}) = \sum_{V_{GS}=-1V}^{V_{GS}=1V} \left\{ log_{10} \left( I_{DS, Eqn(7)} \right) - log_{10} \left( I_{DS, measured} \right) \right\}^2,$$

where,  $I_{DS,Eqn.(7)}$ is the calculated current using Eqn. (7), for a fixed  $V_{DS}$  and  $V_{PG}$ , and  $I_{DS,measured}$ is the measured value for the current. Defining:  $Y_1(V_{DS}, V_{PG})$  $b_0(V_{DS}, V_{PG}),$  $Y_2(V_{DS}, V_{PG})$ =  $c_0(V_{DS}, V_{PG}), \quad Y_3(V_{DS}, V_{PG}) = b_1(V_{DS}, V_{PG}), \quad Y_4(V_{DS}, V_{PG}), \quad Y_4(V_{DS}, V_{PG}) = b_1(V_{DS}, V_{PG}), \quad Y_5(V_{DS}, V_{PG}) = b_1(V_{DS}, V_{PG}) = b_1($  $V_{PG}$ ) =  $c_1(V_{DS}, V_{PG})$ ,  $Y_5(V_{DS}, V_{PG})$  =  $C_T(V_{DS}, V_{PG})$ , and  $Y_6(V_{DS}, V_{PG}) = V_{OFF}(V_{DS}, V_{PG})$ , 3D plots in the  $[V_{DS}, V_{PG}]$  space were then generated for these six functions  $(Y_i, i = 1, \dots, 6)$ . From these 3D plots, the values for the unknown parameters in the  $Y_i$  functions were found by means of fitting the extracted values to the mathematical expressions in Eqns. (8)-(13). Fitting was performed via employing a non nonlinear least squares error method to minimize the following error functions:

$$e_{i} = \sum_{VDS} \sum_{VPG} \left\{ Y_{i,calc,Eqn.(n)} - Y_{i,extracted} \right\}^{2},$$

where  $Y_{i,calc,Eqn.(n)}$  are the calculated values for the  $Y_i$  functions at various  $V_{DS}$  and  $V_{PG}$  using Eqns. (8)-(13) and  $Y_{i,extracted}$  are the extracted values.

Depicted in Fig. 6 are the extracted values of  $c_0$  versus  $V_{PG}$  (Fig. 6(a)) and  $V_{DS}$  (Fig. 6(b)) together with the curves obtained from the best fit of this data to the model in Eqn. (8). Overall, a good qualitative agreement is observed between the physical model in Eqn. (8) and the extracted values of  $c_0$  obtained from fitting the measured device data to the model in Eqn. (7).

At this end, we proceeded to plot the extracted values of  $b_0$  versus  $V_{PG}$  (Fig. 7(a)) and  $V_{DS}$  (Fig. 7(b)). It is observed that these extracted values, in logarithmic scale, show a linear dependence with  $V_{PG}$  as well as  $V_{DS}$ . Moreover,  $b_0$  increases with  $V_{DS}$ , and decreases with  $V_{PG}$ . As discussed in Fig. 2, a smaller  $b_0$  is an indication of a smaller leakage current floor thus a larger ON/OFF current ratio. The fact that  $b_0$  increases with  $V_{DS}$  is an evidence of Drain-Induced Barrier Lowering (DIBL) effect [18], which leads to an increased OFF-current leakage current floor as  $V_{DS}$  is increased. In contrast, the lower  $b_0$  at larger  $V_{PG}$  is a result of a better conduction through the polarity-gate controlled regions, i.e., higher polarity-gate induced doping, which physically translates into a larger ON-current level. Based on the above observations, we employed a model of the form:

$$b_0 = \alpha_1 e^{\beta_1 V_{DS} + \beta_2 V_{PG} + \beta_3 V_{DS} V_{PG} + \beta_4}, \tag{9}$$

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in order to model the  $V_{DS}$  and  $V_{PG}$  dependences in  $b_0$ . Depicted in Fig. 7(a)-(b) are the curves obtained from fitting the extracted values of  $b_0$  to the model in Eqn. (9). A good qualitative agreement is observed between the model in Eqn. (9) and the extracted values of  $b_0$  obtained from fitting the measured device data to the model in Eqn. (7).



**FIGURE 6.** Extracted and fitting of  $c_0$  versus (a)  $V_{PG}$ , and (b)  $V_{DS}$ . Fitting was performed employing the model in Eqn. (8).



**FIGURE 7.** Extracted and fitting of  $b_0$  versus (a)  $V_{PG}$ , and (b)  $V_{DS}$ . Fitting was performed employing the model in Eqn. (9).

In addition, we proceeded to find models for  $b_1$ ,  $c_1$ ,  $C_T$ , and  $V_{OFF}$ . For this purpose we looked at the 3D plots of these parameters in the  $[V_{DS}, V_{PG}]$  space and determine adequate functions for fitting these points with a continuous surface. Since the physical origin of these parameters is less clear than that for  $b_0$  and  $c_0$ , we employed two-variable polynomials and exponential functions in our fitting models. In this regard, for  $b_1$  and  $c_1$  we will employ a 3<sup>rd</sup> and 2<sup>nd</sup> order quadratic equations, respectively, of the form:

$$b_{1} = p_{0} + p_{1}V_{DS} + p_{2}V_{PG} + p_{3}V_{DS}^{2} + p_{4}V_{DS}V_{PG} + p_{5}V_{PG}^{2} + p_{6}V_{DS}^{3} + p_{7}V_{DS}^{2}V_{PG} + p_{8}V_{DS}V_{PG}^{2} + p_{9}V_{PG}^{3},$$
(10)  
$$c_{1} = q_{0} + q_{1}V_{DS} + q_{2}V_{PG} + q_{3}V_{DS}V_{PG} + q_{4}V_{PG}^{2} + q_{5}V_{DS}^{2},$$
(11)

Moreover, for  $C_T$  we will employ a combination of a linear dependence on  $V_{PG}$  and an exponential dependence on  $V_{DS}$  and  $V_{PG}$ , according to the following equation:

$$C_T = \mu_0 + \mu_1 V_{PG} + \mu_2 \ e^{\sigma_0 V_{PG} + \sigma_1 V_{DS} + \sigma_2 V_{PG} V_{DS}}, \quad (12)$$

And, for  $V_{OFF}$  we will employ polynomial dependences on  $V_{PG}$  as well as on  $V_{DS}$  of the form:

$$V_{OFF} = r_0 + r_1 V_{DS} + r_2 V_{PG} + r_3 V_{DS}^2 + r_4 V_{DS} V_{PG} + r_5 V_{PG}^2 + r_6 V_{DS}^3 + r_7 V_{DS}^2 V_{PG} + r_8 V_{DS} V_{PG}^2 + r_9 V_{PG}^3.$$
 (13)

Depicted in Fig. 8 are 3D plots of the extracted values of  $b_1$  (Fig. 8(a)),  $c_1$  (Fig. 8(b)),  $C_T$  (Fig. 8(c)), and  $V_{OFF}$  (Fig. 8(d)), versus  $V_{PG}$  and  $V_{DS}$ , together with the surface plots obtained from fitting these extracted values to the models in Eqns. (10)-(13). A good qualitative agreement is observed between these models and the extracted values of the parameters obtained from fitting the measured device data to the model in Eqn. (7).



**FIGURE 8.** 3D plots of the extracted values of: (a)  $b_1$ , (b)  $c_1$ , (c)  $C_7$ , and (d)  $V_{OFF}$ , versus  $V_{PG}$  and  $V_{DS}$ , together with the surfaces fitting these parameters to the models in Eqns. (10)-(13).

Depicted in Fig. 9 are examples of the measured currentvoltage characteristics for the analyzed DIG-FinFET, whose performance was reported in [6], together with the calculated current-voltage characteristics obtained via fitting to the continuous compact model here discussed in semi-logarithmic (Fig. 9(a)-(b)) and linear scales (Fig. 9(c)-(d)), showing a good agreement between the model and measurements.

## F. IDS-VDS MODEL

Our discussion so far has been focused on the device operating in "saturation" region and on the  $I_{DS} - V_{GS}$  characteristics. When analyzing the device characteristics on

the  $V_{DS}$  domain, it is observed that "saturation" occurs for  $V_{DS} > 2V$ . Otherwise, when  $V_{DS} < 2V$ , the current was found to follow a dependence on  $V_{DS}$  of the form:

$$I \propto \left(1 - e^{-\frac{V_{DS}B}{A}}\right). \tag{14}$$



**FIGURE 9.** Experimental (circles) and modeled  $I - V_{GS}$  characteristic for a DIG FinFET. (a)  $V_{PG} = 5V$  and  $V_{DS} = 2$ , 3, 4, 5V (logarithmic-scale). (b)  $V_{DS} = 3V$  and  $V_{PG} = 3$ , 4, 5V (logarithmic-scale). (c)  $V_{PG} = 5V$  and  $V_{DS} = 2$ , 3, 4, 5V (linear-scale). (b)  $V_{DS} = 3V$  and  $V_{PG} = 3$ , 4, 5V (linear-scale). (c)  $V_{PG} = 5V$  and  $V_{DS} = 3V$  and  $V_{PG} = 3$ , 4, 5V (linear-scale).

So to capture this behavior, we can define a modified current equation as follows:

$$I_{DS} = \left\{ w \ I_{Eqn.(7)} + (1 - w) \ I_{V_{DS} \approx 2V} \right\} \left( 1 - e^{-\frac{V_{DS}^B}{A}} \right).$$
(15)

where, A and B were found to be functions of  $V_{GS}$  and  $V_{PG}$ . The equations employed for representing A and B are of the form:

$$A = s_0 + s_1 V_{GS} + s_2 V_{PG} + s_3 V_{PG}^2 + s_4 V_{GS} V_{PG} + s_5 V_{GS}^2 + s_6 V_{GS}^3 + s_7 V_{GS} V_{PG}^2 + s_8 V_{GS}^2 V_{PG},$$
(16)  
$$B = t_0 + t_1 V_{PG} + t_2 V_{GS} + t_3 V_{GS} V_{PG} + t_4 V_{PG}^2.$$
(17)

Moreover:

$$w = \left\{ \frac{1}{2} + \frac{1}{2} \tanh\left(\frac{V_{DS} - 2V}{50mV}\right) \right\},\tag{18}$$

was used to smoothly connect the "linear" and "saturation" regions.

Depicted in Fig. 10 are examples of the measured  $I_{DS} - V_{DS}$  characteristics and their fitting using the

proposed model. Results are displayed in semi-logarithmic and linear scale showing good agreement in both cases.

Finally, Table 1 summarizes all the parameters used to fit

#### TABLE 1. Fitting parameters.



**FIGURE 10.** Experimental (circles) and modeled *I-V* characteristic for a DIG FinFET, with  $V_{PG} = 5V$  and  $V_{GS} = 0.1, 0.3, 0.5V$ . (a) In semi-logarithmic scale. (b) In linear scale.



**FIGURE 11.** Transconductance  $(g_m)$ : (a) directly extracted from the numerical derivative of the measured device data, and (b) evaluated from the model.

## **IV. PREDICTIVE NATURE**

In order to evaluate the potential of electron devices for analog and RF applications, accurate estimates for the current-voltage characteristics as well as for its derivatives are required. For instance, compact models for TFETs enabled the evaluation of the perspectives of these devices for analog applications [13] as well as RF detectors [19]. In this regard, it is worth mentioning that numerical calculation of the derivatives directly from experimental device data can be prone to significant uncertainty when computing derivatives. From this point of view fitting of experimental data to a continuous compact model, whose derivatives are also continuous, can be therefore beneficial. Depicted in Fig. 11 is a typical plot of the trans-conductance  $(g_m)$ as directly extracted from the numerical derivative of the measured device data, and as evaluated from the model depicting this effect. From this point of view fitting of experimental data to a continuous compact model, whose derivatives are also continuous, can be therefore beneficial.

PARAMETER	PARAMETER	PARAMETER	PARAMETER
	VALUE		VALUE
α <sub>0</sub>	0.05507 V	r <sub>0</sub>	-0.3548 V
$\beta_0$	2.348 x 10 <sup>-8</sup>	$\mathbf{r}_1$	0.0783
γo	2038 V <sup>-1</sup>	$\mathbf{r}_2$	0.3568
$\delta_0$	$2.092 \text{ V}^{-1}$	r <sub>3</sub>	0.07945 V <sup>-1</sup>
$\alpha_1$	1.0 x 10 <sup>-6</sup>	$r_4$	-0.1935 V <sup>-1</sup>
$\beta_{I}$	1.367 V <sup>-1</sup>	<b>r</b> 5	-0.0317 V <sup>-1</sup>
$\beta_2$	-0.04912 V <sup>-1</sup>	r <sub>6</sub>	-0.03445 V <sup>-2</sup>
$\beta_3$	-0.1259 V <sup>-2</sup>	$\mathbf{r}_7$	0.0539 V <sup>-2</sup>
$\beta_4$	-0.6986	$r_8$	-0.004 V <sup>-2</sup>
$\mathbf{p}_0$	-23.25	r9	0.00285 V <sup>-2</sup>
$\mathbf{p}_1$	-45.02 V <sup>-1</sup>	$S_0$	0.09588
$p_2$	75.4 V <sup>-1</sup>	$\mathbf{s}_1$	0.2357 V <sup>-1</sup>
p <sub>3</sub>	5.9 V <sup>-2</sup>	<b>S</b> <sub>2</sub>	-0.05913 V <sup>-1</sup>
$p_4$	14.47 V <sup>-2</sup>	<b>S</b> <sub>3</sub>	0.009813 V <sup>-2</sup>
<b>p</b> <sub>5</sub>	-33.98 V <sup>-2</sup>	S4	-0.1246 V <sup>-2</sup>
$\mathbf{p}_6$	2.667 V <sup>-3</sup>	<b>S</b> 5	-0.2963 V <sup>-2</sup>
p <sub>7</sub>	-6.18 V <sup>-3</sup>	<b>S</b> <sub>6</sub>	-0.3516 V <sup>-3</sup>
$p_8$	1.96 V <sup>-3</sup>	<b>S</b> <sub>7</sub>	0.01688 V <sup>-3</sup>
p9	3.815 V <sup>-3</sup>	$\mathbf{S}_8$	0.2009 V
$q_0$	-1.115	t <sub>0</sub>	3.276
$\mathbf{q}_1$	0.8633 V <sup>-1</sup>	t <sub>1</sub>	$0.1658 \text{ V}^{-1}$
q <sub>2</sub>	$0.5276 \text{ V}^{-1}$	$t_2$	-1.7 V <sup>-1</sup>
$q_3$	-0.1 V <sup>-2</sup>	t <sub>3</sub>	0.2 V <sup>-2</sup>
$q_4$	-0.01571 V <sup>-2</sup>	$t_4$	-0.04583 V <sup>-2</sup>
q <sub>5</sub>	-0.06 V <sup>-2</sup>		
$\mu_0$	2.939 A		
$\mu_1$	0.03424 A/V		
$\mu_2$	3.86 x 10 <sup>-5</sup> A		
$\sigma_0$	$0.8955 V^{-1}$		
$\sigma_1$	1.076 V <sup>-1</sup>		
<b>σ</b> 2	0.001335 V <sup>-2</sup>		

In this section, we will discuss on the application of DIG-FinFETs as RF detectors at millimeter-wave and terahertz frequencies and predict the performance of these devices based on the developed model. Many detection approaches have been developed to-date [20], however, a miniaturized room-temperature detector technology, exhibiting a very sensitive broadband response at room-temperature, is still largely missing. Conventional FETs can operate as efficient detectors far beyond their fundamental cut-off frequency [21]. Indeed, FETs can intrinsically operate as non-resonant broadband detectors with high responsivity, because of self-mixing [22] as described by the Dyakonov-Shur theory [23].

When a RF field is applied between the gate and the source terminal of the transistor, this electric field is rectified, and a source-to-drain DC photo-current, or photo-voltage, dependent on the device configuration and boundary conditions, is induced. Assuming current detection and conjugate matching of an antenna coupling the incoming RF signal into the device so to provide maximum power transfer, the following upper bound for the detector responsivity can be obtained [24]–[26]:

$$R_i = \frac{1}{2\sigma_0} \frac{\partial \sigma_0}{\partial V_{GS}} = \frac{1}{2} \frac{\partial \ln I_{DS}(V_{GS})}{\partial V_{GS}} = \frac{1}{2} \frac{g_m}{I_{DS}}, \qquad (19)$$

where  $\sigma_0$  corresponds to the FET channel conductivity. Moreover,  $g_m$  in Eqn. (19) represents the transistor trans-conductance, and  $g_m/I_{DS}$  is the trans-conductance to current ratio, which is an important factor, related to the slope in the current-voltage characteristics, and widely employed in analog circuit design. In traditional FETs, the  $g_m/I_{DS}$  ratio is theoretically limited to values below 38.5 V<sup>-1</sup>, which sets a maximum attainable responsivity. The ratio between the responsivity in DIG FinFET detectors and the maximum attainable responsivity in traditional FETs is thus given by:

$$E = \frac{R_{i, DIG \ FinFET}}{R_{i, \max MOSFET}} = \frac{kT}{q} \frac{\partial \ln I_{DS} (V_{GS})}{\partial V_{GS}} = \frac{kT}{q} \frac{g_m}{I_{DS}}.$$
 (20)

Plotted in Fig. 12 is the responsivity enhancement factor (E), as defined in Eqn. (20), for a DIG FinFET detector configuration where the polarity-gate voltage is set to  $V_{PG} = 5$ V. E is plotted as a function of  $V_{GS}$  in order to depict the  $V_{GS}$ bias dependence of the response. Moreover, two cases are analyzed, corresponding to  $V_{DS}$  biases of: (i)  $V_{DS} = 5V$ , where impact ionization is significant and the device shows maximum  $SS \ll 60 \text{ mV/dec}$ , and (ii)  $V_{DS} = 2 \text{V}$ , where the SS is always > 60 mV/dec. We observe that whereas in the first case, in the subthreshold region much larger responsivity than that in traditional FET detectors is attainable, in the second case, the responsivity is always smaller than the FETlimit. This is a result of the capability of attaining a smaller SS under large  $V_{PG}$  and  $V_{DS}$  biases because of impact ionization processes. In general, based on Eqns. (7) and (20), it is observed that:

$$E_{max} = \frac{R_{i,maxDIGFinFET}}{R_{i,maxMOSFET}} = \frac{kT}{q} \frac{2}{c_0}.$$
 (21)

Therefore, the smallest  $c_0$ , the largest the detector responsivity. As a general rule,  $c_0$  will benefit from: (a) large applied voltages ( $V_{PG}$  and  $V_{DS}$ ), thus large slopes in the electric field, (b) small polarity-gate lengths, and (c) from the incorporation of direct-gap materials with smaller bandgap than that of Si at the drain-end of the device so to boost carrier multiplication processes as carriers drift from the control-gate-controlled channel to the drain contact.



**FIGURE 12.** Estimated current-responsivity enhancement (*E*) as a function of  $V_{GS}$  bias for two different drain biases for a DIG FinFET configured as RF detector ( $V_{PG} = 5V$ ). The RF responsivity is estimated from the developed DC compact model of the device.

#### **V. CONCLUSION**

We discussed an analytic DIG FinFET continuous compact model capable of capturing the current–voltage characteristics of the DIG FinFET in all the device operation regions. Although our discussion is focused only with regards to *n*type operation, the model can be readily adapted to also describe the *p*-type operation. The model accounts for the bias-dependent super steep subthreshold swing, leakage levels, as well as the above-threshold characteristics of the device. To prove the validity of the model, the model was fitted to experimental results on fabricated devices. An excellent agreement was observed. Moreover, based on the proposed model, we predict on the performance of DIG FinFETs as RF detectors.

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