

# China's First Infrared Occultation Spectrometer (GaoFen-5/AIUS) Monitoring Atmosphere Above Antarctic: Evaluation of Methane Profiles

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**Abstract**—The atmospheric infrared ultraspectral sounder (AIUS) is China's first infrared occultation spectrometer onboard the GaoFen-5 satellite. This instrument provides measurements in the near-infrared and mid-infrared spectral regions with a spectral resolution of  $0.02 \text{ cm}^{-1}$ , monitoring key greenhouse gases like methane ( $\text{CH}_4$ ). In this study, we have developed a physics-based retrieval algorithm for accurately deriving  $\text{CH}_4$  mole fraction profiles from infrared absorption spectra. The retrieval algorithm is applied to the AIUS and successfully obtains  $\text{CH}_4$  mole fraction profiles from January to December 2019 over Antarctic. The validation results show that the retrieval algorithm has a satisfactory performance using AIUS measurements, with average relative standard deviation of the differences less than 7.2% throughout the altitude range of 10–70 km.

**Index Terms**—Antarctic, GaoFen-5 (GF-5)/atmospheric infrared ultraspectral sounder (AIUS), methane ( $\text{CH}_4$ ), retrieval algorithm.

## I. INTRODUCTION

METHANE ( $\text{CH}_4$ ) is a chemically and radiatively active gas in the atmosphere and is the second most important

Manuscript received 30 October 2023; revised 23 January 2024 and 25 February 2024; accepted 21 March 2024. Date of publication 27 March 2024; date of current version 12 April 2024. This work was supported in part by the Chinese Academy of Sciences (CAS) "Pioneering Initiative Talents Program" under Grant E1RC2WB2, in part by the National Key Research and Development Program of China under Grant 2022YFF0606404, and in part by Basic Strengthening Program Technical Field Foundation under Grant 2023KJC-Y-0193. (Corresponding authors: Xiaoying Li; Jian Xu.)

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Digital Object Identifier 10.1109/JSTARS.2024.3381651

anthropogenic greenhouse gas after carbon dioxide ( $\text{CO}_2$ ) [1]. Before the 1970s, scientists measured atmospheric  $\text{CH}_4$  concentrations by analyzing air trapped in polar ice cores. Direct measurements of  $\text{CH}_4$  in the troposphere began in 1978 [2], [3]. Currently,  $\text{CH}_4$  concentrations can be obtained through discrete air samples collected at the ground surface, as well as through remotely sensed measurements of atmospheric  $\text{CH}_4$  columns obtained from the ground surface or space. According to data from the National Oceanic and Atmospheric Administration, the global annual concentration of  $\text{CH}_4$  has been constantly changing since 1983. These data indicate a sustained increase in atmospheric  $\text{CH}_4$  concentrations in the 1980s, a stable period between 1999 and 2005, and generally rapid growth after 2005. In 2022, the concentration of  $\text{CH}_4$  reached  $1911.83 \pm 0.59$  ppb [4].

Compared to traditional methods based on ground-based observations or air sampling, satellite measurements have several advantages, including stability, continuity, large-scale coverage, and convenient access to globally spatiotemporal distribution. Satellite remote sensing has become a major tool over the past decades for monitoring greenhouse gases, such as  $\text{CH}_4$  and  $\text{CO}_2$ . In 1979, the stratospheric and mesospheric sounder on the Nimbus-7 satellite measured the global distribution of stratospheric  $\text{CH}_4$  concentration for the first time [5]. Afterward, several satellite sensors have been launched to measure the vertical column densities and concentration profiles of  $\text{CH}_4$ . The  $\text{CH}_4$  column densities can be retrieved from near-infrared radiance measured by instruments, such as the scanning imaging absorption spectrometer for atmospheric cartography [6] and the tropospheric monitoring instrument [7]. The instruments, such as the Michelson interferometer for passive atmospheric sounding [8], the atmospheric infrared sounder [9], the atmospheric chemistry experiment—Fourier transform spectrometer (ACE-FTS) [10], and the tropospheric emission sounder [11] can observe  $\text{CH}_4$  profiles based on thermal infrared emission. The thermal and near-infrared sensor for Carbon Observation Fourier Transform Spectrometer-2 [12] can retrieve the  $\text{CH}_4$  columns and profiles from SWIR spectra and TIR spectra, respectively [13].

GaoFen-5 (GF-5) is a high-resolution remote sensing satellite that was launched on May 9, 2018 (Beijing local time). The satellite was designed to fly in a high-inclination ( $98.2^\circ$ ) circular low-Earth orbit, located 705 km above the Earth's surface, with

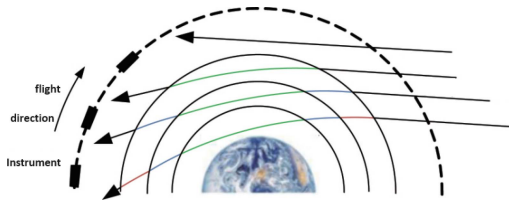


Fig. 1. Illustration of the AIUS solar occultation observation [24].

the local time of ascending node at 13:00 [14]. GF-5 is equipped with six payloads designed to achieve various scientific objectives and has a projected lifespan of eight years [15]. However, the GaoFen-5 satellite lost contact with the ground station and has been out of service since April 2020. The atmospheric infrared ultraspectral sounder (AIUS) on board the GF-5 satellite, designed by the Beijing Institute of Space Mechanics and Electricity, is China's first solar occultation infrared spectrometer. The AIUS was designed to detect  $O_3$  and other atmospheric compositions related to ozone depletion, with a specific focus on the atmosphere above Antarctic [16], [17]. The main scientific objectives of the AIUS are to provide high-precision observations for atmospheric environmental monitoring and climate change research, and to promote the development of China's occultation/limb measurement technology [17].

Presently, only a few studies on AIUS trace gas retrievals have been conducted. Wang et al. [18] used simulated spectra to retrieve the profiles of  $N_2O$ ,  $NO_2$ , and HF in order to validate the suitability of the OEM algorithm using the parameters of GF-5 AIUS. Cao et al. [19] applied a channel selection algorithm to the AIUS and retrieved temperature profiles from the AIUS. Li et al. [16], [20] evaluated the obtained AIUS Level-1 data and performed the first retrievals of  $O_3$ ,  $H_2O$ , and HCl profiles from AIUS spectra.  $CH_4$  is an important greenhouse gas and its spatiotemporal variations over Antarctic strongly associate with global warming and climate change.

In addition to the AIUS, the ACE-FTS is the only instrument that is currently available for measuring  $CH_4$  profiles over Antarctic. The AIUS is similar to the ACE-FTS in terms of observation geometry and can serve as a complement to the ACE-FTS. It is an important data source for monitoring  $CH_4$  in polar regions.

Currently, the AIUS have demonstrated its potential for retrieving  $O_3$ ,  $H_2O$ , HCl, and other gases. Considering the role of  $CH_4$  in global warming and climate change [21], [22], the vertical distribution of  $CH_4$  over Antarctic has not yet been adequately observed. Therefore, this article focuses on the AIUS  $CH_4$  inversion. Here, we develop a retrieval algorithm based on the radiative transfer calculations and optimal estimation for deriving  $CH_4$  profiles from AIUS spectra during the whole year of 2019 and analysis the spatiotemporal variations over Antarctic. This study will, for the first time, focus on retrieving  $CH_4$  profiles using AIUS measurements and expand the application scope of the first domestic infrared occultation satellite.

The rest of this article is organized as follows. In Section II, the measurements and methodology are introduced, followed by a brief description of the AIUS instrument and Level-1 data,

the retrieval method, and the channel selection. The results of AIUS  $CH_4$  retrieval over the Antarctic, the error analysis, and the comparison with the ACE-FTS  $CH_4$  data are discussed in Section III. Finally, Section IV concludes this article.

## II. MEASUREMENTS AND METHODOLOGY

### A. Overview of the AIUS Instrument and Level-1 Data

During sunrise in orbit, the AIUS measured the solar transmittance spectrum of the tangential atmosphere from the horizon to the outer atmosphere (as shown in Fig. 1) within a bandwidth of  $750$  to  $4100\text{ cm}^{-1}$  along the altitude from  $8$  to  $100\text{ km}$ . It carried out this measurement with a high spectral resolution of  $0.02\text{ cm}^{-1}$ . The pressure, temperature, and concentrations of atmospheric species (such as  $O_3$ ,  $H_2O$ , HCl,  $NO_2$ ,  $N_2O$ , HF, etc.) were retrieved from the specific spectral absorption lines of different atmospheric components. Its latitudinal coverage ranged from  $55^\circ\text{S}$  to  $90^\circ\text{S}$ .

The AIUS was implemented with two types of photovoltaic detectors: Mercury Cadmium Telluride (MCT) and Indium Antimonide (InSb). The spectral calibration precision met the requirements for atmospheric trace gas inversion [23]. Their spectral ranges were  $750\text{--}1850\text{ cm}^{-1}$  and  $1850\text{--}4100\text{ cm}^{-1}$ , respectively. The AIUS utilized ultraresolution Fourier transform spectrometer technology and sun-tracking system technology to provide observations with a high spectral resolution [24].

Converting the AIUS Level-0 data (raw measurements) to Level-1 data (irradiance spectra) required a series of processing and calibration steps, including data preparation, interferogram reconstruction, spectrum computation, spectral calibration, and transmittance calculation [16], [24].

The first step was to acquire and process the auxiliary data. The Level-0 data consisted of interferometric measurements downloaded from the spacecraft, along with the acquisition time, sun position, and satellite position. These parameters were used to calculate the geometric parameters, altitude, and geographic location (latitude and longitude).

The AIUS operated in low Earth orbit. The interferogram experienced spikes due to the influence of charged particles. The second step was to calculate the interference values at the spikes in order to reconstruct the interferogram using interpolation or statistical methods.

In the third step, the fast Fourier transform was used to calculate the spectrum from the reconstructed interferogram. Accurate calibration is an essential technology for obtaining precise measurements of  $O_3$ ,  $CO_2$ , and  $CH_4$  concentrations. The fourth step involved spectral calibration, which included correcting for Doppler shift, selecting suitable lines, and accurately determining peak positions [24].

The final step was to calculate the transmittance. The Level-1 data product includes transmittance, instrument line shape (ILS), as shown in Fig. 2, as well as observation geometry and other processed auxiliary data. The relationship between transmittance and the digital number is shown as follows [16]:

$$\tau(h, \lambda) = \frac{D(h, \lambda) - B(h, \lambda)}{S(h, \lambda) - B(h, \lambda)} \quad (1)$$

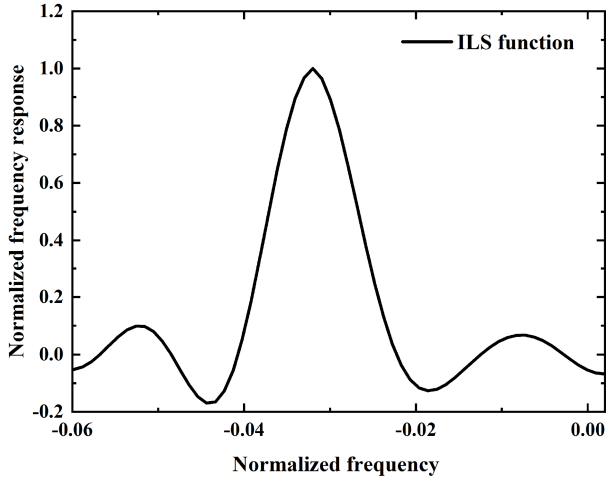


Fig. 2. ILS function of the AIUS.

where  $h$  represents the tangent;  $\lambda$  denotes the wavenumber; and  $D(h, \lambda)$ ,  $B(h, \lambda)$ , and  $S(h, \lambda)$  refer to the digital number value, the deep space signal, and the solar radiation outside the atmosphere, respectively.

### B. Retrieval Method

The retrieval method comprises a forward model based on the reference forward model (RFM) and an inversion procedure based on the optimal estimation method (OEM). Ignoring scattering and assuming local thermodynamic equilibrium, the radiative transfer equation for the radiation intensity at wavenumber  $\nu$  received by an instrument at position  $s$  is as follows [25], [26]:

$$I(\nu, s) = I(\nu, s_0)e^{-\tau(\nu; s_0, s)} + \int_{s_0}^s B(\nu, T(s'))e^{-\tau(\nu; s', s)}\alpha(\nu, s')ds' \quad (2)$$

where  $I(\nu, s_0)$  represents the background contribution at wavenumber  $\nu$  and position  $s_0$ ,  $\tau$  denotes the optical depth,  $B(\nu, T)$  signifies the radiance emitted by a black body at temperature  $T$ , and  $\alpha(\nu, s')$  represents the volume absorption coefficient from position  $s_0$  to position  $s$ . On the right side of (2), the first term represents the attenuated radiation, while the second term represents the increase in radiance caused by atmospheric thermal emission. In this equation,  $B(\nu, T)$  is represented by the Planck function, shown as follows:

$$B(\nu, T) = \frac{2hc^2\nu^3}{e^{h\nu/k_B T} - 1} \quad (3)$$

where  $h$  represents the Planck constant and  $k_B$  denotes the Boltzmann constant, respectively. The transmittance calculated by the radiation transfer model is shown as follows [27]:

$$\begin{aligned} \tau &= e^{-\sum \chi} \\ &= e^{-\sum k\mu} \\ &= e^{-\sum k(\nu, p, T, e) \int \nu \rho ds} \end{aligned} \quad (4)$$

where  $\chi$  is the optical thickness,  $k$  is the absorption coefficient, which is in general a function of the wavenumber  $\nu$ , pressure  $p$ , temperature  $T$ , and partial pressure  $e$ .  $\mu$  is the absorber amount, which is related to the volume mixing ratio (VMR) of the absorber  $\nu$ , the (molecular) air density  $\rho$ , and the distance along the path  $s$ .

We typically used a discrete version of the nonlinear data model for the radiative transfer equation. Equation (2) can be replaced with

$$\mathbf{y} = \mathbf{F}(\mathbf{x}, \mathbf{b}) + \boldsymbol{\epsilon} \quad (5)$$

where  $\mathbf{y}$  represents the measurement vector, which refers to the measurement of transmittance [see (1)]; and  $\mathbf{F}$  denotes the radiative transfer model, which represents the simulated transmittance [see (4)]. In this study, the RFM [27], [28] is used for line-by-line radiative transfer calculations.  $\mathbf{x}$  refers to the state vector;  $\mathbf{b}$  represents the forward model parameters; and  $\boldsymbol{\epsilon}$  denotes the vector of measurement noise.

The retrieval method utilizes the OEM to find a solution estimate that minimizes a cost function, which can be written as [29]

$$\begin{aligned} \chi^2 &= (\mathbf{y} - \mathbf{F}(\mathbf{x}, \mathbf{b}))^T \mathbf{S}_\epsilon^{-1} (\mathbf{y} - \mathbf{F}(\mathbf{x}, \mathbf{b})) \\ &\quad + (\mathbf{x} - \mathbf{x}_a)^T \mathbf{S}_a^{-1} (\mathbf{x} - \mathbf{x}_a) \end{aligned} \quad (6)$$

where  $\mathbf{x}_a$  represents a priori profile,  $\mathbf{S}_a$  denotes a priori covariance matrix,  $E|(\mathbf{x} - \mathbf{x}_a)(\mathbf{x} - \mathbf{x}_a)^T|$ ,  $\mathbf{S}_\epsilon$  signifies the covariance matrix for measurement errors,  $\mathbf{K}\mathbf{S}_a\mathbf{K}^T$ , and  $\mathbf{K}$  represents the Jacobian matrix,  $\frac{\partial \mathbf{F}(\mathbf{x})}{\partial \mathbf{x}}$ .

Newtonian iteration is a straightforward numerical technique to find the root of the gradient of (6) when the problem is not highly nonlinear. For the general vector equation,  $\mathbf{g}(\mathbf{x}) = 0$ . The function  $\mathbf{g}(\mathbf{x})$  is the derivative of the cost function (6) and  $\nabla_{\mathbf{x}}\mathbf{g}$  is the second derivative. The iteration is similar to Newton's method for the scalar case and can be expressed as

$$\mathbf{x}_{i+1} = \mathbf{x}_i - [\nabla_{\mathbf{x}}\mathbf{g}(\mathbf{x}_i)]^{-1}\mathbf{g}(\mathbf{x}_i). \quad (7)$$

Calculate  $\mathbf{g}(\mathbf{x})$  and  $\nabla_{\mathbf{x}}\mathbf{g}$  based on (6), and substitute into the Newtonian iteration (7). In the Gauss–Newton method, the next iteration is given by

$$\begin{aligned} \mathbf{x}_{i+1} &= \mathbf{x}_i + (\mathbf{S}_a^{-1} + \mathbf{K}_i^T \mathbf{S}_\epsilon^{-1} \mathbf{K}_i)^{-1} \\ &\quad (\mathbf{K}_i^T \mathbf{S}_\epsilon^{-1} (\mathbf{y} - \mathbf{F}(\mathbf{x}_i)) - \mathbf{S}_a^{-1} (\mathbf{x}_i - \mathbf{x}_a)). \end{aligned} \quad (8)$$

When the true solution is significantly far from the current iteration point, the Newton iteration will be invalid. In the retrieval of  $\text{O}_3$ ,  $\text{H}_2\text{O}$ , and  $\text{HCl}$  profiles from AIUS measurements, Li et al. [16] used the Levenberg–Marquardt (LM) method to solve the underlying least-squares fitting problem. In the LM method, the cost function is minimized by introducing a constraint factor, denoted as  $\gamma$ . A new iterate is as follows:

$$\begin{aligned} \mathbf{x}_{i+1} &= \mathbf{x}_i + ((1 + \gamma) \mathbf{S}_a^{-1} + \mathbf{K}_i^T \mathbf{S}_\epsilon^{-1} \mathbf{K}_i)^{-1} \\ &\quad (\mathbf{K}_i^T \mathbf{S}_\epsilon^{-1} (\mathbf{y} - \mathbf{F}(\mathbf{x}_i)) - \mathbf{S}_a^{-1} (\mathbf{x}_i - \mathbf{x}_a)) \end{aligned} \quad (9)$$

where  $\gamma$  represents a smoothing factor. When  $\gamma \rightarrow 0$ , (8) and (9) are similar. For more details on the retrieval procedure, refer to [16], [29], [30], [31], [32], [33].

TABLE I  
MODEL PARAMETERS AND EXTERNAL DATA SOURCES

Atmospheric parameters	Description
Temperature profile	MLS retrievals
Pressure profile	MLS retrievals
Interfering species	MLS, ACE-FTS and AFGL
CH <sub>4</sub>	Li et al. [16]
Spectroscopic parameters	Description
Spectroscopic data	HITRAN2016
Instrument parameters	Description
Viewing geometry	Occultation
Tangent Height	the tangent height of AIUS
Instrument line shape	see Fig. 2
Field of view	1.25 mrad
Observer Altitude	705 KM

The error analysis is essential for retrieving the atmospheric composition. There are three types of errors that can be characterized by their respective error covariance: the covariance of the smoothing error ( $\mathbf{S}_s$ ), the covariance of the forward model error ( $\mathbf{S}_p$ , which includes spectroscopy, temperature, and pressure error), and the covariance of the retrieval random noise error ( $\mathbf{S}_m$ ) [29], given as follows:

$$\mathbf{S}_s = (\mathbf{A} - \mathbf{I})\mathbf{S}_a(\mathbf{A} - \mathbf{I})^T \quad (10)$$

$$\mathbf{S}_p = \mathbf{G}\mathbf{K}_b\mathbf{S}_b\mathbf{K}_b^T\mathbf{G}^T \quad (11)$$

$$\mathbf{S}_m = \mathbf{G}\mathbf{S}_\epsilon\mathbf{G}^T \quad (12)$$

$$e = \sqrt{\mathbf{S}_s + \mathbf{S}_p + \mathbf{S}_m} \quad (13)$$

where  $\mathbf{A}$  represents the Averaging Kernel (AK) matrix,  $\mathbf{A} = \mathbf{G}\mathbf{K}$ ;  $\mathbf{I}$  denotes the identity matrix;  $\mathbf{G}$  refers to the gain matrix,  $\mathbf{G} = (\mathbf{K}^T\mathbf{S}_\epsilon^{-1}\mathbf{K} + \mathbf{S}_a^{-1})^{-1}\mathbf{K}^T\mathbf{S}_\epsilon^{-1}$ ;  $\mathbf{S}_b$  represents the error covariance matrix of  $b$ ; and  $\mathbf{K}_b$  denotes the Jacobian matrix with respect to the parameters of the forward model. The calculation of matrices  $\mathbf{G}$  and  $\mathbf{A}$  is realized by coding, whereas  $\mathbf{K}$  is obtained from the forward model. The smoothing error ( $e_s$ ), the forward model error ( $e_p$ ), and the retrieval random noise error ( $e_m$ ) can be represented by the square roots of  $\mathbf{S}_s$ ,  $\mathbf{S}_p$ , and  $\mathbf{S}_m$ , respectively. The total error ( $e$ ), is calculated as the square root of the sum of the aforementioned covariance.

The parameters in the forward model include the CH<sub>4</sub> retrieval channel, spectral resolution, the absorbing species, atmospheric profiles, viewing geometry, spectroscopic data, ILS, and the field of view [27] (see Table I). Within the wavelength range used in this experiment, there are no significant differences between the CH<sub>4</sub> versions of HITRAN 2020 and HITRAN 2016. The spectroscopic parameters for CH<sub>4</sub>, N<sub>2</sub>O, H<sub>2</sub>O, and other species were taken from the HITRAN 2016 database [34].

The interfering species profiles were constructed using atmospheric models from the Air Force Geophysics Laboratory, Microwave Limb Sounder (MLS) Level-2 products, and the ACE-FTS Climatology Version 4.1. The atmospheric profiles cover the altitude range from 0 to 102 km with a vertical grid of 1, 2, and 3 km in the ranges of 0–50 km, 50–90 km, and above 90 km, respectively. The a priori information for the CH<sub>4</sub> VMR was constructed following the methodology outlined by Li et al. [16].

### C. Channel Selection

The AIUS has 167 501 spectral bands that contain a significant amount of information but cannot be directly used for CH<sub>4</sub> retrieval. High spectral resolution measurements allow for the separation of individual bands. We use microwindows to enhance computational efficiency and reduce errors caused by interfering species [35].

First, we analyzed the absorption of atmospheric molecules in the 750–4100 cm<sup>-1</sup> range using data on the atmospheric molecular absorption spectrum from the HITRAN website.<sup>1</sup> Combining the four spectral features with the significant absorption points of CH<sub>4</sub> in the infrared spectrum (2913 cm<sup>-1</sup>, 1533.3 cm<sup>-1</sup>, 3018.9 cm<sup>-1</sup>, and 1305.9 cm<sup>-1</sup>) and considering the absorption of interfering components, we initially screened three broad spectral absorption ranges: 1130–1660 cm<sup>-1</sup>, 2470–3160 cm<sup>-1</sup>, and 3820–3910 cm<sup>-1</sup>. And then, based on the absorption windows of potential interfering components in the radiative transfer model (refer to Table II), we selected seven initial inversion spectral windows for CH<sub>4</sub>, namely 1200–1230 cm<sup>-1</sup>, 1330–1340 cm<sup>-1</sup>, 1380–1500 cm<sup>-1</sup>, 1550–1660 cm<sup>-1</sup>, 2645–2980 cm<sup>-1</sup>, 3820–3822 cm<sup>-1</sup>, and 3865–3920 cm<sup>-1</sup>. These windows were selected to minimize interference from other atmospheric species.

The seven aforementioned microwindows still contain 33 107 spectral bands. Using these bands for retrieval is still inefficient in terms of data management and computational resources. Finally, We selected the retrieved channels based on the measure of information entropy. The channel selection method mainly includes the following four main stages.

*Step 1:* Calculate the background error covariance matrix ( $\mathbf{S}_a$ ).

*Step 2:* Calculate the observed covariance matrix ( $\hat{\mathbf{S}}$ ). The observation covariance matrix can be written as [36]

$$\hat{\mathbf{S}} = \mathbf{S}_a - \mathbf{S}_a\mathbf{K}^T(\mathbf{K}\mathbf{S}_a\mathbf{K}^T + \mathbf{S}_\epsilon)^{-1}\mathbf{K}\mathbf{S}_a. \quad (14)$$

The third step is to calculate the information capacity ( $\mathbf{H}$ ). If  $\mathbf{S}_a$  and  $\hat{\mathbf{S}}$  are known, the information capacity included in the observation process is recorded as follows:

$$\mathbf{H} = \frac{\ln|\mathbf{S}_a| - \ln|\hat{\mathbf{S}}|}{2}. \quad (15)$$

The final step is to obtain the channels through iterative processes. At each iteration, only the channel with the highest value of  $\mathbf{H}$  is selected.

After conducting an analysis on disturbance in atmospheric components and performing experiments to channel selection, we were able to obtain 120 spectral bands for CH<sub>4</sub> retrieval. These bands are shown in Fig. 3. On most CH<sub>4</sub> microwindows, the spectra were of high quality with a signal-to-noise ratio (SNR) of 200–350 [16].

We compared the AIUS CH<sub>4</sub> retrieval bands used in this experiment with the officially announced ACE-FTS CH<sub>4</sub> retrieval bands, as shown in Fig. 3. The gray line represents the simulated transmittance, while the blue dot denotes the ACE-FTS

<sup>1</sup>[Online]. Available: <https://hitran.org/>

TABLE II  
FITTING WINDOWS

atmospheric parameter	sensitivity wavelength ( $\text{cm}^{-1}$ )		
Temperature	1300, 1525	-	3850
Pressure	1250 – 1300	3000 – 3100	-
CO <sub>2</sub>	1230 – 1270, 1340 – 1380	2480 – 2530, 2580 – 2645	3830
H <sub>2</sub> O	1500 – 1550	2980 – 3160	-
N <sub>2</sub> O	1140 – 1200, 1240 – 1330	2470 – 2490, 2525 – 2600	3822 – 3835, 3840 – 3860
O <sub>3</sub>	1130 – 1170	3000 – 3060	3825 – 3850, 3855 – 3865

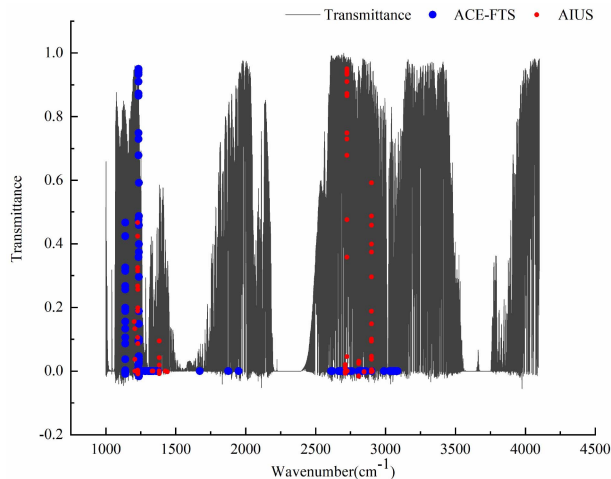


Fig. 3. Comparison of the AIUS and ACE-FTS CH<sub>4</sub> retrieval channels.

CH<sub>4</sub> retrieval bands, and the red dot indicates the AIUS CH<sub>4</sub> retrieval bands. The CH<sub>4</sub> channel selection results from this experiment are primarily included in the ACE-FTS CH<sub>4</sub> inversion channel, except for the spectral range of 1500–2000  $\text{cm}^{-1}$ . ACE-FTS CH<sub>4</sub> experiments use a spectral range spanning from 1500 to 2000  $\text{cm}^{-1}$  to enhance CH<sub>4</sub> retrieval. However, in our experiment, we excluded this microwindow to eliminate H<sub>2</sub>O interference on the retrieval performance.

CH<sub>4</sub> retrieval experiments were conducted using two different channel schemes, as illustrated in Fig. 4. The CH<sub>4</sub> concentration profiles retrieved using two different channel schemes show a similar variation trend with height, but the retrieval efficiency of the AIUS retrieval scheme 1 (using 120 channels) is 24 times higher than that of the channel scheme 2 (using 1115 channels).

Fig. 5 illustrates the weight function values of the 120 selected channels in the order of their selection. The horizontal axis represents the selected channels, the vertical axis represents the tangent height, and the color bar represents the weight. The 120 selected channels covered the altitude range from the troposphere to the mesosphere, and the tangent heights at which the extreme values of each channel weight function varied sequentially.

### III. RESULTS AND DISCUSSION

#### A. AIUS CH<sub>4</sub> Retrievals

Unfortunately, the AIUS instrument has been out of service since April 2020. The spectral measurements after 2019 were lost due to malfunctions at the ground station. To obtain a more

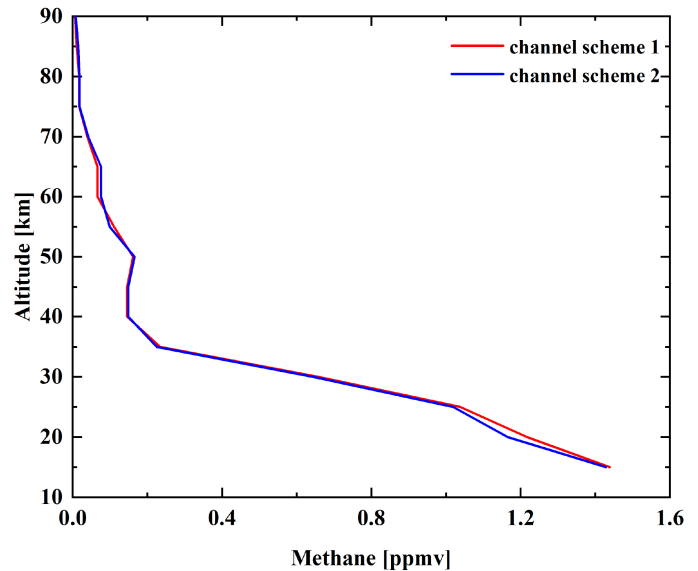


Fig. 4. CH<sub>4</sub> profiles were retrieved from two different channel schemes. The retrieval was carried for AIUS measurements recorded on May 16, 2019. The red line represents the CH<sub>4</sub> profile retrieved from the 120 channels, while the blue line refers to the CH<sub>4</sub> profile retrieved from 1115 channels.

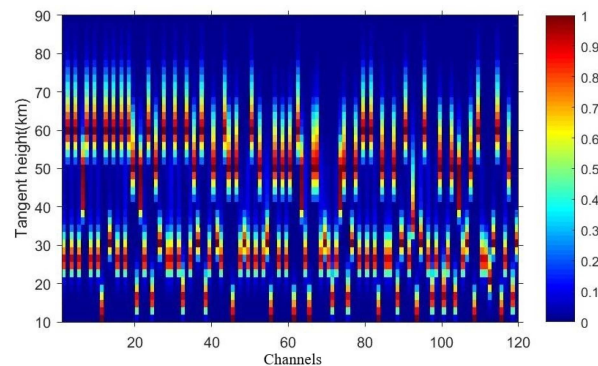


Fig. 5. Weight function matrix of AIUS channels.

solid outcome, we have included Level-1 data from the commissioning phase (from January to March 2019) for additional retrieval experiments. In this study, we used 246 orbits of AIUS spectra observed in different months for retrieval. Table III presents the temporal and spatial characteristics of the AIUS data we used in this experiment, which were categorized into five latitude zones: 60–65°S, 65–70°S, 70–75°S, 75–80°S, and 80–85°S.

The value of the averaging kernel (AK) represents the contribution of the measurement data in the retrieval. The AK shape describes the vertical resolution of the retrieved profile [37]. The

TABLE III  
INFORMATION ON THE AIUS DATA USED IN THIS RESEARCH

Latitude Zone	Month
60 – 65°S	Jan, Nov, Dec
65 – 70°S	Jan, May, Jun, Jul, Nov
70 – 75°S	Feb, May, Aug
75 – 80°S	Feb, Apr, Aug, Oct
80 – 85°S	Mar, Apr, Sep, Oct

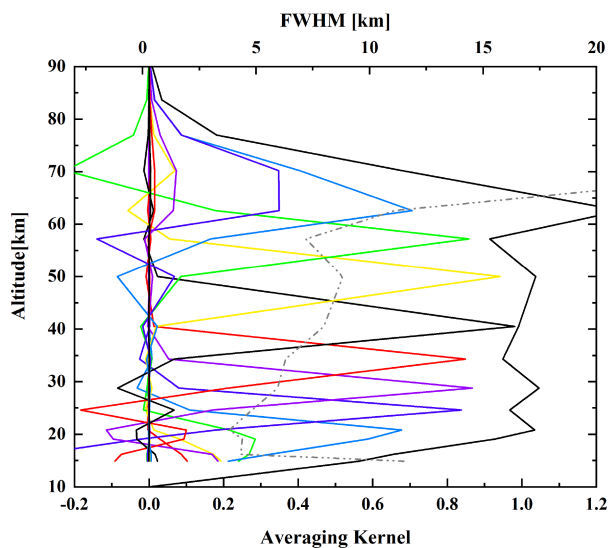


Fig. 6. Averaging kernel of a single  $\text{CH}_4$  profile. The retrieval was carried for AIUS measurements recorded on May 16, 2019. The colorful solid line represents the averaging kernel, while the gray line indicates the FWHM. The black solid line expresses the sum of the average kernel at each altitude.

AK and the full-width-at-half-maximum (FWHM) of AIUS data on May 16, 2019 are shown in Fig. 6. When the altitude is below 15 km or above 65 km, the AK value is less than 0.3. When the altitude is between 15 and 65 km, the AK value falls within the range of 0.7–1. It indicates that the information on  $\text{CH}_4$  retrieval mainly comes from the AIUS measurement data in this experiment. The FWHM reaches 4–11 km across the altitude range from 10 to 65 km. Above 65 km, the FWHM exceeds 20 km.

The residuals between the measured spectra and the simulated spectra are shown in Fig. 7. The spectral residuals between the simulated and measured spectra are within  $\pm 0.2$  and  $\pm 0.03$  at the tangent heights of 40.5 and 70.2 km, respectively. This indicates that the residual transmittance of the measured data is lower at the higher tangent height compared to the simulated spectra.

We calculated the monthly average value of the  $\text{CH}_4$  profiles, as shown in Fig. 8. It illustrates that  $\text{CH}_4$  concentrations show significant seasonal variations over Antarctica. In each season, the monthly mean  $\text{CH}_4$  profile is similar during consecutive months, such as April–May, July–August, and November–December.

The monthly average VMR in November and December is higher than that in other months for altitudes ranging from 21 to 76 km. These two months are the warm season in Antarctica, during which there may be an increase in the release of  $\text{CH}_4$  from ocean sediment and permafrost. This can result in an increase in the  $\text{CH}_4$  VMR in the atmosphere [38]. The average VMR in

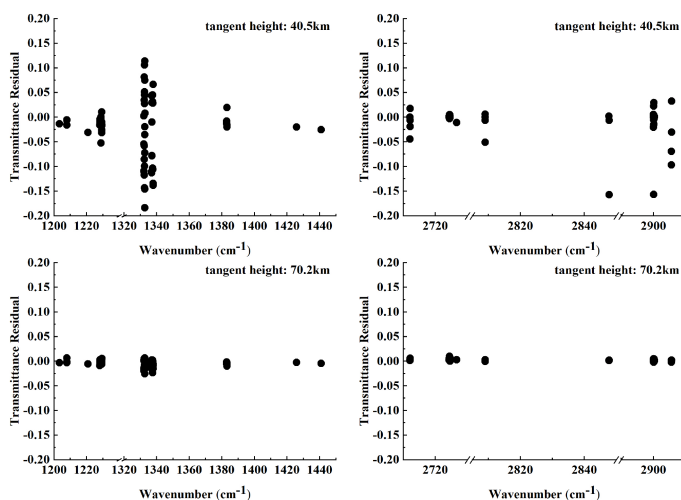


Fig. 7. Comparison of residual transmittance between measured and simulated spectra. The AIUS measurements were recorded on May 16, 2019. Four spectra are plotted for tangent heights at 40.5 and 70.2 km, respectively.

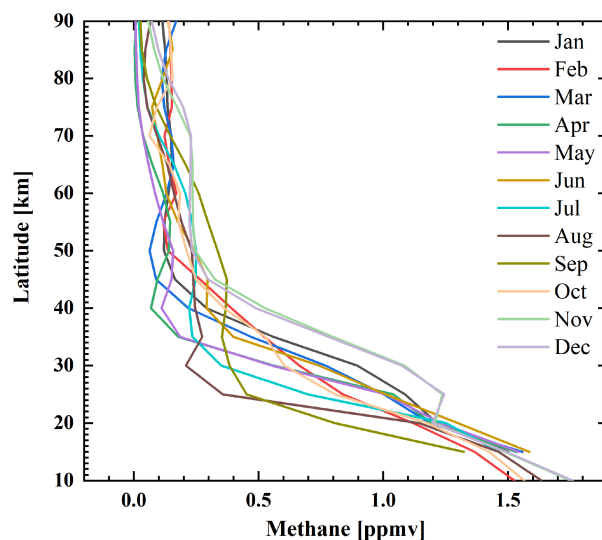


Fig. 8. Comparison of the AIUS  $\text{CH}_4$  profiles in different months.

April and May is lower than in other months for altitudes above 57 km. Below 57 km, as altitude increases, the monthly average VMRs in September, August, April, and March are lower than those in other months. The VMR value of  $\text{CH}_4$  decreases as altitude increases. The  $\text{CH}_4$  VMR is approximately 1.17–1.65 ppmv between 10 and 20 km, and decreases to less than 0.2 ppmv at 35 km in April and May.

Due to the characteristics of AIUS measurements, it is impossible to observe the same latitude zone every month, and we cannot obtain measurement data for the same latitude zone each month. The aforementioned analysis results may be influenced by latitude. So, we also analyzed the average  $\text{CH}_4$  profiles in different latitude zones, as shown in Fig. 9.

Fig. 9 illustrates that  $\text{CH}_4$  concentrations show significant latitude variations over Antarctica. The average VMR in 60 – 65° S is higher than that in other months, ranging from 21 to 90 km. The VMR of  $\text{CH}_4$  decreases with increasing altitude.

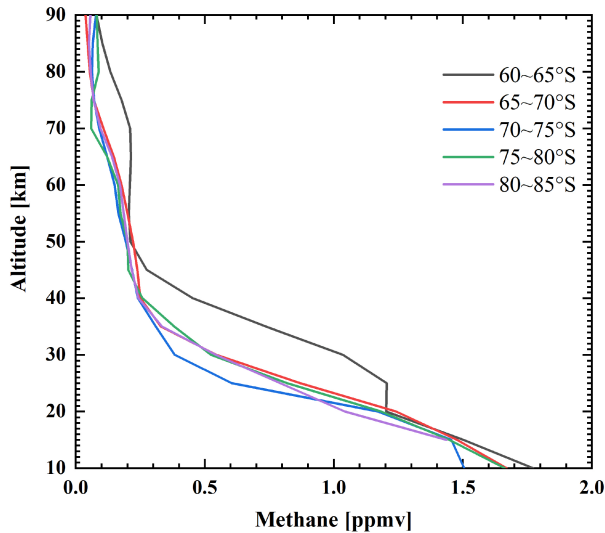


Fig. 9. Comparison of the AIUS CH<sub>4</sub> profiles in different latitude zones.

TABLE IV  
MODEL PARAMETER ERRORS CONSIDERED IN THE AIUS CH<sub>4</sub> RETRIEVALS.

Atmospheric parameters	Perturbation
Temperature	1K
Pressure	1%
Line strength	
CH <sub>4</sub>	4%
H <sub>2</sub> O	1%
CO <sub>2</sub>	5%
N <sub>2</sub> O	1%
O <sub>3</sub>	1%

As can be seen from Figs. 8 and 9, certain profiles lack numerical values at the 10-km mark. The tangent height measured by the AIUS is sometimes higher than 10 km, which makes it impossible to retrieve the CH<sub>4</sub> concentration at this altitude.

### B. Error Characterization

In the retrieval experiment, it is necessary to assess the uncertainties of atmospheric parameters, such as temperature and pressure, among others, that are used in the experiment. In this study, we consider the potential sources of error, including smoothing error, forward model error, and random noise error. The sources of error, and the corresponding perturbation amounts are shown in the Table IV.

The precision of the MLS version 4 temperature and pressure measurements is better than 1 K and 1% from the upper troposphere to the stratosphere [39], [40], [41], respectively. In this study, we found errors caused by the uncertainties in temperature and pressure to be 4.82E-04 ppmv (0.15%), and 2.70E-03 ppmv (0.44%), respectively.

The CH<sub>4</sub> intensities in HITRAN 2016 [34] are derived from the HITRAN 2012 and the McCaSDa database below 1370 cm<sup>-1</sup>, and they remain unchanged from the HITRAN 2012 intensities in the 1370–4000 cm<sup>-1</sup> region. The uncertainty of CH<sub>4</sub> line intensities used in this experiment is smaller than 4% [42], [43], [44], which can cause a bias of 6.17E-03 ppmv (1.27%). The uncertainty of H<sub>2</sub>O and CO<sub>2</sub> line intensity is

1% and 5% [45], respectively. Unfortunately, we did not find official information on the uncertainty of N<sub>2</sub>O and O<sub>3</sub> line strength within the AIUS CH<sub>4</sub> retrieval channels; we have chosen to apply a 1% perturbation as a conservative estimate of the actual uncertainty. The uncertainties in line intensity of the four atmospheric components mentioned previously (H<sub>2</sub>O, CO<sub>2</sub>, N<sub>2</sub>O, and O<sub>3</sub>) contribute to a deviation of 1.38E-05 ppmv (1.90E-03%).

Individual estimates of various errors and relative errors for the CH<sub>4</sub> retrieval are shown in Fig. 10. Below 25 km, the total error is mainly attributed to random noise and smoothing error, while above 25 km, the total error is primarily due to spectroscopy and random noise error. The uncertainties in temperature and pressure have a relatively minor impact on the total error.

### C. Validation Using ACE-FTS Observations

The ACE-FTS is a solar occultation measurement instrument on board SCISAT-1, which was launched in August 2003. The satellite is in a circular orbit at an altitude of 650 km and an inclination angle of 74° [46]. The ACE-FTS covers the 750–4400 cm<sup>-1</sup> range with a high spectral resolution of 0.02 cm<sup>-1</sup>. It provides vertical profiles of trace gas VMRs and temperature across latitudes from 85° S to 85°N. The vertical resolution is approximately 4 km from the cloud tops up to about 150 km [47].

The retrieval of ACE-FTS CH<sub>4</sub> products from atmospheric spectra involves two steps. The first step is deriving the pressure and temperature profiles using the spectral lines of carbon dioxide. Second, CH<sub>4</sub> VMR profiles are retrieved using a modified global fitting approach, in which all parameters are determined with the LM nonlinear least-squares method [48]. The retrieval product is not sensitive to the initial profiles and does not include averaging kernels. The uncertainties reported in the data files are the statistical fitting errors resulting from the least-squares process and do not account for systematic components or parameter correlations [48], [49].

The HITRAN 2016 spectroscopic line parameters and the Voigt line shape were used in the ACE-FTS V4.1 retrieval calculations, employing a similar methodology to that used for AIUS retrievals. More detailed explanations of the retrieval algorithms can be found in [48], [50], and [51].

The accuracy of the version v2.2 CH<sub>4</sub> data is within 10% from the upper troposphere to the lower stratosphere, and within 25% in the middle and higher stratosphere up to the lower mesosphere (<60 km) when compared with correlative satellite, balloon-borne, and ground-based data [10]. In Version 4 (with version 4.1), which represents the most recent update, the retrieval utilizes the latest spectroscopic information and features improved accuracy in forward model calculations [51]. In this article, we compared ACE-FTS CH<sub>4</sub> data V4.1 with AIUS data. The data are available online.<sup>2</sup>

For the comparison experiment, we limited the spatial coincidence to within 1.5° (approximately 150 km) and the temporal coincidence to within 36 h between AIUS and ACE-FTS. As a

<sup>2</sup>[Online]. Available: [https://database.scisat.ca/level2/ace\\_v4.1\\_v4.2/](https://database.scisat.ca/level2/ace_v4.1_v4.2/)

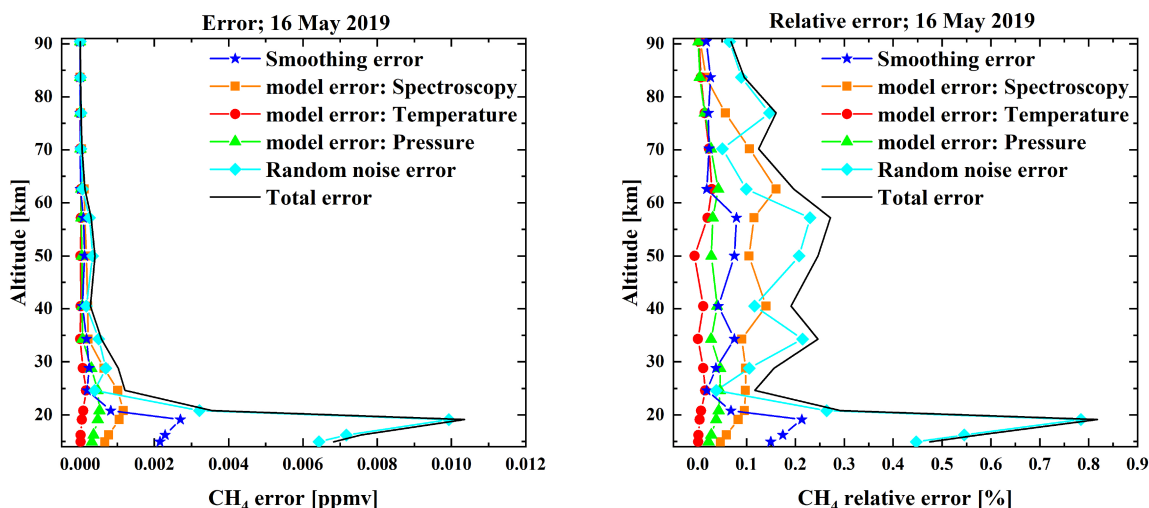


Fig. 10. Individual estimates of smoothing, forward model, and noise errors for  $\text{CH}_4$  retrieval. The absolute and relative errors correspond to the retrieval was carried for AIUS measurements recorded on May 16, 2019. Assumed uncertainties in the forward model errors are given in Table IV. The solid black line (Total error) represents the root sum squares of the smoothing error, forward model error, and random noise error.

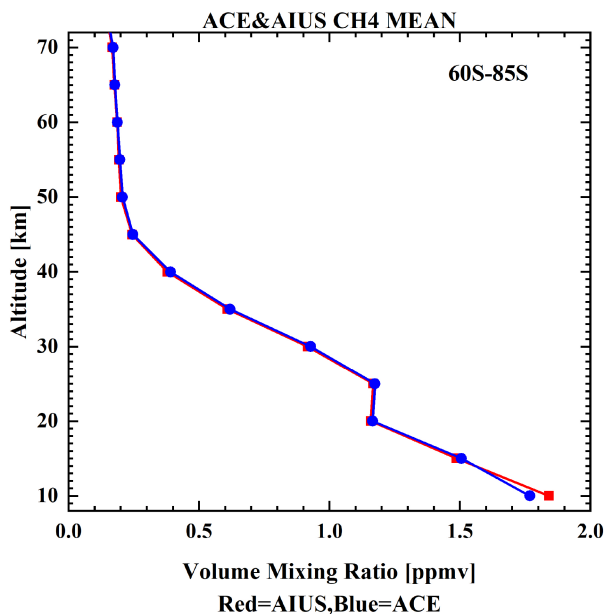


Fig. 11. Mean AIUS (red line) and ACE-FTS (blue line) profiles of the 42 matching pairs in the 60–85°S latitude zone are included in the comparison.

result, we obtained 42 orbits of AIUS data that met the matching criteria.

The range of altitude in which both instruments have good sensitivity is limited by the AIUS data at the lower boundary and by the ACE-FTS data at the upper boundary. We have set the lower altitude at 10 km and the upper altitude at 70 km. The inversion errors in this altitude range are very low, resulting in an uncertainty of 2–3% for ACE-FTS  $\text{CH}_4$ , while it may be more than 10% at the lowest and highest altitudes of the retrieval [10].

We needed to consider the difference in vertical resolution between the AIUS and the ACE-FTS when comparing their  $\text{CH}_4$  products. Indeed, the AIUS has a vertical resolution of approximately 4–11 km (with an altitude range of 10–60 km), whereas ACE-FTS has a resolution of around 4 km. In order to facilitate

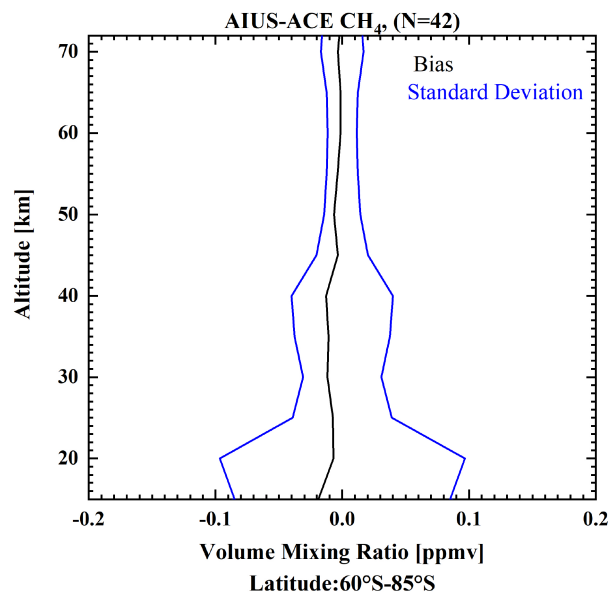


Fig. 12. Difference between AIUS and ACE-FTS profiles of  $\text{CH}_4$ .

comparison, the AIUS and ACE-FTS data are interpolated into a unified grid with intervals of 10, 15, 20, 25, 30, 35, 40, 45, 50, 55, 60, 65, and 70 km. The comparisons between AIUS and ACE-FTS  $\text{CH}_4$  profiles are shown in Figs. 11–13.

Fig. 11 shows the mean  $\text{CH}_4$  profiles of the AIUS and ACE-FTS between 10 and 70 km. It indicates that the AIUS and ACE-FTS  $\text{CH}_4$  mean profiles are in good agreement. The  $\text{CH}_4$  concentrations retrieved by the AIUS are slightly higher than those retrieved by ACE-FTS at 10 km.

The quantitative differences between AIUS and ACE-FTS are illustrated in Figs. 12 and 13. These differences are presented in absolute units and as percentages relative to ACE-FTS, respectively. The average relative standard deviation of the differences is smaller than 7.2% below 70 km.



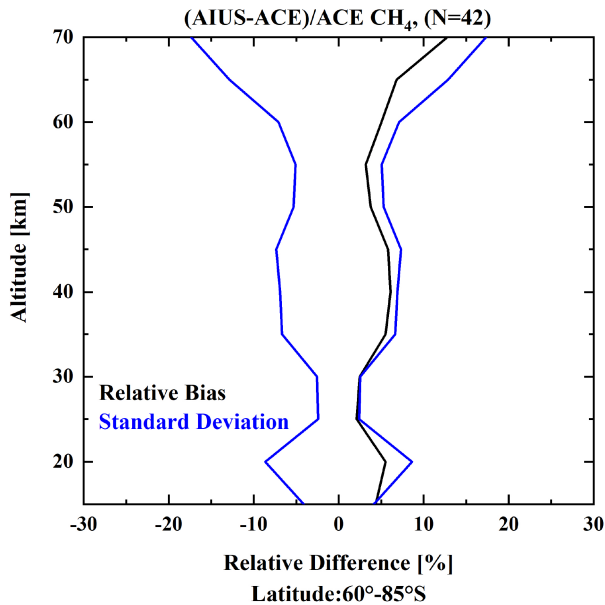


Fig. 13. Percentage relative difference between AIUS and ACE-FTS profiles of CH<sub>4</sub>.

#### IV. CONCLUSION

This article provides the CH<sub>4</sub> retrievals and an assessment of the atmospheric CH<sub>4</sub> profiles obtained during the second year of the AIUS mission, in order to quantify the level of agreement with other available instruments. The AIUS CH<sub>4</sub> products are retrieved from solar occultation measurements, using 120 spectral bands located between 1203 and 2906 cm<sup>-1</sup>.

The inversion results indicate that CH<sub>4</sub> mole fractions show significant variations in latitude and season over Antarctica. The CH<sub>4</sub> VMR is approximately 1.17–1.65 ppmv between 10 and 20 km, and decreases to less than 0.2 ppmv at 35 km in April and May.

The results of the error analysis indicate that the majority of the total error is caused by spectroscopic, random noise and smoothing errors. The uncertainties in temperature and pressure have a relatively minor impact on the total error. The comparison between AIUS CH<sub>4</sub> profiles and ACE-FTS CH<sub>4</sub> profiles reveals that the average relative standard deviation of the differences is smaller than 7.2% below 70 km.

In future work, the retrieval algorithm will be improved to retrieve more trace gases, such as HCN, CCl<sub>4</sub>, SF<sub>6</sub>, etc. Comparisons with ACE-FTS, ground stations, and other datasets are required for further validation.

#### ACKNOWLEDGMENT

The AIUS, ACE-FTS, and MLS data used in this study were acquired from the China Center for Resources Satellite Data and Application<sup>3</sup>, ACE/SCISAT Database<sup>4</sup>, and EARTH-DATA<sup>5</sup>, respectively. The authors would like to thank the Global Monitoring Laboratory for providing the trends in CH<sub>4</sub>, and

the HITRAN group members for providing the HITRAN online data.

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