# Improved Wavefront Reconstruction and Correction Strategy for Adaptive Optics System With a Plenoptic Sensor

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*Abstract***—Wavefront reconstruction and correction under strong turbulence has become a crucial issue for adaptive optics systems. An approach using a plenoptic sensor has attracted high attention recently. However, the existing reconstruction and correction algorithms based on the plenoptic sensor are sometimes unsatisfactory. Therefore, we propose a modified approach based on the checkerboard algorithm to reconstruct the wavefront. By means of a weighted summation, the reconstruction accuracy is clearly increased. We propose using the direct-gradient method to correct the disturbed wavefront after reconstruction, and our experimental results demonstrate the effectiveness with only one correction step.**

*Index Terms***—plenoptic sensor, wavefront reconstruction, wavefront correction.**

# I. INTRODUCTION

**A**DAPTIVE optics technology has been widely used to correct the wavefront distortion caused by atmospheric turbulence [1]. A wavefront sensor is often employed for wavefront detection, and a deformable mirror is employed for the correction. The Shack–Hartmann(S–H) wavefront sensor, which is the most widely used wavefront sensor, has been proven effective under weak turbulence conditions. However, the S-H wavefront sensor will encounter challenges in strong turbulence scenarios [2]. It will have considerate trouble reconstructing wavefronts with problems such as scintillation and branch points caused by strong turbulence.

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Various approaches have been proposed to overcome this problem. The use of a plenoptic camera for wavefront sensing is one option [3], [4]. With this solution, the wavefront can be estimated by measuring each individual micro-lens image. With the developments of the plenoptic sensing, a modified plenoptic sensor has been proposed by researchers at the University of Maryland College Park for wavefront detection [5]–[15]. It can simultaneously acquire the amplitude and phase information of the incident wavefront. Thereafter, the wavefront can be reconstructed under any turbulence conditions. The researchers have proposed two algorithms called checkerboard reconstruction algorithm and graph reconstruction algorithm to reconstruct and correct the wavefront [1]. However, the checkerboard algorithm does a good job, but it is time-consuming. The graph algorithm works quickly, but accuracy cannot be guaranteed. Therefore, there exists an untenable trade-off between speed and accuracy.

In this paper, we propose a modified checkerboard algorithm to determine the wavefront phase information. Subsequently, the direct-gradient method is used to correct the wavefront distortions. As a result, the wavefront can be reconstructed and corrected quickly with high accuracy.

The remainder of this article is organized as follows. The basic principles of the plenoptic sensor are introduced first in Section II. Then, we discuss the modified reconstruction algorithm, and present experimental results in Section III. In Section IV, the direct-gradient method is introduced to correct the wavefront. The study conclusions are summarized in Section V.

## II. PRINCIPLES OF THE PLENOPTIC SENSOR

## *A. Basic Structure and Design*

Several configurations of the plenoptic sensor have been developed for imaging purposes. Fig. 1 shows the basic structure of the plenoptic sensor we used. As can be seen, the plenoptic sensor contains an objective lens, a micro-lens array (MLA), and an image sensor. In Fig. 1, *fobj* and *dobj* represent the focal length and the diameter of the objective lens, respectively. The focal length and the pitch of the MLA are expressed as *fMLA* and *dMLA*, respectively. The back focal plane of the objective lens coincides with the front focal plane of the MLA. The image sensor is placed at the back focal plane of the MLA.

From Fourier optics, we can determine that the objective lens performs a 2-dimensional Fourier transform of the incident

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Fig. 1. Structure of the plenoptic sensor.

wavefront,  $t_1(x,y)$ , located at the front focal plane to the transformed wavefront,  $t_2(u,v)$ , located at the back focal plane [16]. Then the MLA performs another Fourier transform of  $t_2(u,v)$  to the back focal plane of the MLA,  $t_3(s,t)$ . Finally we can acquire the plenoptic images on the image sensor. The two steps are expressed by (1) and (2) as follows:

$$
t_2(u,v) = \frac{1}{j\lambda f_{obj}} \iint t_1(x,y) \exp\left\{\frac{-j2\pi}{\lambda f_{obj}}(ux+vy)\right\} dx dy
$$
\n(1)

$$
t_3(s,t) = \frac{1}{j\lambda f_{MLA}} \iint t_2(u,v) \sum_{M,N}
$$

$$
rect\left(\frac{u+s-2Md_{MLA}}{d_{MLA}}\right) rect\left(\frac{v+t-2Nd_{MLA}}{d_{MLA}}\right)
$$

$$
\exp\left\{\frac{-j2\pi}{\lambda f_{MLA}}(u - Md_{MLA})(s - Md_{MLA}) + (v - Nd_{MLA})(t - Nd_{MLA})\right\}dudv.
$$
 (2)

In (2), *(M, N)* represents the coordinates of each MLA lens. To avoid overlapping of the MLA sub-images, the *f*-number of the objective lens should be no less than that of the MLA lenslet [9], [12], which can be expressed as:

$$
\frac{d_{obj}}{f_{obj}} \le \frac{d_{MLA}}{f_{MLA}}.\tag{3}
$$

As Fig. 2 shows, if the incident beam is plane, the whole beam will only be imaged by the central MLA sub-aperture. The features of the incident beam will change if it is disturbed. It will be separated into different MLA sub-apertures according to the angular spectra of the wavefront [12]. Furthermore, the incident beam will be separated into more MLA sub-apertures when the turbulence gets stronger.

## *B. Performance Analysis*

To make full use of a wavefront sensor, performance characteristics such as sensitivity and dynamic range must be considered. As is discussed in Ref [12], the sensitivity and dynamic range of the plenoptic sensor can be calculated by (4) and (5):

$$
\left(\Delta x_{\min}, \Delta y_{\min}\right) = \frac{d_{MLA}}{f_{obj}}\tag{4}
$$

$$
(\Delta x_{\text{max}}, \Delta y_{\text{max}}) = \frac{Sensor\_Width}{2 \times d_{MLA}} \times \frac{d_{MLA}}{f_{obj}}
$$

$$
= \frac{Sensor\_Width}{2 \times f_{obj}}
$$
(5)

Equation (4) represents the calculation of the plenoptic sensor's sensitivity, which designates the minimum quantized wavefront tilt that can be measured by the plenoptic sensor. Eq. (5) represents the dynamic range of the plenoptic sensor, where *Sensor\_Width* represents the width of the image sensor. By multiplying the sensitivity by the half number of the MLA sub-apertures that the image sensor owns in one direction, the maximum measurable tilt can be calculated.

It can be easily found that a small MLA pitch can improve the sensitivity of the plenoptic sensor. But we cannot reduce the pitch up to a point, where diffraction of the small lens may cause the loss of local features, such that reconstruction cannot be performed. Furthermore, a long focal length of the objective length can also improve the sensitivity, but the dynamic range will be reduced. As a result, a trade-off between the sensitivity and dynamic range should be made to obtain a high wavefront sensing ability.

#### III. WAVEFRONT RECONSTRUCTION

# *A. Checkerboard Reconstruction Algorithm*

According to the previous theoretical analyses, each MLA sub-aperture represents a specific angular spectrum area of the wavefront. Based on this knowledge, researchers at the University of Maryland College Park proposed a checkerboard reconstruction algorithm [1] that divided each MLA cell into smaller "checkerboard" units. They further divided those units in each MLA sub-aperture to represent different spatial positions of the incident wavefront.

Taking Fig. 3 as an example, we can start with a simple  $3 \times 3$ MLA cell. Each MLA cell is then divided into four checkerboard units as "A", "B", "C", and "D" for simplicity. Each unit "A" in different sub-apertures represents the same spatial position of the incident wavefront regardless of the different angular spectra represented. Fig. 3 also shows the layout of a plane beam. The incident beam is only imaged by the center MLA cell as discussed before.

Fig. 4 describes the process of the checkerboard algorithm. When the incident beam is distorted, the entire beam will be separated into various MLA cells, just as shown in the left image of Fig. 4. In order to reconstruct the wavefront, we will examine each checkerboard unit to determine the local gradient at the corresponding location. Taking unit "A" as an example, we search for each MLA sub-aperture first and select the MLA sub-aperture having the largest pixel sum value in unit "A" area, which is cell  $(-1,0)$  in Fig. 4, to represent the spatial wavefront information at the corresponding position. The same operation is performed for the remaining units. A map of the local gradients



**Plane Wavefront** 

Fig. 2. Imaging diagram of the plenoptic sensor.



Fig. 3. Diagram of the checkerboard division and layout of plane beam.

for each checkerboard unit can then be acquired. By using the Zernike modal method, we can finally obtain the reconstructed wavefront.

#### *B. Modified Checkerboard Reconstruction Algorithm*

Sometimes the pixel-sum values for the same unit of the two or more MLA sub-apertures are extremely similar. In Fig. 5, unit "C" has a non-zero pixel value in three sub-apertures, which is cell $(-1,1)$ , cell $(1,0)$  and cell $(0,-1)$ . The pixel-sum values are 180, 55 and 186, respectively. In this case, it would be inaccurate to select the largest intensity to represent the wavefront information. Therefore, we propose a modified approach.

After searching for all MLA sub-apertures, we select the largest two cells for each unit and perform a weighted summation, which can be expressed as:

$$
M = \frac{V_1}{V_1 + V_2} M_1 + \frac{V_2}{V_1 + V_2} M_2
$$
 (6)

$$
N = \frac{V_1}{V_1 + V_2} N_1 + \frac{V_2}{V_1 + V_2} N_2.
$$
 (7)

In (6) and (7),  $(M_1, N_1)$  and  $(M_2, N_2)$  represent the coordinates of the MLA cells we selected, and  $V_1$  and  $V_2$  are their pixel sum values respectively. Instead of simply using the cell with



**Disturbed Wavefront** 



- 6: Calculate the wavefront gradients for each checkerboard units through (*M, N*).
- 7: Reconstruct the wavefront with the use of gradients information by the Zernike modal method

the largest intensity, we choose the weighted value to represent the wavefront gradient information.

If the values of  $V_1$  and  $V_2$  are differed by a large amount, the coordinates *(M, N)* calculated by (6) and (7) will be extremely close to the sub-aperture coordinates having the largest intensity. Therefore, the wavefront gradient information will also be very similar. If the values of  $V_1$  and  $V_2$  are close, a large number of errors will be produced using the original checkerboard algorithm, the calculated wavefront gradient information will be more accurate after the weighted summation. Based on the above analysis, our approach effectively improves the accuracy of the reconstruction without increasing too much running time, which may be less than 0.1 seconds in our experiment.

Subsequently, the same operation as we discussed before is performed again here to reconstruct the wavefront. The overall procedure of our approach, which can be called the "top-2" method, is shown in the following algorithm:

#### *C. Experimental Layout*

To demonstrate that our approach is effective, we built an experimental platform based on the plenoptic sensor. The schematic and actual experimental layout are shown in Figs 6





Fig. 4. Diagram of the checkerboard reconstruction.



Fig. 5. Layout of special condition.





Fig. 7. Actual experimental layout.

Fig. 6. Diagram of the experimental platform setup.

and 7, respectively. A 633nm HeNe laser passes through a beam splitter into Lens1, creating a parallel beam that spreads to our 21-element deformable mirror. By controlling the deformable mirror, the reflected beam with distortion returns and then passes through the Lens1 and Lens2 shrink-beam group to avoid overlapping. Finally, the beam is spread to the plenoptic sensor.

Because the focal length and the pitch of our MLA are 7mm and 0.2mm, respectively, and the focal length of our objective lens is 500mm, the width of the incident beam should be no greater than 14mm. Fig. 8 shows the layout diagram of our deformable mirror. As it shows, the distance between neighbouring actuator is 7mm, and the diameter of our deformable mirror is 35mm, hence we use a shrink-beam group comprised by Lens1 having a focal length of 500mm and Lens2 having a focal length of 200mm to satisfy the restriction. The camera we used is DMK



Fig. 8. Layout diagram of our 21-element deformable mirror actuators.

TABLE I RECONSTRUCTED RESULTS OF 0.5λ DEFORMATION

|                       | Aberration             | Defocus         | Astigmatism    | Coma           | Trefoil        |
|-----------------------|------------------------|-----------------|----------------|----------------|----------------|
| Our<br>approach       | Zernike<br>coefficient | $0.518\lambda$  | $0.469\lambda$ | $0.511\lambda$ | $0.520\lambda$ |
|                       | <b>RMS</b>             | $0.137\lambda$  | $0.157\lambda$ | $0.189\lambda$ | $0.162\lambda$ |
| Original<br>algorithm | Zernike<br>coefficient | 0.608 $\lambda$ | $0.456\lambda$ | $0.619\lambda$ | $0.551\lambda$ |
|                       | <b>RMS</b>             | በ 235እ          | $0.160\lambda$ | $0.265\lambda$ | $0.199\lambda$ |
|                       |                        |                 |                |                |                |

TABLE II RECONSTRUCTED RESULTS OF 1λ DEFORMATION

|                       | Aberration             | Defocus        | Astigmatism    | Coma            | Trefoil            |
|-----------------------|------------------------|----------------|----------------|-----------------|--------------------|
| Our<br>approach       | Zernike<br>coefficient | $1.174\lambda$ | $0.964\lambda$ | 0.996እ          | 0.988 <sub>A</sub> |
|                       | <b>RMS</b>             | $0.274\lambda$ | $0.163\lambda$ | $0.138\lambda$  | $0.149\lambda$     |
| Original<br>algorithm | Zernike<br>coefficient | $0.958\lambda$ | $1.036\lambda$ | 1.065 $\lambda$ | $0.911\lambda$     |
|                       | <b>RMS</b>             | $0.245\lambda$ | $0.164\lambda$ | $0.253\lambda$  | $0.208\lambda$     |

TABLE III RECONSTRUCTED RESULTS OF 1.5λ DEFORMATION



23UM021 from the Imaging Source, which contains a  $480\times480$ resolution with  $7.5\mu$ m pixel width in the binning mode.

#### *D. Results and Analysis*

Various Zernike polynomials can be applied to the deformable mirror to create different Zernike deformations [17]. The polynomials we used are described as follows:

$$
Defocus: Z_2^0(r, \theta) = \sqrt{3}(2r^2 - 1)
$$
 (8)

$$
Astigmatism: Z_2^2(r,\theta) = \sqrt{6}r^2\sin(2\theta)
$$
 (9)

$$
Coma: Z_3^1(r, \theta) = \sqrt{8}(3r^3 - 2r)\sin\theta \qquad (10)
$$

$$
Trefoil : Z_3^3(r, \theta) = \sqrt{8}r^3 \sin(3\theta). \tag{11}
$$

In (8)–(11),  $(r, \theta)$  represent the normalized radius and angle, respectively. Fig. 9 shows the images of the different Zernike

TABLE IV RECONSTRUCTED RESULTS OF 2λ DEFORMATION

|                       | Aberration                    | Defocus            | Astigmatism        | Coma              | Trefoil            |
|-----------------------|-------------------------------|--------------------|--------------------|-------------------|--------------------|
| Our<br>approach       | <b>Zernike</b><br>coefficient | 2292 <sub>h</sub>  | 1.999 <sup>A</sup> | 1877 <sub>A</sub> | $1.913\lambda$     |
|                       | <b>RMS</b>                    | 0.307 <sub>A</sub> | $0.207\lambda$     | $0.248\lambda$    | $0.185\lambda$     |
| Original<br>algorithm | Zernike<br>coefficient        | $2.345\lambda$     | $2.445\lambda$     | $1.701\lambda$    | 1.773 <sub>A</sub> |
|                       | <b>RMS</b>                    | $0.408\lambda$     | $0.454\lambda$     | $0.405\lambda$    | $0.383\lambda$     |

TABLE V RESULTS OF 0.5λ DEFORMATION AFTER CORRECTION

|            | Aberration | Defocus        | Astigmatism    | Coma           | Trefoil        |
|------------|------------|----------------|----------------|----------------|----------------|
| <b>RMS</b> | Our        | $0.078\lambda$ | $0.097\lambda$ | $0.091\lambda$ | $0.089\lambda$ |
|            | approach   |                |                |                |                |
|            | Original   | $0.116\lambda$ | $0.107\lambda$ | $0.209\lambda$ | $0.118\lambda$ |
|            | algorithm  |                |                |                |                |

TABLE VI RESULTS OF 1λ DEFORMATION AFTER CORRECTION



deformations, but only a smaller  $250\times250$  pixel section is displayed.

To demonstrate the effectiveness of our approach, both algorithms were used to reconstruct the wavefront of the images obtained previously. In our experiment, the total plenoptic image is  $480\times480$  pixel, and the MLA we used is  $15\times15$ . We divide each MLA sub-aperture into 64 checkerboard units, and the size of each checkerboard unit is  $3 \times 3$  pixel.

Fig. 10 shows the reconstructed results of Astigmatism for two algorithms. A two-wavelength deformation was applied to the deformable mirror. Fig. 9(c) shows the image obtained from the plenoptic sensor. Fig.  $10(a)$  is the result of the reconstructed wavefront using the original checkerboard reconstruction algorithm. Fig. 10(c) separates the reconstructed result into different Zernike modes. Figs 10(b) and (d) represent the same results after improvement.

It is obvious from Figs 10(a) and (b) that the reconstructed wavefront is the result of Astigmatism. In Figs 10(c) and (d), we can see that Zernike mode 5, which represents the Zernike deformation of Astigmatism, is the dominant mode. This also demonstrates the effectiveness of both reconstruction algorithms. Moreover, the value of the Zernike coefficient is 2.445λ and the root-mean-square (RMS) phase reconstruction error is 0.454λ using original checkerboard algorithm. The RMS phase reconstruction error is calculated by:

$$
RMS = \sqrt{\sum_{i=1}^{10} (Z_{rec}(i) - Z_{input}(i))^2}
$$
 (12)

where *Zinput* represents the input coefficients of the various Zernike modes, *Zrec* represents the reconstructed coefficients of the various Zernike modes, and *i* represents the index of the



Fig. 9. Plenoptic images of different Zernike deformations.



Fig. 10. Reconstruction of Astigmatism.

TABLE VII RESULTS OF 1.5λ DEFORMATION AFTER CORRECTION

|            | Aberration            | Defocus        | Astigmatism    | Coma            | Trefoil        |
|------------|-----------------------|----------------|----------------|-----------------|----------------|
| <b>RMS</b> | Our<br>approach       | $0.101\lambda$ | $0.153\lambda$ | $0.154\lambda$  | $0.102\lambda$ |
|            | Original<br>algorithm | $0.264\lambda$ | $0.295\lambda$ | 0.338 $\lambda$ | $0.329\lambda$ |
|            |                       |                |                |                 |                |

TABLE VIII RESULTS OF 2λ DEFORMATION AFTER CORRECTION



Zernike coefficients. The results respectively drop to 1.999λ and  $0.207\lambda$  after improvement, which demonstrates that a higher reconstruction accuracy is obtained with our approach.

More reconstructed results can be found in Tables I–IV. It can be concluded that our approach results in a lower error for most of the cases and comparable results for the corner cases.

## IV. WAVEFRONT CORRECTION

#### *A. Direct-Gradient Method*

Because the reconstructed wavefront gradient reflects the wavefront distortion information, through correction of the wavefront gradient error, the overall distortion can be effectively corrected. Therefore, after the wavefront gradient is determined, the direct-gradient method can be used to correct the distortion [18]–[21]. Generally, there is a linear relationship between the local gradient and the drive voltages of the deformable mirror actuators, which can be expressed as:

$$
G = D \cdot V \tag{13}
$$

where  $G = [g_{1x}, g_{1y}, \dots g_{mx}, g_{my}]^T$  is the vector of the wavefront gradients, and *m* is the number of the checkerboard units, which is 64 in our experiment.  $V = [v_1, v_2, \dots v_n]^T$  is the vector of the drive voltages, and  $n$  is the number of the actuators employed for our deformable mirror, which is 21 in our experiment. *D* is a  $2m \times n$  constant matrix called the influence matrix, which bridges the deformable mirror to our plenoptic sensor.

By recording the gradients response and the driving voltages for the various determined deformations, the influence matrix *D* can be easily acquired. After the influence matrix *D* is determined, we can easily obtain the correction voltage for any incoming wave by:

$$
V = D^+ \cdot G \tag{14}
$$

where  $D^+$  is the pseudo inverse of matrix *D*.

## *B. Results and Analysis*

To demonstrate that the above discussions are effective, we allow the direct-gradient method to correct the disturbed plenoptic images we get before. Some experimental results using our modified algorithm are shown in Fig. 11. Fig. 11(b) shows the



Fig. 11. Correction of different Zernike deformations.

result of Astigmatism after one correction step. We can see that the entire beam almost locates in single MLA cell and the RMS error drops to 0.100λ, which means a nearly plane wavefront. Compared with the work in Ref [1], it costs 5 correction steps to remove the major distortion. More correction results can be found in Tables V–VIII, which demonstrates the quickness and the effectiveness of our method.

To further prove the effectiveness of our modified algorithm, original checkerboard algorithm is also used for correction. Tables V–VIII also present the detailed results after one correction step using the original algorithm. It can be seen that with a higher reconstruction accuracy obtained by our approach, a better correction result can also be acquired.

## V. CONCLUSION

In this paper, we propose a modified checkerboard reconstruction algorithm based on the plenoptic sensor. A weighted summation is added to acquire the wavefront gradient information. Subsequently, we build an experimental platform to validate the effectiveness of our approach. The experimental results demonstrate that a higher reconstruction accuracy can be obtained after improvement. Moreover, we propose to use the direct-gradient method to correct the wavefront distortion. After determination of the influence matrix and the wavefront gradients by calculation, distortions can then be corrected quickly and effectively. The experimental results indicate that most distortions can be removed after only one step of correction.

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