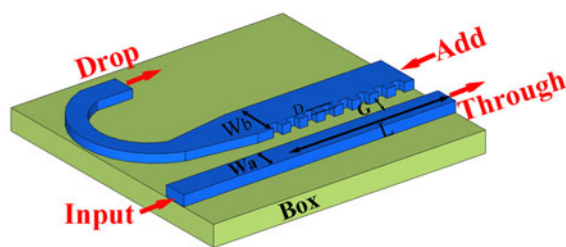


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**Abstract:** We demonstrate the design, fabrication, and characterization of a two-mode silicon add-drop filter, which is based on multimode grating-assisted contra-directional couplers. The mode conversions between two asymmetry strip waveguides are realized by using spatially periodic corrugations in one side wall of a multimode waveguide at the phase-match condition. Experimental results show that the fundamental mode of a single-mode waveguide is successfully coupled to the first two modes of the multimode waveguide in the opposite direction. The corresponding mode conversion bandwidths and extinction ratios are also measured. The demonstrated add-drop filter may find their applications in a hybrid multiplexing system by combing wavelength-division multiplexing technology and mode-division multiplexing technology for on-chip optical interconnect.

**Index Terms:** Integrated optics device, gratings, wavelength filtering devices, Bragg reflectors, waveguides.

## 1. Introduction

Silicon photonics devices are desirable for the realization of on-chip optical interconnections with high capacity for their compatibility with a standard complementary metal-oxide-semiconductor (CMOS) technology [1]. As an important component of wavelength division multiplexing (WDM) systems, various add-drop filters have been developed for silicon-on-insulator (SOI) platform, such as Bragg gratings, ring-resonators and contra-directional couplers (contra-DCs). Bragg gratings have perfect spectral characteristics for (de)multiplexing in a WDM system [2]–[5]. However, they are two-port devices and require a circulator in the system. Ring-resonator add-drop filters have received much attention for their compact size and various applications [6], [7], but they are limited by relatively small free spectral ranges (FSR). Contra-DCs are grating assisted add-drop filters and have shown their good properties as an add-drop filter in a WDM system [8]–[12]. However, conventional Contra-DCs usually consist of two asymmetry single-mode waveguides, which support only two fundamental compound modes and thus one mode channel.

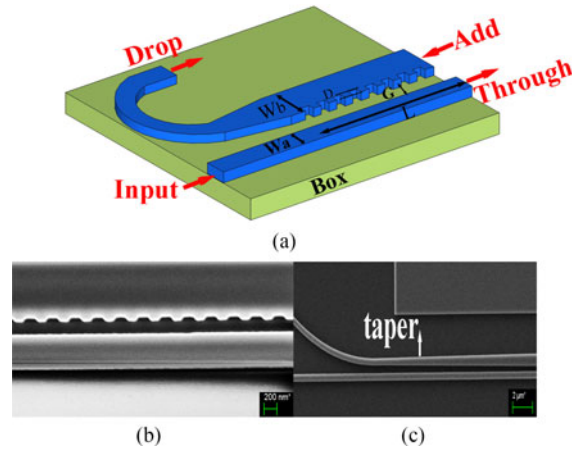


Fig. 1. (a) Three-dimensional view of the proposed MGACCs. The SEM image of the (b) coupling part (c) tapered waveguide.

Mode-division multiplexing (MDM) technology is a good way to expand the capacity of a single photonics link without increasing too much footprint and complexity of the system [13], [14]. Recently, Bragg gratings based on multimode silicon waveguides have attracted a lot of interests due to their unique features [15]–[17]. The fundamental mode of a single mode waveguide is coupled into an adjacent multi-mode waveguide grating at the wavelength of interest by using Contra-DCs [16]. An integrated space-division multiplexing (SDM) is proposed on a chip with small device footprint and high packing density by cascading multiple contra-DCs [17]. In this paper, we design and fabricate a two-mode silicon add-drop filter based on multimode grating-assisted contra-directional couplers (MGACCs) on the SOI platform. In our device, the fundamental mode of a single-mode strip waveguide is coupled to the  $TE_0$  and  $TE_1$  modes of a multimode strip waveguide with different designed Bragg grating periods, respectively. A tapered waveguide is designed to distinguish and convert the  $TE_0$  and  $TE_1$  modes. In order to drop the  $TE_1$  mode in the multimode waveguide, another MGACCs is cascaded. The experimental results show that the modes conversions between the two waveguides are realized successfully. The proposed device can combine its WDM functionalities with MDM capabilities to further augment the multiplexing capacity of the device.

## 2. Device Design and Analysis

The proposed MGACCs is schematically illustrated in Fig. 1(a). The device consists of two asymmetry parallel strip waveguides on the top of  $SiO_2$  box layer. The cladding layer is air. Waveguide *a* supports only  $TE_0$  mode, while waveguide *b* supports two TE modes. The corrugations in waveguide *b* are placed in one sidewall, which faces the straight waveguide *a*. The corrugated waveguide *b* is used as the add-drop port, and the straight waveguide *a* is used as the input-through port. The widths of the two strip waveguides, which are separated by a gap with a width of  $G$ , are  $W_a$  and  $W_b$ , respectively. An adiabatic taper connects the MGACCs with a single-mode waveguide at the drop-port.

When light is input from the straight waveguide *a* (see Fig. 1), the forward cross-coupling between the two waveguides is negligible due to the phase mismatch caused by the width difference. Meanwhile, the contra-directional coupling can be realized by the spatially periodic corrugation at the phase-match condition [15]

$$\lambda/\Lambda = (n_a + n_{bi}) \quad (i = 0, 1) \quad (1)$$

where  $n_a$  and  $n_b$  represent the TE mode effective indices of waveguides *a* and *b*, respectively,  $i$  is the mode order,  $\Lambda$  is the grating period, and  $\lambda$  is the Bragg wavelength.

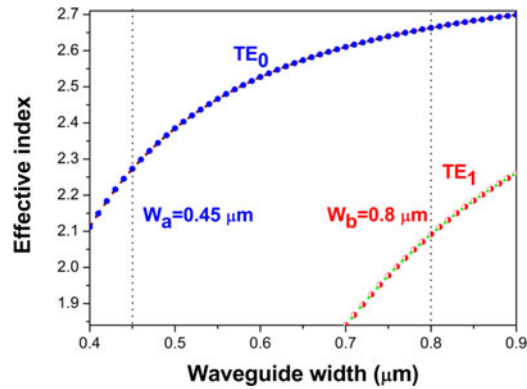


Fig. 2. Relationships between effective index and waveguide width.

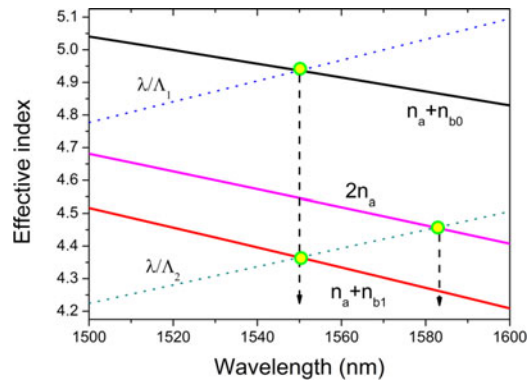


Fig. 3. Calculated effective indices of the modes with the phase-match condition. ( $n_a$ ,  $n_{b0}$ , and  $n_{b1}$  are the effective indices of the fundamental mode of the input waveguide and the first two modes of the multimode waveguide.)

By designing the grating period, one can couple the  $TE_0$  mode of the input waveguide to the  $TE_0$  and  $TE_1$  modes of the multimode waveguide, respectively. When the dropped  $TE_0$  mode in the multimode waveguide propagates through the tapered waveguide, it will then be adiabatically transitioned into a forward propagating fundamental mode of the drop waveguide. On the contrary, the dropped  $TE_1$  mode will be radiated once it enters into the drop waveguide, as the drop waveguide supports only the fundamental mode. Thus we can ensure that only the  $TE_0$  mode could be output at the drop-port. To drop the  $TE_1$  mode, we could cascade another MGACCs before the taper.

In our device, the thickness of the silicon strip waveguides is chosen to be 220 nm. Fig. 2 shows the effective index as a function of waveguide width for the first two mode orders at the wavelength of 1550 nm. The straight waveguide has a width of 0.45  $\mu\text{m}$ , while the width of the multimode waveguide is chosen to be 0.8  $\mu\text{m}$  to support two TE modes. The coupling gap  $G$  is set to be 200 nm, the corrugation width  $D$  is 50 nm, the duty cycle of the Bragg gratings is 50%, and the coupling length  $L$  is 500  $\mu\text{m}$ .

The solid and dotted lines in Fig. 3 are the plots of the right-hand side and left-hand side of the (1), as a function of the wavelength, respectively. Therefore, the cross points of the solid and dotted lines represent the phase-match condition. We choose the values of  $\Lambda$  for effective contra-directional coupling at the wavelength of 1550 nm. It is observed that in the device, a larger value of  $\Lambda$  is needed for higher order mode as the sum of  $n_a$  and  $n_b$  decreases. The calculation results show that the grating periods for coupling from the fundamental mode of the input waveguide to the  $TE_0$  and  $TE_1$  modes of the multimode waveguide are  $\Lambda_1 = 314$  nm and  $\Lambda_2 = 356$  nm, respectively. The phase-match condition for the inter-coupling in the input waveguide, i.e.,  $\lambda/\Lambda = 2n_a$ , is also plotted

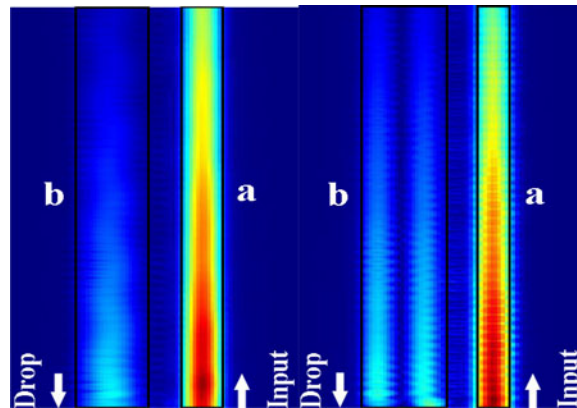


Fig. 4. Three-dimensional FDTD simulations of the MGACCs for light coupling from the straight input waveguide *a* at the wavelength of 1550 nm to (a) the TE<sub>0</sub> mode and (b) the TE<sub>1</sub> mode in the multimode waveguide *b*.

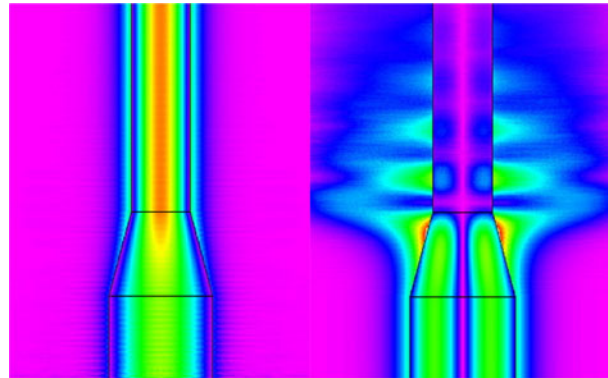


Fig. 5. Light propagation along the designed taper for (a) the TE<sub>0</sub> mode and (b) the TE<sub>1</sub> mode in the multimode waveguide.

in Fig. 3. It can be observed that the Bragg wavelength is far beyond 1550 nm, when we select the grating periods of 314 nm and 356 nm, respectively. Therefore, we can design two grating couplers to couple the fundamental modes of two input waveguides to the TE<sub>0</sub> and TE<sub>1</sub> modes of the bus multimode waveguide at the same wavelength, respectively. In our work, we fabricated two devices for the two modes conversions, respectively. It is worth pointing out that the inter-coupling in the input waveguide is greatly suppressed, as the corrugations are designed far away from the input waveguide [8].

In order to better understand this concept, we perform 3-D finite-difference time-domain (FDTD) simulations. In view of the computation time and memory, we set a short coupling length of only 10  $\mu\text{m}$ . The mode conversions between the input waveguide and the multimode waveguide are calculated by a 3-D-FDTD simulation tool, as presented in Fig. 4. Fig. 4(a) and 4(b) show that the fundamental mode of the input waveguide is contra-directionally coupled to the TE<sub>0</sub> and TE<sub>1</sub> mode in the multimode waveguide at the wavelength of 1550 nm, respectively.

A 20- $\mu\text{m}$ -length tapered waveguide with the core width varying from 0.8  $\mu\text{m}$  to 0.45  $\mu\text{m}$  is designed to modify the mode size at the multimode waveguide. The simulation of the light propagating from the multimode waveguide to a single-mode drop waveguide via the taper is illustrated in Fig. 5 by using beam propagation method (BPM). As seen from Fig. 5(a), when the TE<sub>0</sub> mode passes through the tapered waveguide, it adiabatically transits into the fundamental mode of the drop waveguide. On the other hand, when the TE<sub>1</sub> mode light reaches the tapered waveguide, it totally transits into the radiation mode and dissipates, as shown in Fig. 5(b). It is due to that the TE<sub>1</sub> mode

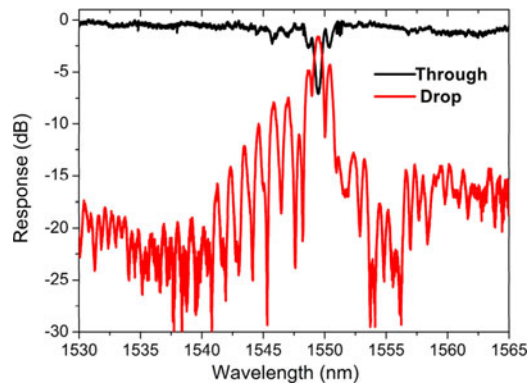


Fig. 6. Normalized through-port and drop-port spectra for contra-directional coupling from the fundamental of the input waveguide to the  $TE_0$  mode of the multimode waveguide.

and the fundamental mode of the drop waveguide are mismatched. Therefore, we can determine the order of the dropped mode in the multimode waveguide by measuring the power at the through port and drop port.

### 3. Experimental Results and Discussion

The proposed GMACCs are fabricated on an SOI wafer with a 220-nm-thick top silicon layer and a 3- $\mu\text{m}$ -thick buried oxide layer. In the experiment, E-beam lithography (Vistec EBP 5200) was used to define the structures on the ZEP520A resist. Then the patterns were transferred to the top silicon layer by inductively coupled plasma (ICP) dry etching using SF<sub>6</sub> and C<sub>4</sub>F<sub>8</sub> gases. To share the same etching depth as the strip waveguide, we employ photonics crystal based grating couplers. The scanning electron microscopy (SEM) image of the fabricated GMACCs and tapered waveguide are shown in Fig. 1(b) and (c). In the measurements, the TE-polarized lights from a tunable laser (SANTEC, TSL-510) were coupled into/out of the chip by grating couplers. The output spectra were measured by an optical power detector and power meter (HP, 8153a).

To verify the operation principle, we at first characterize the mode conversion between the fundamental mode of the input waveguide and the  $TE_0$  mode of the multimode waveguide. Fig. 6 shows the measured normalized through port and drop port spectra with the designed grating period  $\Lambda = 314$  nm. The coupling loss of the two grating couplers has been normalized. As seen from Fig. 6, when light is launched at the input port, the contra-directional coupling occurs at the central wavelength of 1550 nm. As the device has a slight corrugation depth of 50 nm and large coupling gap of 200 nm, the coupling strength is weakened. Therefore, the add drop filter shows good wavelength selectivity with a measured bandwidth of  $\sim 0.8$  nm. One could obtain a broad bandwidth by enhancing the coupling coefficient, such as larger corrugation depth or narrower coupling gap [11]. The extinction ratio is -7.2 dB at the resonant wavelength, which can be increased by extending the coupling length since the extinction ratio is a function of both the coupling coefficient and the coupling length [18]. The insertion loss of -1.2 dB at the central wavelength for the drop port spectrum is also measured. Owing to the uniform corrugations, the device has undesirable sideband lobes in the spectra, which can be compressed in future by apodization [19]. There is no other notch in the spectrum throughout the scanned wavelength range, which implies no inter-coupling in the input waveguide occurs.

Then, the mode coupling between the fundamental mode of the input waveguide and the  $TE_1$  mode of the multimode waveguide is demonstrated with the designed grating period of 356 nm. As shown in Fig. 7, in the through-port spectrum, there is a resonant notch with a bandwidth of  $\sim 2.3$  nm and an extinction ratio of -18 dB at the central wavelength of 1560 nm. However, the relevant peak is not measured at the drop port spectrum. In order to investigate the coupling mode at the resonant wavelength, an optical fiber circulator is used to measure the reflection power at

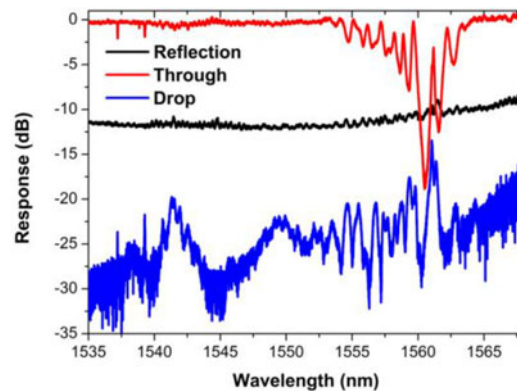


Fig. 7. Normalized through port and drop port spectra for contra-directional coupling from the fundamental mode of the input waveguide to the  $TE_1$  mode of the multimode waveguide.

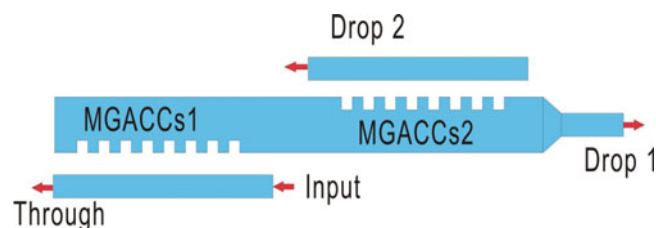


Fig. 8. Pair of GMACCs for the  $TE_1$  mode coupling in the multimode waveguide.

the input port. It can be seen that there is no resonance in the reflection spectrum (see the black line in Fig. 7). Therefore, we believe that the resonance in the through-port spectrum (is not caused by the intra-waveguide Bragg reflection. From the drop port spectrum, one can see the power around the resonant wavelength is lower than -20 dB. According to the analysis, we confirm that the fundamental mode of the input waveguide is coupled to the  $TE_1$  mode of the multimode waveguide and the  $TE_1$  mode power finally dissipates when passing through the taper. In addition, one can see that the bandwidth and extinction ratio of the  $TE_1$  mode coupling in the through port spectrum are larger than those in Fig. 6, which means the coupling coefficient between the fundamental mode and the  $TE_1$  mode of the multimode waveguide is larger. The results agree well with our conclusions in [15].

Higher order modes of multimode waveguides are usually converted to the fundamental mode of a single-mode waveguide before they are measured. There are a lot of ways to achieve this conversion, such as directional couplers (DCs) [20], [21]. Contra-DCs [15], asymmetric Y-junctions [22]–[24], or adiabatic couplers [25]. In our device, a pair of MGACCs is design to couple the fundamental mode of the input waveguide to the  $TE_1$  mode of the multimode waveguide, and then the  $TE_1$  mode is converted to the fundamental mode of a single-mode drop waveguide, as shown in Fig. 8. Both of the two MGACCs have the same parameters. In order to enhance the coupling strength, the MGACCs have a large corrugation depth of 100 nm and a narrow coupling gap of 150 nm, respectively.

Fig. 9 illustrates the measured transmission spectra at the through and drop ports after normalizing the insertion loss of the grating coupler. From the spectrum of the through port (black line in Fig. 9), one can see that the mode coupling between the fundamental mode of the input waveguide and the  $TE_1$  mode of the multimode by MGACCs1, is greatly strengthened, as a broad bandwidth of about 13 nm and an high extinction ratio of over -21 dB are measured. The spectrum at the second drop port demonstrates that the  $TE_1$  mode in the multimode waveguide is coupled to the fundamental mode in the second drop waveguide. The spectrum does not match well with that at the

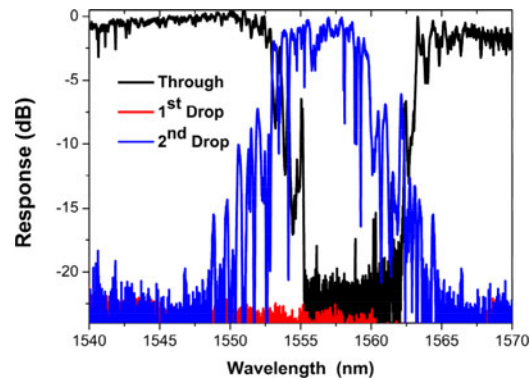


Fig. 9. Normalized through port and drop port spectra for the mode conversions from the fundamental mode of the input waveguide to the  $TE_1$  mode of the multimode waveguide and finally to second drop waveguide.

through-port with a  $\sim 3$  nm blue shift. Therefore, the bandwidth of the spectrum decreases to about 10 nm. We attribute this to that the resonant Bragg wavelength of MGACCs2 shifts to the blue wavelength. The difference of the Bragg wavelengths is because that the two fabricated MGACCs are not totally same. There are some undesired abnormal oscillations in the spectrum, which is due to the phase noise caused by the fabrication errors and waveguide sidewall roughness. We could improve our device by more careful fabrication. The spectrum at the first drop-port is also measured, and the result shows that the crosstalk between the two drop ports is larger than 22 dB at the wavelength of 1559 nm. It is worth to point out that more mode channels can be achieved by cascading more pairs of GMACCs.

#### 4. Conclusion

In summary, we have proposed a two-mode add-drop filter based on GMACCs. The design, analysis and fabrication of the device are demonstrated. By using the GMACCs, the fundamental mode of the input waveguide could be coupled to the  $TE_0$  and  $TE_1$  modes of the multimode waveguide, respectively. The  $TE_1$  mode conversion is observed to have a stronger coupling coefficient with larger bandwidths and higher extinction ratios than that of  $TE_0$  mode. The measured results show that the designed filter have the potentiality to be used in WDM and MDM systems for on-chip optical interconnect.

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