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#### Abstract

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#### Abstract

We report a low-loss $\operatorname{InP}$ wire waveguide monolithically integrated with an $\mathrm{SiO}_{x}$ waveguide on an Si platform. By means of directly bonding InP to $\mathrm{SiO}_{2}$, we realized a submicrometer-scale InP wire waveguide with high optical confinement. In addition, a several-micrometer-scale $\mathrm{SiO}_{x}$ waveguide with a refractive index difference of $\sim 3 \%$ is integrated on the InP wire waveguide via a spot-size converter (SSC). Experimental results indicate that the InP wire waveguide has a propagation loss of $5 \mathrm{~dB} / \mathrm{cm}$ and that the $\mathrm{InP}-\mathrm{SiO}_{x} \mathrm{SSC}$ has an insertion loss of 0.7 dB and a reflectance of less than -50 dB . By using the SSC, a low coupling loss of 0.9 dB is achieved between the optical fiber and the $\operatorname{lnP}$ wire waveguide.


Index Terms: $\operatorname{InP}$ wire waveguide, Si photonics, spot-size converter, $\mathrm{SiO}_{x}$ waveguide.

## 1. Introduction

Over the past few decades, Si photonics technology has evolved as a promising photonic integration platform. With the goal of complete integration of passive and active photonic devices, Si and non-Si materials are being integrated on the Si photonics platform to compensate for performance disadvantage due to the physical properties of Si. In particular, for active devices, such as laser light sources and amplifiers, monolithic integration of group III-V materials on the Si platform is promising and is being intensively investigated worldwide [1]-[3]. Recently, by utilizing the high-index contrast between group III-V materials and $\mathrm{SiO}_{2}$, III-V membrane photonic devices with a very thin III-V layer on $\mathrm{SiO}_{2} / \mathrm{Si}$ substrate are realized [4]-[7]. For integration, the membrane devices are connected with output InP waveguides and then connected to Si photonic devices. In addition, III-V wire waveguides can be applied to nonlinear photonics thanks to high optical confinement similar to that in Si wire waveguides [8], [9].

Up to now, III-V waveguides and the integration of III-V passive and active devices have been demonstrated on an Si platform. However, there are few material options; it is either all-III-V or Si-III-V integration [1], [7]. Within their material systems, the performance of passive devices based on high-index-contrast waveguides can easily degrade because the fabrication error of core size leads to considerable variation of the effective index. In addition, for such high-indexcontrast waveguides, efficient fiber coupling is an important issue. Integrations of other materials


Fig. 1. Schematic of the $\operatorname{InP}-\mathrm{SiO}_{x}$ waveguide integration structure. The $\mathrm{SiO}_{2}$ overcladding of the $\mathrm{SiO}_{x}$ waveguide is not drawn.
can provide promising solutions to these problems. In particular, a key technology is the integration of a low-refractive-index-difference (low- $\Delta$ ) [10] waveguides made of $\mathrm{SiO}_{x} \mathrm{~N}_{y}$-based materials, which are suitable for high-performance passive devices [11].

In this paper, we report the monolithic integration of InP wire waveguides and low- $\Delta \mathrm{SiO}_{x}$ waveguides. We form $\operatorname{InP}$ waveguides with sub- $\mu \mathrm{m}$ core dimensions on an $\mathrm{SiO}_{2} / \mathrm{Si}$ substrate using the direct bonding technique. Then, using low-temperature fabrication methods, we monolithically fabricate $\mathrm{SiO}_{x}$ waveguides with $3-\% \Delta$. To connect these two different waveguides with low loss, we use a spot-size converter (SSC) with an InP inverse-taper structure. With this integration technique, we achieve low-loss fiber coupling between a high-numerical-aperture (highNA ) single-mode fiber and the InP wire waveguide.

## 2. Design

Fig. 1 shows a schematic of the structure of the $\operatorname{InP}-\mathrm{SiO}_{x}$ integration using the SSC. The $\operatorname{InP}$ core ( $n=3.169$ ) is configured on an $\mathrm{SiO}_{2}$ undercladding ( $n=1.444$ ). An $\mathrm{SiO}_{x}$ layer ( $n=1.505$ ) covers the InP core as an overcladding. In this configuration, the InP core is completely embedded in $\mathrm{SiO}_{x}$. The $\mathrm{SiO}_{x}$ layer is partly removed to pattern the $\mathrm{SiO}_{x}$ waveguide core. On top, a $\mathrm{SiO}_{2}$ layer ( $n=1.468$ ) covers the overall structure as an overcladding of the $\mathrm{SiO}_{x}$ waveguide.

The InP waveguide has $\Delta$ of $39 \%$, which is quite high compared with other conventional dielectric waveguides, such as silica glass, and similar to that of an Si wire waveguide (41\%). Hence, the dimensions satisfying the single-mode condition are on the order of sub $\mu \mathrm{m}$. Fig. 2 shows calculated effective refractive indices ( $n_{\text {eff }}$ ) at 1550 -nm wavelength for transverse-electric (TE) -like modes as a function of $\operatorname{InP}$ core width ( $w_{\mathrm{InP}}$ ) with various core heights ( $h_{\mathrm{InP}}$ ). The mode notations are taken from ref. [12]. For $h_{\text {lnP }}$ of 200, 300, and 400 nm , the single-mode condition can be satisfied when $w_{\text {inp's }}$ s are less than 580,460 , and 380 nm , respectively. The dimensions satisfying the single-mode condition are slightly larger than those of Si [12].

The $\mathrm{SiO}_{x}$ core is $2 \times 4 \mu \mathrm{~m}$, which is one order of magnitude larger than the InP core. To connect them with low loss, in this work, we used an SSC consisting of an InP adiabatic taper and a $\mathrm{SiO}_{x}$-waveguide core covering the taper, as shown in Fig. 1. This SSC structure is similar to the SSC for $\mathrm{Si}_{\mathrm{SiO}}^{x}$ connection, which functions as a mode matching section between the $\operatorname{InP}$ wire and the following $\mathrm{SiO}_{x}$ waveguides [11], [13]. By tuning the taper length ( $L$ ) and the tip width ( $w_{\text {tp }}$ ), the propagating mode inside the $\operatorname{InP}$ wire could adiabatically evolve into the mode inside the $\mathrm{SiO}_{x}$ waveguide. Fig. 3(a)-(c) show calculated SSC losses as functions of $w_{\mathrm{tp}}$ and $L$ with $h_{\text {InP }}$ of 200, 300, and 400 nm , respectively. Here, the SSC loss is defined as the optical power ratio between the $\mathrm{SiO}_{x}$ waveguide output to the InP waveguide input for the fundamental TE mode at $1550-\mathrm{nm}$ wavelength. For the calculation, we used a commercial software based on a two-dimensional eigenmode solver and eigenmode-expansion calculation [14]. For each $h_{\text {inp }}$, the SSC efficiency tends to become higher as $L$ increases and $w_{\text {tp }}$ decreases. When $w_{\text {tp }}$ is too large,


Fig. 2. Calculated effective indices ( $n_{\text {eff }}$ ) of the $\operatorname{InP}$ waveguide modes as a function of $\operatorname{InP}$ core width ( $w_{\mathrm{InP}}$ ) with various $\operatorname{InP}$ core heights ( $h_{\operatorname{lnP}}$ ).


Fig. 3. Calculated SSC conversion loss as functions of taper length $(L)$ and $\operatorname{InP}$ tip width ( $w_{\text {tp }}$ ) with various InP core heights ( $h_{\mathrm{InP}}$ ). (a) $h_{\mathrm{lnP}}=200 \mathrm{~nm}$, (b) $h_{\mathrm{lnP}}=300 \mathrm{~nm}$, and (c) $h_{\mathrm{InP}}=400 \mathrm{~nm}$, respectively.
the SSC efficiency can not increase even though $L$ increases, and it reaches a low-efficiency plateau, the value of which depends on $w_{\text {tp }}$. This is because, with $w_{\text {tp }}$ that is too large, the efficiency is dominated by mode mismatch at the boundary between the InP -taper tip and the $\mathrm{SiO}_{x}$ waveguide. The coupling loss between the $\mathrm{SiO}_{x}$ waveguide and the high-NA fiber was also estimated to be 0.2 dB by the calculation.

Another concern is reflection. When the SSC is integrated with a laser diode, reflected light increases oscillating-mode instability in it. Thus, the reflection at the SSC should be carefully suppressed. Fig. 4(a)-(c) show calculated reflectance at the SSC when the light was input from the InP waveguide as functions of tip width $\left(w_{\text {tp }}\right)$ and taper length $(L)$ with $h_{\text {InP }}$ of 200, 300, and 400 nm , respectively. In order to achieve less than $-50-\mathrm{dB}$ reflectance, $w_{\text {tp }}$ should be set to $<150,<120$, and $<100 \mathrm{~nm}$ for $h_{\mathrm{lnP}}=200,300$, and 400 nm , respectively. As $w_{\mathrm{tp}}$ increases, the reflection at the InP tip increases. In addition, the $L$ should be set to more than $100 \mu \mathrm{~m}$ regardless of the $h_{\text {InP }}$ value. This is because, with $L$ of less than $100 \mu \mathrm{~m}$, the adiabatic mode conversion can not be achieved at the taper section, which also increases reflectance.

On the basis of these results, to confirm the single-mode propagation in the InP waveguide, we designed the InP core with dimensions of $h_{\mathrm{InP}}=200 \mathrm{~nm}$ and $w_{\mathrm{InP}}=500 \mathrm{~nm}$. For the SSC, to confirm the low-loss connection between the InP and $\mathrm{SiO}_{x}$ waveguides, we used $w_{\mathrm{tp}}=80 \mathrm{~nm}$ and $L=300 \mu \mathrm{~m}$.

## 3. Fabrication

In advance, we prepared two substrates: a III-V bonding substrate layered as $\operatorname{InP} / \mathrm{InGaAs} / \mathrm{InP}$ and a base substrate layered as $\mathrm{SiO}_{2} / \mathrm{Si}$. To prepare the III-V bonding substrate, we grew an etchstop $\operatorname{InGaAs}$ layer and an $\operatorname{InP}$ layer on $\operatorname{InP}$ wafer using the metal-organic vapor phase


Fig. 4. Calculated SSC reflectance as functions of taper length $(L)$ and $\operatorname{InP}$ tip width ( $w_{\text {tp }}$ ) with various $\operatorname{lnP}$ core heights $\left(h_{\mathrm{InP}}\right)$. (a) $h_{\mathrm{InP}}=200 \mathrm{~nm}$, (b) $h_{\mathrm{lnP}}=300 \mathrm{~nm}$, and (c) $h_{\mathrm{InP}}=400 \mathrm{~nm}$, respectively.


Fig. 5. SEM images of (a) InP taper tip. (b) InP core. (c) Optical microscope image of a fabricated chip.
epitaxy (MOVPE) method. The thickness of the top InP layer was set to 200 nm . The base substrate was a Si wafer with a top thermal-oxide layer of $2 \mu \mathrm{~m}$. After surface cleaning and pre treatment, these two wafers were directly bonded by using an $\mathrm{O}_{2}$ plasma assisted bonding technique. Then, after InP support substrate and InGaAs etch-stop layer had been removed, we obtained the InP on $\mathrm{SiO}_{2}$ substrate. On this substrate, we deposited $\mathrm{SiO}_{2}$ as a hard mask for InP etching and formed a resist pattern using electron-beam lithography to define the $\operatorname{InP}$ core and taper. Then, after the $\mathrm{SiO}_{2}$ hard-mask pattern had been formed by reactive-ion etching (RIE), the InP core was formed by inductive-coupled-plasma (ICP) RIE. Fig. 5(a) and (b) show SEM images of a fabricated $\operatorname{InP}$ taper tip and $\operatorname{InP}$ core. The $\operatorname{InP}$ layer was formed with $w_{\text {tp }}$ and $w_{\text {inP }}$ of 137 and 551 nm , respectively. The InP width is slightly larger than designed due to the multiple etching process. Nevertheless, the single-mode condition is satisfied and SSC conversion loss is expected to be less than 0.2 dB .

Next, we fabricated $\mathrm{SiO}_{x}$ waveguides. First, a $2-\mu \mathrm{m}$-thick $\mathrm{SiO}_{x}$ layer was deposited by electroncyclotron resonance (ECR) plasma-enhanced chemical vapor deposition (PECVD) [11], this was followed by definition of the low-refractive-index core of SSC to have a width of $4 \mu \mathrm{~m}$ by photolithography and RIE. In the definition of $\mathrm{SiO}_{x}$ core, the offset between InP and $\mathrm{SiO}_{x}$ cores can occur due to misalignment of our stepper lithography, which leads to polarization rotation in the SSC. Fig. 6(b) shows a calculated transmittance for the TE-like mode input to the TE-like mode output and the TE-like mode input to the TM-like mode output, as a function of misalignment defined in Fig. 6(a). Here, $w_{\text {tp }}$ and $L$ are set to 140 nm and $300 \mu \mathrm{~m}$, respectively. We numerically confirmed that the polarization rotation is less than -20 dB even if the misalignment of 100 nm exists, which is a typical value on our stepper lithography. Thus, the polarization rotation is not a critical issue. Then, the $\mathrm{SiO}_{2}$ overcladding was deposited by ECR PECVD. The processed wafer was cut by dicing to obtain facets for the fiber coupling. The facet lines cross only the $\mathrm{SiO}_{x}$ waveguide and no polishing was performed. Fig. 5(c) shows an optical microscope image of a fabricated chip. The meandering structures are InP wire waveguides with different length.


Fig. 6. (a) Definition of the misalignment. (b) Calculated transmittance for the TE-like mode input to the TE-like mode output and the TE-like mode input to the TM-like mode output as a function of misalignment.


Fig. 7. Propagation characteristics of InP waveguides. (a) Waveguide transmittance as a function of $\operatorname{InP}$ length at a wavelength of 1550 nm . (Inset) Transmission spectra with various waveguide lengths. (b) Propagation loss spectrum estimated by the slope of fitted line in (a).

## 4. Experiments and Discussion

First, we characterized the propagation loss of the InP waveguide. In experiments, we used cleaved high-NA fibers that were butt-coupled with the input and output facets of the $\mathrm{SiO}_{x}$ waveguide. The input polarization was set to the TE mode. We obtained transmission spectra of the fabricated waveguides with InP -waveguide lengths of $0.83,1.56$, and 3.16 cm , respectively. Fig. 7(a) shows waveguide transmittance as a function of InP-waveguide length at 1550 nm , and the inset shows the transmission spectra of waveguides with various InP lengths. From the slope of the fitted line in Fig. 7 (a), we estimated the propagation loss of $5.3 \mathrm{~dB} / \mathrm{cm}$ at a wavelength of 1550 nm . In a similar way at each wavelength, we also acquired the propagation-loss spectra and obtained about $5 \mathrm{~dB} / \mathrm{cm}$ within the C-band, as shown in Fig. 7(b). Note that the propagation loss of the monolithically fabricated $\mathrm{SiO}_{x}$ waveguide is $0.8 \mathrm{~dB} / \mathrm{cm}$, which is almost same performance as an unintegrated $\mathrm{SiO}_{x}$ waveguide [11].

Next, for the loss estimation of the SSC connecting the InP and $\mathrm{SiO}_{x}$ waveguides, as shown in Fig. 8(a), we measured waveguides with various numbers of $\mathrm{InP}-\mathrm{SiO}_{x}$ connections to the SSCs. We estimated the SSC loss from the slope of the relationship between the waveguide transmittance and number of SSCs. At 1550 nm , we estimated the SSC loss to be $0.7 \mathrm{~dB} / \mathrm{SSC}$. Fig. 8(b) shows measured SSC loss as a function of wavelength. We confirmed that the SSC loss ranges from 0.6 to $1.0 \mathrm{~dB} / \mathrm{SSC}$ in the C-band. The loss is slightly higher than our calculation shown in Fig. 3. This is due to scattering loss that occurred because of the surface roughness of the InP taper, which we ignored in the calculation. In addition, the fabricated taper tip was wider than the design value, which increased the SSC loss. Then, we characterized the spatial distribution of reflection in the $\operatorname{InP}-\mathrm{SiO}_{x}$ connected waveguide using an optical low-coherence reflectometer (OLCR). For the measurement, we used a cleaved dispersion-compensated fiber (DCF) at the input facet. with index matching oil to suppress the reflection at the


Fig. 8. Transmission characteristics of the SSC. (a) Relationship between number of SSCs and transmittance. (Inset) Test-pattern layout. (b) Estimated SSC loss spectrum.


Fig. 9. OLCR trace of an $\operatorname{InP-} \mathrm{SiO}_{x}$ connected waveguide.

DCF- $\mathrm{SiO}_{x}$ facet as much as possible. Fig. 9 shows a measured OLCR trace obtained with spatial resolution of $20 \mu \mathrm{~m}$. The vertical axis was normalized by reflectance at the DCF-air boundary. No significant reflection is observed at the SSC, and the intensity of the reflection appears to be almost the same as that of back scattering from the InP waveguide. Considering DCF-SiO ${ }_{x}$ coupling loss and round-trip propagation loss to the reflection point, we estimated the reflectance is less than -50 dB at the SSC. As a result, we clearly confirmed that the SSC can connect the $\operatorname{InP}$ waveguide and the $\mathrm{SiO}_{x}$ waveguide with low loss and low reflectance. This technique can be further adopted for the monolithic integration of InP membrane active devices on $\mathrm{Si}-, \mathrm{Ge}-, \mathrm{SiO}_{x^{-}}$, and electronics-integrated photonic devices [15], [16].

In addition, we examined the coupling loss between the InP waveguide and a high-NA fiber. Fig. 10 shows measured coupling loss. By means of the standard method, we obtained 1.6 dB at 1550 nm by simply dividing in half the intercept on the graph of the relationship between waveguide transmittance and length as shown in Fig. 7(a). However, this 1.6 dB can be further broken down to four parts as follows: facet loss between the high-NA fiber and the $\mathrm{SiO}_{x}$ waveguide, $\mathrm{SiO}_{x}$ waveguide loss, SSC loss, and InP bending loss. Among them, facet loss is estimated to be 0.2 dB by the calculation. The $\mathrm{SiO}_{x}$ waveguide loss was estimated to be less than 0.1 dB and negligible because the length of $\mathrm{SiO}_{x}$ at the input and output is less than 1 mm and the $\mathrm{SiO}_{x}$ propagation loss is $0.8 \mathrm{~dB} / \mathrm{cm}$. The SSC loss is $0.7 \mathrm{~dB} / \mathrm{SSC}$ as shown in Fig. 8. These values are estimated at 1550 nm , and their variations within the C-band are negligible on the physical basis. We therefore consider that the residual 0.7 dB originates from 12 bends with a radius of $10 \mu \mathrm{~m}$, which leads to $0.058 \mathrm{~dB} / 90 \mathrm{deg}$-bend. Subtracting the bending loss, we estimated the fiber-to-InP-waveguide coupling loss to be only 0.9 dB at wavelength of 1550 nm as shown in Fig. 10.

Finally, we compare the obtained characteristics with existing reports using Table I. Concerning to the propagation loss, our waveguide exhibits preferably low loss among sub- $\mu \mathrm{m}$-wide


Fig. 10. Measured coupling loss from a high-NA fiber to the $\operatorname{InP}$ waveguide as a function of wavelength.

TABLE I
Comparison of III-V Wire Waveguides on Si Platform

| Bonding method | Core | Core dimensions <br> $w \times h[\mathrm{~nm}]$ | Prop. loss <br> $[\mathrm{dB} / \mathrm{cm}]$ | Fiber coup. loss <br> [dB/coup.] | Ref. |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Direct $\left(\mathrm{InP} / \mathrm{SiO}_{2}\right)$ | $\operatorname{lnP}$ | $560 \times 200$ | 5.2 | 0.9 (high-NA fiber) | This work |
| Direct $\left(\mathrm{SiO}_{2} / \mathrm{SiO}_{2}\right)$ | GaAs | $400 \times 400$ | 45 | 13.5 (lense fiber) | $[17]$ |
| Direct $\left(\operatorname{lnGaAsP/\mathrm {SiO}_{2})}\right.$ | $\operatorname{InGaAsP}$ | $2000 \times 315$ | 30 | 4.2 | $[19][20]$ |
| Direct $\left(\mathrm{Al}_{2} \mathrm{O}_{3} / \mathrm{Al}_{2} \mathrm{O}_{3}\right)$ | $\operatorname{InP} / \operatorname{InGaAsP}$ | $2000 \times 375$ | 4 | - | $[7]$ |
| Adhesive $\left(\mathrm{BCB} / \mathrm{SiO}_{2}\right)$ | $\operatorname{InP}$ | $400 \times 250$ | 3.3 | 6.4 | $[6][18]$ |
| Adhesive $\left(\mathrm{BCB} / \mathrm{SiO}_{2}\right)$ | GalnAsP | $410 \times 150$ | 4 | - | $[21]$ |
| Adhesive $\left(\mathrm{BCB} / \mathrm{SiO}_{2}\right)$ | $\operatorname{InGaP}$ | $630 \times 250$ | 12 | 7.5 | $[8]$ |

waveguides fabricated by the direct bonding technique. Among the waveguides fabricated by the direct bonding technology, a lower propagation loss is reported [7]. However, this structure has a much larger core width of 2000 nm , whereas we use a sub- $\mu \mathrm{m}$ core width, which makes a direct comparison of the loss values difficult. We believe the propagation loss of our waveguide can be further decreases by improving the bonding process, bonding structure, and the techniques for III-V nanofabrication on the Si platform. For adhesively bonded waveguides, lower loss than this work, comparable to that of Si wire waveguides in early days has been reported [6], [18], [21]. However, the fact that our concept relies on direct bonding instead of adhesive bonding makes it highly promising for cost-effective mass fabrication of III-V devices monolithically integrated on $\geq 300-\mathrm{mm}$ Si wafers and subsequent processing using standard microelectronic CMOS manufacturing facilities [4]. For that purpose, the high-temperature processes needed for III-V crystal growth and activation annealing are performed in the front-end-of-line (FEOL). If adhesive bonding technology is used for III-V integration it would be difficult to carry out such high-temperature processes. Therefore, we consider the direct bonding technology is promising.

As for the fiber coupling loss, to the best of our knowledge, the lowest coupling loss from a standard single-mode fiber (SMF) to the InP waveguide on the Si platform is $4.2 \mathrm{~dB} /$ coupling with a grating coupler [20]. Even though we used a high-NA fiber in this work, we expect the SMF-to-InP-waveguide coupling of around 1 dB can be achieved thanks to the low-loss $\operatorname{InP}-\mathrm{SiO}_{x} \mathrm{SSC}$. Because a high-NA fiber can be connected with SMF with only $0.2-\mathrm{dB}$ loss by the thermally-expanded-core process and the fiber-fusion splicing. We have previously reported an SSC for the SMF to Si -wire waveguide coupling [23]. Moreover, the SSC provides a wide operation bandwidth of over 40 nm as shown in Fig. 8 and this bandwidth potentially exceeds over 100 nm [22].

We believe this low-loss and broadband fiber coupling technique is very useful for the membrane lasers on Si and their wavelength-division-multiplexing application and for the nonlinear application of the III-V wire waveguides.

## 5. Conclusion

By means of the direct bonding technique, an $\operatorname{InP}$ wire waveguide with propagation loss of $5 \mathrm{~dB} / \mathrm{cm}$ in the C-band was developed on an Si platform. In addition, we monolithically integrated the $\operatorname{InP}$ wire waveguide and an $\mathrm{SiO}_{x}$ waveguide and achieved a $0.7-\mathrm{dB}$ connection between InP wire and $\mathrm{SiO}_{x}$ waveguides via a SSC with low reflectance of less than -50 dB . Moreover, using the SSC, we achieved low-loss fiber coupling. This work is a first step for the integration of III-V active devices with the $\mathrm{Si}-, \mathrm{Ge}-, \mathrm{SiO}_{x} \mathrm{~N}_{y}-$, and electronics-integrated photonic circuits.

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## References

[1] G. Roelkens et al., "III-V/silicon photonics for on-chip and intra-chip optical interconnects," Laser Photonics Rev., vol. 4, no. 6, pp. 751-779, Nov. 2010.
[2] A. Descos et al., "Heterogeneously integrated III-V/Si distributed Bragg reflector laser with adiabatic coupling," in Proc. ECOC, 2013, pp. 1-3, Th.1.B.2.
[3] C. Zhang et al., "Low threshold and high speed short cavity distributed feedback hybrid silicon lasers," Opt. Exp., vol. 22, no. 9, pp. 10202-10209.
[4] S. Matsuo et al., "Directly modulated buried heterostructure DFB laser on $\mathrm{SiO}_{2} / \mathrm{Si}$ substrate fabricated by regrowth of InP using bonded active layer," Opt. Exp., vol. 22, no. 10, pp. 12139-12147, 2014.
[5] S. Matsuo et al., "Directly modulated DFB laser on $\mathrm{SiO}_{2} / \mathrm{Si}$ substrate for datacenter networks", J. Lightwave Technol., vol. 33, no. 6, pp. 1217-1222, Mar. 2015.
[6] J. van der Tol et al., "Photonic integration in Indium-Phosphide Membranes on Silicon (IMOS)," in Proc. SPIE, vol. 8988, Mar. 2014, Art ID. 89880M.
[7] Y. Ikku et al., "Low-driving-current InGaAsP photonic-wire optical switches using III-V CMOS photonics platform," Opt. Exp., vol. 20, no. 26, pp. B357-B364, Dec. 2012.
[8] U. Dave et al., "Nonlinear properties of dispersion engineered InGaP photonic wire waveguides in the telecommunication wavelength range," Opt. Exp., vol. 23, no. 4, pp. 4650-4657, Feb. 2015.
[9] M. Pu et al., "AlGaAs-on-insulator nanowire with 750 nm FWM bandwidth, -9 dB CW conversion efficiency, and ultrafast operation enabling record Tbaud wavelength conversion," in Proc. OFC, 2015, vol. 1, pp. 1-3, Th5A.3.
[10] K. Okamoto, Fundamentals of Optical Waveguides, 2nd ed. London, U.K.: Academic, 2006.
[11] T. Tsuchizawa et al., "Monolithic integration of silicon-, germanium-, and silica-based optical devices for telecommunications applications," IEEE J. Sel. Top. Quantum Electron., vol. 17, no. 3, pp. 516-525, May 2011.
[12] K. Yamada, "Silicon photonic wire waveguides: Fundamentals and applications," in Silicon Photonics II, D. Lockwood and L. Pavesi, Eds. Berlin, Germany: Springer-Verlag, Ch. 1.
[13] T. Shoji, T. Tsuchizawa, T. Watanabe, K. Yamada, and H. Morita, "Low loss mode size converter from 0.3 um square Si wire waveguides to singlemode fibres," Electron. Lett., vol. 38, no. 25, pp. 1669-1670, Dec. 2002.
[14] FIMMWAVE/FIMMPROP, Photon Design Inc., Oxford, U.K.
[15] H. Nishi et al., "Monolithic integration of a silica-based arrayed waveguide grating filter and silicon variable optical attenuators based on p-i-n carrier-injection structures," Appl. Phys. Exp., vol. 3, no. 10, Oct. 2010, Art ID. 102203.
[16] H. Nishi et al., "Monolithic integration of a silica AWG and Ge photodiodes on Si photonic platform for one-chip WDM receiver," Opt. Exp., vol. 20, no. 8, pp. 9312-9321, Apr. 2012.
[17] R. Kafouf et al., "Ultrahigh relative refractive index contrast GaAs nanowire waveguides," Appl. Phys. Exp., vol. 1, no. 12, 2008, Art ID. 122101.
[18] Y. Jiao et al., "Fullerene-assisted electron-beam lithography for pattern improvement and loss reduction in InP membrane waveguide devices," Opt. Lett., vol. 39, no. 6, pp. 1645-1648, Mar. 2014.
[19] M. Takenaka, M. Yokoyama, M. Sugiyama, Y. Nakano, and S. Takagi, "InGaAsP photonic wire based ultrasmall arrayed waveguide grating multiplexer on Si wafer," Appl. Phys. Exp., vol. 2, no. 12, 2009, Art ID. 122201.
[20] M. Takenaka, M. Yokoyama, M. Sugiyama, Y. Nakano, and S. Takagi, "InGaAsP grating couplers fabricated using complementary-metal-oxide-semiconductor-compatible III-V-on-insulator on Si," Appl. Phys. Exp., vol. 6, no. 4, 2013, Art ID. 042501.
[21] J. Lee et al., "Low-loss GalnAsP wire waveguide on Si substrate with benzocyclobutene adhesive wafer bonding for membrane photonic circuits," Jpn. J. Appl. Phys., vol. 51, no. 4R, Mar. 2012, Art ID. 042201.
[22] T. Tsuchizawa et al., "Microphotonics devices based on silicon microfabrication technology," IEEE J. Sel. Top. Quantum Electron., vol. 11, no. 1, pp. 232-240, Jan. 2005.
[23] H. Nishi, T. Tsuchizawa, H. Fukuda, H. Okayama, and K. Yamada, "Connection of Si-wire waveguides with singlemode fibers by inverse-taper spot-size converters," in Proc. ISPEC, 2014, P-50.

