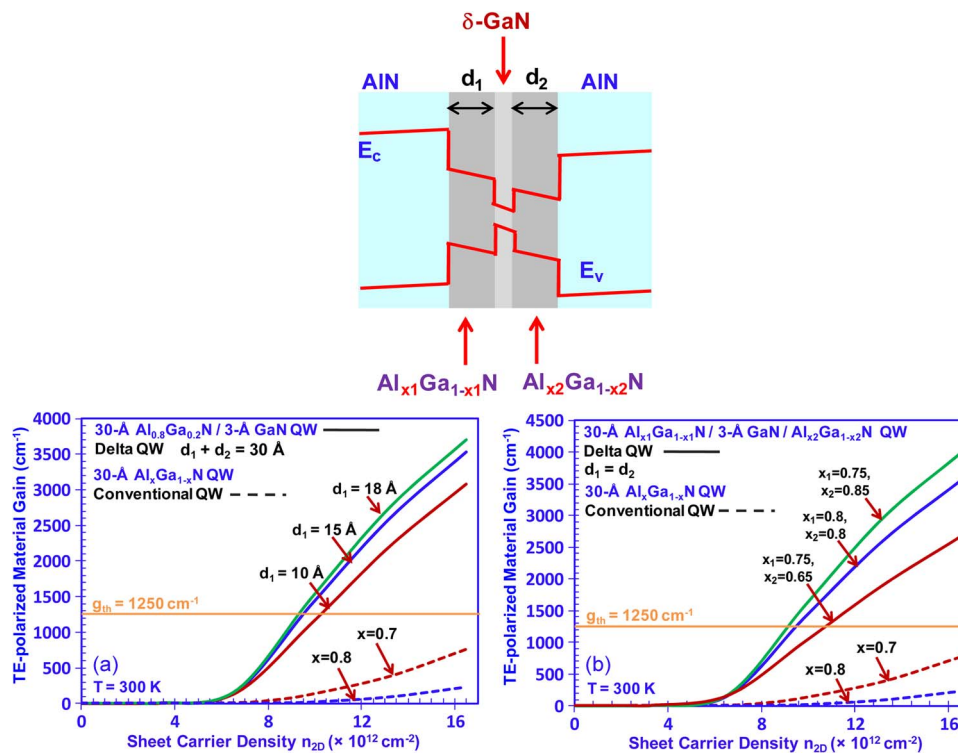


Engineering of AlGa_xN-Delta-GaN Quantum-Well Gain Media for Mid- and Deep-Ultraviolet Lasers

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Engineering of AlGa_N-Delta-GaN Quantum-Well Gain Media for Mid- and Deep-Ultraviolet Lasers

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Abstract: The gain characteristics of AlGa_N-delta-GaN quantum wells (QWs) with varying delta-GaN positions and AlGa_N QW compositions are analyzed. The use of optimized AlGa_N-delta-GaN QWs resulted in ~7-times increase in material gain over that of conventional AlGa_N QWs for gain media emitting at ~240 nm. By employing asymmetric AlGa_N-delta-GaN QWs, the optimized optical gain can be achievable for AlGa_N-delta-GaN QW structure with realistic design applicable for mid- and deep-ultraviolet (UV) lasers. The threshold properties and differential gains are also studied for optimized AlGa_N-delta-GaN QWs UV lasers.

Index Terms: III-Nitride, AlGa_N-delta-GaN quantum wells (QWs), deep UV (UV) lasers, optical gain, laser diodes.

1. Introduction

The field of III-Nitride semiconductors has applications covering lasers and light-emitting diodes (LEDs) [1]–[16], power electronics [17], thermoelectric [18], [19], photovoltaics [20], and terahertz photonics [21]. Significant advances had been achieved in visible LEDs, especially in addressing the internal quantum efficiency [1]–[4], material epitaxy [5]–[10], extraction efficiency [11], [12], and efficiency droop [13]–[16]. In contrast to the progress in visible emitters [1]–[16], the realization of electrically injected mid- ($\lambda \sim 250$ – 320 nm) and deep-ultraviolet (UV) ($\lambda \sim 220$ – 250 nm) AlGa_N quantum wells (QWs) lasers has been limited to $\lambda \sim 320$ – 360 nm [22]–[39], while only optically pumped deep-UV lasers had been realized for shorter emission wavelength [31]. The challenges in realizing the electrically injected mid- and deep-UV AlGa_N QWs lasers are attributed to the difficult growths of high Al-content AlGa_N and p-AlN and the lack understanding on the physics of high Al-content AlGa_N QWs. The optical properties for low [40]–[42] and high Al-content [43], [44] AlGa_N QWs lasers have been reported, while the studies are relatively lacking on gain properties of high Al-content AlGa_N QWs.

Our recent work [43] revealed the valence subbands crossover in high Al-content AlGa_N QW leads to strong conduction (C)-crystal-field split-off hole (CH) transition, which results in large transverse-magnetic (TM)-polarized optical gain for $\lambda \sim 220$ – 230 nm. Recent experimental works have confirmed the dominant TM-polarized emission for shorter wavelength with higher Al-content

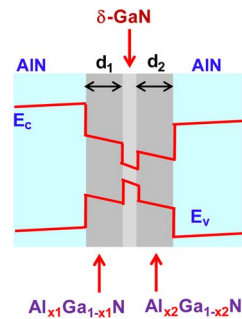


Fig. 1. Schematic energy band lineup of $\text{Al}_{x_1}\text{Ga}_{1-x_1}\text{N}/\delta\text{-GaN}/\text{Al}_{x_2}\text{Ga}_{1-x_2}\text{N}$ QW with AlN barriers.

AlGa_xN QWs LEDs [45], [46], in agreement with our prediction [43]. The subsequent paper has studied about the temperature and barrier effect on the polarization properties of AlGa_xN QW [47]. Recent theoretical work [48] has confirmed the crossover of the valence subbands for AlGa_xN alloys on $\text{Al}_{0.85}\text{Ga}_{0.15}\text{N}$ substrate, in agreement with our finding [43]. The subsequent theoretical work [49] has discussed the strain effect by using various AlGa_xN substrates on the crossover in AlGa_xN. Garrett and coworkers have experimentally confirmed our prediction of the CH/HH crossover for AlGa_xN QWs from the transverse-electric (TE) ($\lambda \sim 253$ nm) to TM ($\lambda \sim 237$ nm) photoluminescence polarization switching [50].

In order to address the issue of low optical gain at $\lambda \sim 250\text{--}300$ nm, we recently proposed the AlGa_xN-delta-GaN QW structure [51] by inserting an ultrathin GaN layer into high Al-content (x) $\text{Al}_x\text{Ga}_{1-x}\text{N}$ -delta-GaN QW active region. The use of AlGa_xN-delta-GaN QWs resulted in strong valence subband mixing, which led to large TE-polarized gain at $\lambda \sim 240\text{--}300$ nm attributed to the dominant C–HH transition. Recent experimental work [52] has reported the polarization properties of the deep-UV emission from AlN/GaN short-period superlattices, which revealed that stronger TE-polarized emission from the use of a very thin GaN in AlN active region similar to our prediction [51]. Note that the use of InGa_xN-delta-InN QW had been used for achieving improved overlap design in visible LEDs [53], [54].

In this paper, we present a comprehensive study on the optical gain and threshold characteristics of AlGa_xN-delta-GaN QWs with various delta-GaN positions and AlGa_xN QW compositions for mid- and deep-UV lasers. The band structure and wave function calculations are based on the six-band $\mathbf{k} \cdot \mathbf{p}$ formalism for wurtzite semiconductors taking into account the valence band mixing, strain effect, polarization fields, and carrier screening effects [55]–[59], with the III-Nitride band parameters and band offsets ratio ($\Delta E_c/\Delta E_v = 70/30$) obtained from references [59]–[61]. The spontaneous and piezoelectric polarizations follow the treatments in [62] and [63], respectively.

2. Concept and Band Structure for AlGa_xN-Delta-GaN QWs

Fig. 1 shows the schematic of the AlGa_xN-delta-GaN QW. Our previous work [51] revealed that the use of symmetric high Al-content AlGa_xN-delta-GaN QWs led to large TE-polarized gain at $\lambda \sim 240\text{--}300$ nm. In this paper, we present a comprehensive study of the asymmetric AlGa_xN-delta-GaN QWs. The asymmetric QWs can be introduced by engineering the thicknesses of the AlGa_xN layers (d_1, d_2) and/or the Al contents of the AlGa_xN layers (x_1, x_2) surrounding the delta-GaN layer.

For the AlGa_xN-delta-GaN QWs, the energy levels of the HH and light hole (LH) subbands are higher than those of the CH subband, which leads to the dominant C–HH transition. The difference in the effective masses of the electrons and holes leads to different spreading of the wave functions for electrons and holes in III-Nitride-based QWs. In the case of identical effective masses for both electrons and holes, the optimized overlap will occur with the delta position at the center of the QWs. However, the effective masses for holes are larger than those of electrons in nitride-based QWs; thus, the optimization of the overlap can be achieved by using engineering of the delta layer

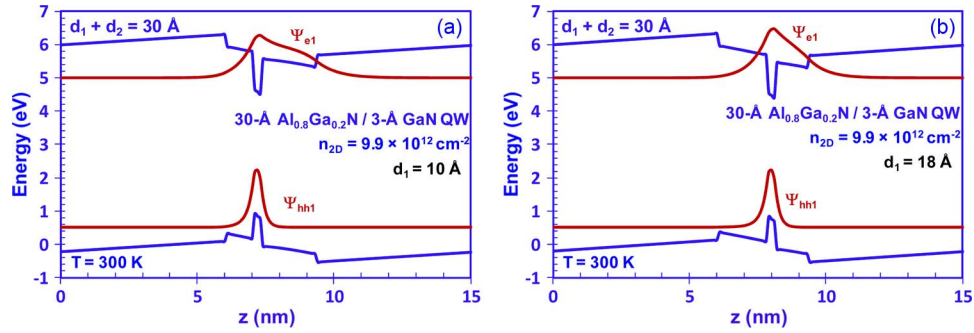


Fig. 2. Energy lineups with electron wavefunction EC1 and hole wave function HH1 for the 30-Å $\text{Al}_{0.8}\text{Ga}_{0.2}\text{N}/3\text{-}\text{\AA}$ GaN QW with (a) $d_1 = 10$ Å, and (b) $d_1 = 18$ Å.

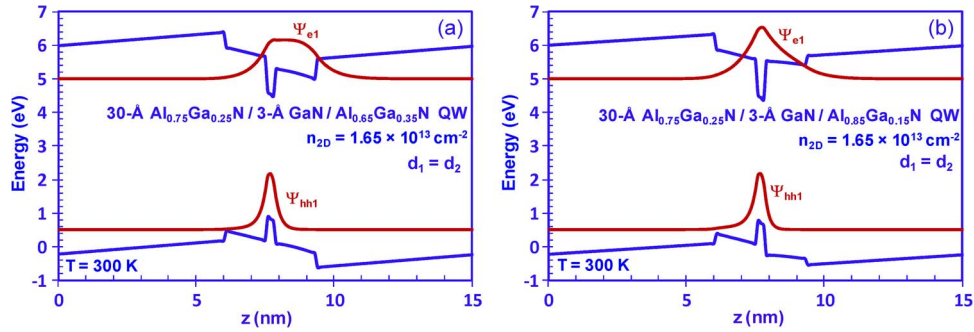


Fig. 3. Energy lineups with electron and hole wavefunctions for (a) the 30-Å $\text{Al}_{0.75}\text{Ga}_{0.25}\text{N}/3\text{-}\text{\AA}$ GaN/ $\text{Al}_{0.65}\text{Ga}_{0.35}\text{N}$ QW, and (b) the 30-Å $\text{Al}_{0.75}\text{Ga}_{0.25}\text{N}/3\text{-}\text{\AA}$ GaN/ $\text{Al}_{0.85}\text{Ga}_{0.15}\text{N}$ QW.

position (d_1, d_2) [Fig. 2(a) and (b)] or different compositions of AlGaN QW sublayers (x_1, x_2) [see Fig. 3(a) and (b)].

Fig. 2(a) and (b) shows the energy lineups with electron wavefunction EC1 and hole wave function HH1 for the 30-Å $\text{Al}_{0.8}\text{Ga}_{0.2}\text{N}/3\text{-}\text{\AA}$ GaN QW with different d_1 thicknesses ($d_1 + d_2 = 30$ Å) at the sheet carrier density (n_{2D}) of $9.9 \times 10^{12} \text{ cm}^{-2}$ ($T = 300$ K). For the 30-Å $\text{Al}_{0.8}\text{Ga}_{0.2}\text{N}/3\text{-}\text{\AA}$ GaN QW structure with $d_1 = 10$ Å in Fig. 2(a), the electron–hole wavefunction overlap ($\Gamma_{e,hh}$) is obtained as $\sim 63.51\%$. By engineering the delta QW structure with $d_1 = 18$ Å, as shown in Fig. 2(b), the electron and hole wavefunctions are strongly localized toward the center of the QW active region, attributing to the engineered spreading of the wavefunctions, which leads to the enhanced $\Gamma_{e,hh} \sim 73.92\%$.

Fig. 3(a) and (b) shows the energy band lineups with electron wavefunction EC1 and hole wave function HH1 for the 30-Å $\text{Al}_{0.75}\text{Ga}_{0.25}\text{N}/3\text{-}\text{\AA}$ GaN/ $\text{Al}_{0.65}\text{Ga}_{0.35}\text{N}$ QW and $\text{Al}_{0.75}\text{Ga}_{0.25}\text{N}/3\text{-}\text{\AA}$ GaN/ $\text{Al}_{0.85}\text{Ga}_{0.15}\text{N}$ QW with $n_{2D} = 1.65 \times 10^{13} \text{ cm}^{-2}$ at $T = 300$ K, and the thickness d_1 is kept as identical with the thickness d_2 . For the 30-Å $\text{Al}_{0.75}\text{Ga}_{0.25}\text{N}/3\text{-}\text{\AA}$ GaN/ $\text{Al}_{0.65}\text{Ga}_{0.35}\text{N}$ QW, the design consisting of asymmetric QW with different Al contents (75% and 65%) results in $\Gamma_{e,hh}$ of $\sim 60.31\%$. By engineering the Al contents of the AlGaN layers as $x_1 = 0.75$ and $x_2 = 0.85$ [see Fig. 3(b)], the electron and hole wavefunctions are pushed toward the center of the QW active region, as the band gap of $\text{Al}_{0.85}\text{Ga}_{0.15}\text{N}$ is larger than that of $\text{Al}_{0.75}\text{Ga}_{0.25}\text{N}$, which results in the improved $\Gamma_{e,hh} \sim 78.83\%$. Therefore, the matrix element will be enhanced attributing to the improved $\Gamma_{e,hh}$, which will contribute to the optical gain by using the optimized asymmetric AlGaN-delta-GaN QW structures.

3. Optical Gain Characteristics of Asymmetric AlGaN-Delta-GaN QWs

Fig. 4(a) shows the TE-polarized optical gain spectra for 30-Å $\text{Al}_{0.8}\text{Ga}_{0.2}\text{N}/3\text{-}\text{\AA}$ GaN QWs with $n_{2D} = 1.65 \times 10^{13} \text{ cm}^{-2}$ at $T = 300$ K with varying d_1 thicknesses. For the asymmetric QW with

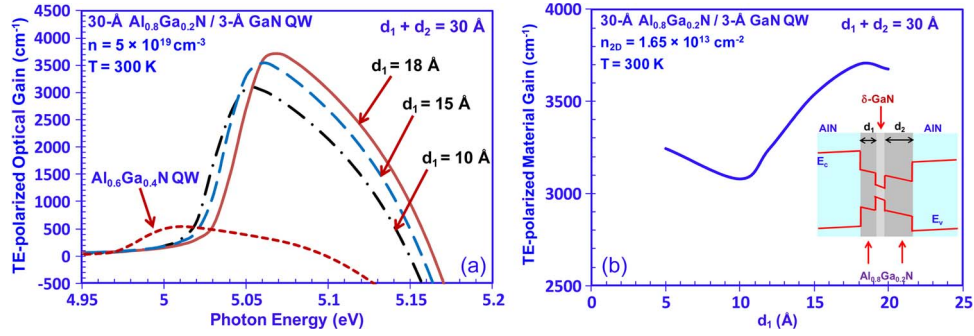


Fig. 4. (a) TE-polarized optical gain spectra for 30-Å $\text{Al}_{0.8}\text{Ga}_{0.2}\text{N}/3\text{-Å GaN}$ QW with varying d_1 thicknesses, and (b) TE-polarized material peak gain as a function of the d_1 thickness for the 30-Å $\text{Al}_{0.8}\text{Ga}_{0.2}\text{N}/3\text{-Å GaN}$ QW with $n_{2D} = 1.65 \times 10^{13} \text{ cm}^{-2}$ at $T = 300 \text{ K}$.

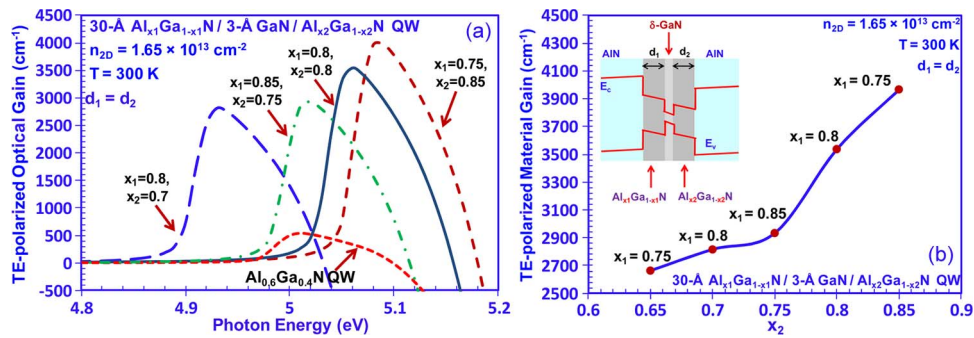


Fig. 5. (a) TE-polarized optical gain spectra for 30-Å $\text{Al}_{x_1}\text{Ga}_{1-x_1}\text{N}/3\text{-Å GaN}/\text{Al}_{x_2}\text{Ga}_{1-x_2}\text{N}$ QW with varying combinations of the Al contents x_1 and x_2 , and (b) TE-polarized material peak gain for the 30-Å $\text{Al}_{x_1}\text{Ga}_{1-x_1}\text{N}/3\text{-Å GaN}/\text{Al}_{x_2}\text{Ga}_{1-x_2}\text{N}$ QW with $n_{2D} = 1.65 \times 10^{13} \text{ cm}^{-2}$ at $T = 300 \text{ K}$.

$d_1 = 18 \text{ Å}$, the TE-polarized optical gain is larger than that of the symmetric QW with $d_1 = 15 \text{ Å}$, attributing to the enhanced $\Gamma_{e,hh}$ from the optimized AlGa_N-delta-GaN QW structure. For the asymmetric QW with $d_1 = 10 \text{ Å}$, the optical gain is slightly lower due to the weaker confinement of the electron and hole wavefunctions. Note that the peak emission wavelengths (λ_{peak}) for the asymmetric QWs with $d_1 = 10 \text{ Å}$ and 18 Å are $\sim 246 \text{ nm}$ and $\sim 245 \text{ nm}$, respectively, which are very similar with that of the symmetric QW. In comparison with the conventional $\text{Al}_{0.6}\text{Ga}_{0.4}\text{N}$ QW, the optimized asymmetric delta QW ($d_1 = 18 \text{ Å}$) shows ~ 6 times enhancement in the TE-polarized optical gain for mid-UV spectral regime.

Fig. 4(b) shows the TE-polarized material peak gain ($g_{\text{peak}}^{\text{TE}}$) as a function of the d_1 thickness for the 30-Å $\text{Al}_{0.8}\text{Ga}_{0.2}\text{N}/3\text{-Å GaN}$ QW with $n_{2D} = 1.65 \times 10^{13} \text{ cm}^{-2}$ at $T = 300 \text{ K}$. With the d_1 thicknesses ranging from 5 Å up to 20 Å , the TE-polarized material gain varies from $\sim 3078 \text{ cm}^{-1}$ ($d_1 = 10 \text{ Å}$) up to $\sim 3703 \text{ cm}^{-1}$ ($d_1 = 18 \text{ Å}$). Therefore, very large optical gain can be maintained for emission wavelength $\sim 245 \text{ nm}$ for the asymmetric AlGa_N-delta-GaN QW structures with different delta-GaN positions.

Fig. 5(a) illustrates the TE-polarized gain spectra for 30-Å $\text{Al}_{x_1}\text{Ga}_{1-x_1}\text{N}/3\text{-Å GaN}/\text{Al}_{x_2}\text{Ga}_{1-x_2}\text{N}$ QW with $n_{2D} = 1.65 \times 10^{13} \text{ cm}^{-2}$ at $T = 300 \text{ K}$ with combinations x_1 and x_2 . For the optimized delta QW with Al contents of $x_1 = 0.75$ and $x_2 = 0.85$, very large TE-polarized gain ($g_{\text{peak}}^{\text{TE}} \sim 3967 \text{ cm}^{-1}$) can be obtained at $\lambda_{\text{peak}} \sim 244 \text{ nm}$, attributed to the improved $\Gamma_{e,hh}$ from the asymmetric QW. For the asymmetric cases of $\text{Al}_{0.85}\text{Ga}_{0.15}\text{N}/3\text{-Å GaN}/\text{Al}_{0.75}\text{Ga}_{0.25}\text{N}$ QW ($x_1 = 0.85, x_2 = 0.75$) and $\text{Al}_{0.8}\text{Ga}_{0.2}\text{N}/3\text{-Å GaN}/\text{Al}_{0.7}\text{Ga}_{0.3}\text{N}$ QW ($x_1 = 0.8, x_2 = 0.7$), large optical gains can be obtained with $\lambda \sim 244\text{--}252 \text{ nm}$, which are $\sim 6\text{--}7$ times larger than that of the conventional QW emitting at similar wavelengths.

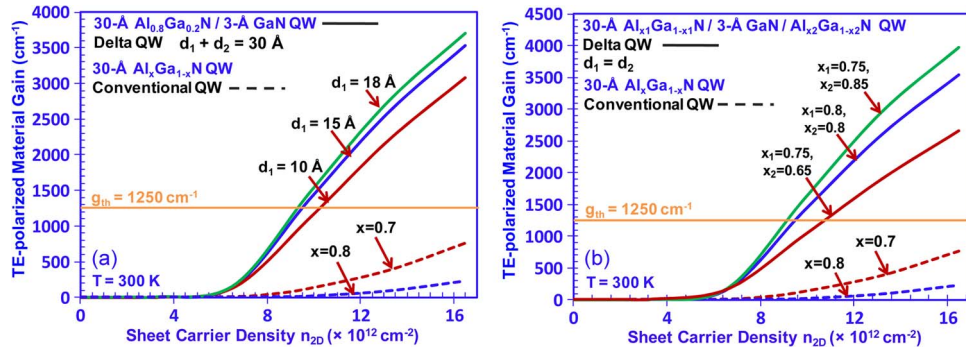


Fig. 6. TE-polarized material gain as a function of sheet carrier density for (a) 30-Å Al_{0.8}Ga_{0.2}N/3-Å GaN QW with varying d_1 thicknesses, and (b) 30-Å Al _{x_1} Ga _{$1-x_1$} N/3-Å GaN/Al _{x_2} Ga _{$1-x_2$} N QWs, in comparison with conventional 30-Å Al _{x} Ga _{$1-x$} N QW ($x = 0.7, 0.8$).

Fig. 5(b) shows the TE-polarized material peak gain for the 30-Å Al _{x_1} Ga _{$1-x_1$} N/3-Å GaN/Al _{x_2} Ga _{$1-x_2$} N QW with $n_{2D} = 1.65 \times 10^{13} \text{ cm}^{-2}$ at $T = 300 \text{ K}$. By varying the Al contents of the AlGa_N layers surrounding the delta-GaN layer, the TE-polarized material peak gains range from $\sim 2700 \text{ cm}^{-1}$ up to $\sim 4000 \text{ cm}^{-1}$. Therefore, the optimized asymmetric QW structures with different Al contents can lead to enhanced optical gain. Also, large TE-polarized gain can be maintained by different combinations of the Al contents of the AlGa_N QWs, which will provide flexibility in the experimental realizations of the AlGa_N-delta-GaN QWs mid- and deep-UV lasers.

4. Threshold Properties and Differential Gains of AlGa_N-Delta-GaN QWs

Fig. 6(a) shows the TE-polarized material gain as a function of sheet carrier density for both 30-Å Al_{0.8}Ga_{0.2}N/3-Å GaN QWs with various d_1 thicknesses, and conventional 30-Å Al _{x} Ga _{$1-x$} N QW ($x = 0.7, 0.8$) at $T = 300 \text{ K}$. The $g_{\text{peak}}^{\text{TE}}$ values of the AlGa_N-delta-GaN structure ($\sim 3100\text{--}3700 \text{ cm}^{-1}$) are found to be significantly larger than that of the conventional high Al-content AlGa_N QWs ($\sim 200\text{--}800 \text{ cm}^{-1}$) at high $n_{2D} = 1.65 \times 10^{13} \text{ cm}^{-2}$. The optimized asymmetric QW structure with $d_1 = 18 \text{ Å}$ achieves ~ 1.2 times larger material gain than that of the symmetric QW structure, attributing to its improved matrix element.

To illustrate the effect of different Al contents of the AlGa_N QWs, Fig. 6(b) shows the TE-polarized material gain as a function of sheet carrier density for 30-Å Al _{x_1} Ga _{$1-x_1$} N/3-Å GaN/Al _{x_2} Ga _{$1-x_2$} N QWs and conventional 30-Å Al _{x} Ga _{$1-x$} N QWs ($x = 0.7, 0.8$) at $T = 300 \text{ K}$. Attributing to the reduced charge separation effect, the optimized asymmetric 30-Å Al_{0.75}Ga_{0.25}N/3-Å GaN/Al_{0.85}Ga_{0.15}N QW achieves ~ 1.4 times larger material gain than that of the symmetric QW. Thus, the TE-polarized lasing is feasible for asymmetric AlGa_N-delta-GaN QW lasers with $\lambda \sim 240\text{--}250 \text{ nm}$.

To analyze the threshold properties of mid- and deep-UV lasers, AlGa_N QW lasers with optical confinement factor (Γ_{opt}) of 0.02 [41] were used in the analysis. Based on the transfer matrix method [64], the Γ_{opt} for the asymmetric and symmetric QWs are calculated as almost identical. Thus, the modal gain comparison for the QWs will be governed by the difference in the material gains. The internal loss (α_i) in typical AlGa_N lasers is 14 cm^{-1} [35]. The laser cavity length is assumed as $500 \mu\text{m}$ [35], [41] with end-facet reflectivities of 95% and 35%, which corresponds to mirror loss (α_m) of 11 cm^{-1} and threshold gain (g_{th}) of $\sim 1250 \text{ cm}^{-1}$. The threshold sheet carrier density (n_{2D}^{th}) is $9.273 \times 10^{12} \text{ cm}^{-2}$ for the symmetric 30-Å Al_{0.8}Ga_{0.2}N/3-Å GaN QW. For the optimized structure with $d_1 = 18 \text{ Å}$ [see Fig. 6(a)], the n_{2D}^{th} is $9.042 \times 10^{12} \text{ cm}^{-2}$. Similarly, for optimized 30-Å Al_{0.75}Ga_{0.25}N/3-Å GaN/Al_{0.85}Ga_{0.15}N QW, the n_{2D}^{th} is $8.976 \times 10^{12} \text{ cm}^{-2}$. Thus, the n_{2D}^{th} for both optimized AlGa_N-delta-GaN QW structures with different delta-GaN positions and varying Al contents are reduced compared with the symmetric QW.

The spontaneous emission rates and peak modal gains (with $\Gamma_{\text{opt}} = 0.02$ [41]) comparisons for both symmetric and asymmetric AlGa_N-Delta-GaN QWs are presented in Figs. 7 and 8, respectively. The comparison studies were performed for AlGa_N-Delta-GaN QWs with different delta layer

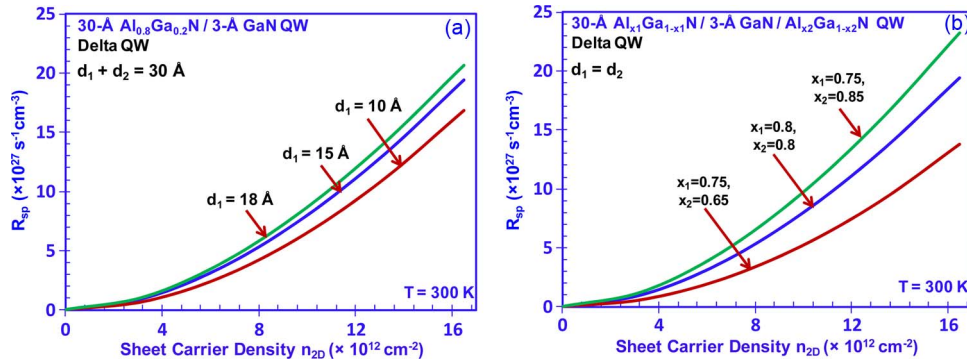


Fig. 7. R_{sp} as a function of sheet carrier density for both symmetric and asymmetric AlGa_xN-Delta-GaN QWs to illustrate (a) effect of delta positions, and (b) effect of Al contents in AlGa_xN QWs.

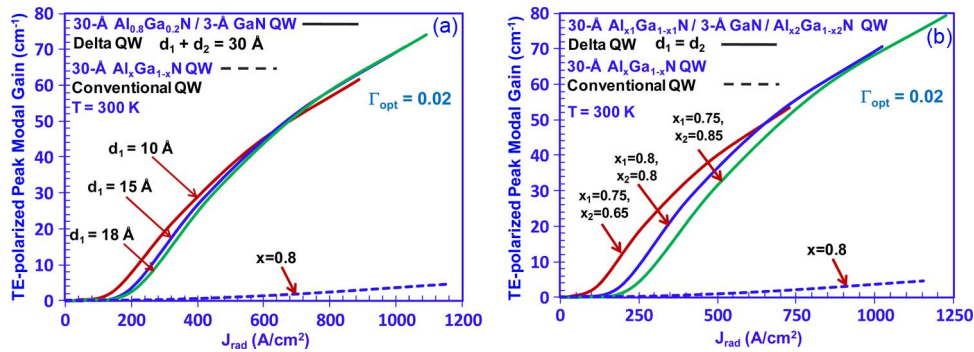


Fig. 8. Modal gains as a function of radiative current density for AlGa_xN-Delta-GaN QWs and conventional Al_{0.8}Ga_{0.2}N QW with different (a) delta positions, and (b) Al contents in AlGa_xN QWs.

position (d_1, d_2) [see Figs. 7(a) and 8(a)] and different AlGa_xN QW sublayer compositions [see Figs. 7(b) and 8(b)]. By comparing the R_{sp} for different delta-layer positions [see Fig. 7(a)], the asymmetric delta QW with $d_1 = 18 \text{ \AA}$ shows higher R_{sp} ($\sim 2.07 \times 10^{28} \text{ s}^{-1} \text{ cm}^{-3}$ with $n_{2D} = 1.65 \times 10^{13} \text{ cm}^{-2}$) than those of the other structures, attributed from the improved overlap ($\Gamma_{e,hh} \sim 73.92\%$). Fig. 7(b) shows that the optimized R_{sp} was obtained from the use of that for Al_{0.75}Ga_{0.25}N / 3-Å GaN / Al_{0.85}Ga_{0.15}N QWs attributed to the optimized C-HH overlap of 78.83%. The delta QWs exhibited significantly higher modal gain for any radiative current density (J_{rad}) injection level [see Fig. 8(a) and (b)], in comparison with those the conventional QW. The total recombination current density J_{tot} in the QW active region includes both the radiative and nonradiative current densities ($J_{tot} = J_{rad} + J_{non-rad}$) [58], and $J_{non-rad}$ ($\sim A \cdot n_{th} + C \cdot n_{th}^3$) represents the dominant part of the J_{tot} in AlGa_xN-based QW [41], [58]. Thus, the reduction in n_{th} is important for suppressing monomolecular ($\sim A \cdot n_{th}$) and Auger ($\sim C \cdot n_{th}^3$) recombination currents at threshold.

The differential gain properties are also analyzed for both symmetric and asymmetric AlGa_xN-Delta-GaN QWs to illustrate effect of delta positions [see Fig. 9(a)], as well as effect of Al contents [see Fig. 9(b)]. For both cases, the differential gains first increase with higher carrier densities, which can be attributed to the carrier screening effect. After reaching the maximum, the differential gains start to decrease with increasing carrier densities, which is due to the band filling effect. In Fig. 9(a) and (b), the peak differential gains of the optimized asymmetric AlGa_xN-delta-GaN QW structures are higher than those of the symmetric QW structures, which indicates that the optimized AlGa_xN-delta-GaN QW structures will be applicable for high-speed modulation lasers.

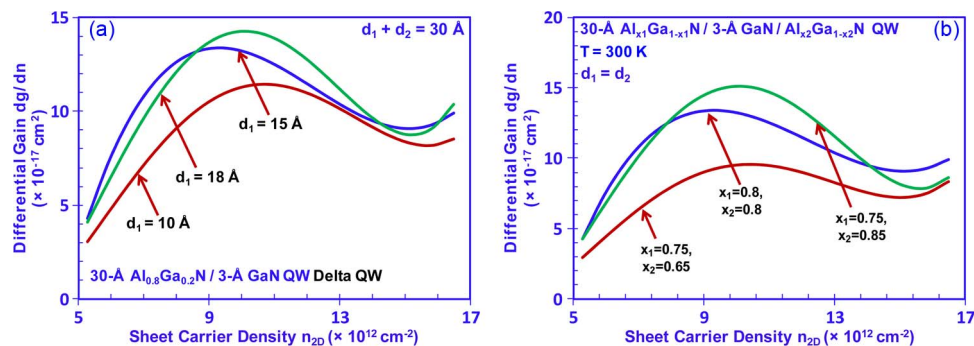


Fig. 9. Differential gain as a function of sheet carrier density for AlGa_N-Delta-GaN QWs to illustrate the effect of (a) delta positions, and (b) Al contents of AlGa_N QWs

5. Summary

In summary, the comprehensive optimization studies on the gain characteristics of AlGa_N-delta-GaN QWs with various delta-GaN positions and AlGa_N compositions are analyzed. The optimized asymmetric AlGa_N-delta-GaN QWs result in ~ 7 times increase in material gain, in comparison with that of the conventional QWs. Recent works had reported the TE-polarized gain from the following: 1) Al_{0.7}Ga_{0.3}N/AlN QWs with low Al-content nanocluster “quantum dot” features within the Al_{0.7}Ga_{0.3}N QW layers [65], [66] and 2) thin GaN “quantum dot” structures embedded in AlN barriers [67], which are consistent with the dominant TE-polarized gain from our prediction for high Al-content AlGa_N QWs consisting of low band-gap GaN delta layer embedded in the active region [51]. Large material gains can be maintained at $\lambda \sim 240\text{--}250 \text{ nm}$ for the asymmetric AlGa_N-delta-GaN QWs [see Fig. 4(b)]. Despite the improved material gain for the asymmetric QWs, the finding shows that large material gain can be obtained for both symmetric and asymmetric AlGa_N-delta-GaN QWs [see Fig. 4(b)], which indicates the flexibility and robustness in the experimental implementations of this concept in devices. Thus, by employing asymmetric QW design with optimized delta-GaN layer position and asymmetric AlGa_N-composition layers, the optimized optical gain and lower threshold carrier densities are achievable for the AlGa_N-delta-GaN QW with realistic design applicable for mid- and deep-UV lasers. Recent works by metalorganic chemical vapor deposition (MOCVD) had reported the growths of AlN/ GaN superlattice structures with thin GaN layers in the order of $\sim 0.9\text{--}2.5$ monolayers [52]. Thus, the growths of the AlGa_N-delta-GaN QW is expected to be practical for implementation in deep/mid-UV LEDs or lasers.

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