

Four-Component Scattering Power Decomposition With Extended Volume Scattering Model

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Abstract—In the three- or four-component decompositions, polarimetric scattering properties and corresponding physical scattering models play essential roles for power decomposition. This letter proposes an improved four-component scattering power decomposition method that employs a suitable volume scattering model for single- or double-bounce scattering in the polarimetric synthetic aperture radar image analysis. The cross-polarized HV component is created by both single-bounce object (such as vegetation) and double-bounce structures (such as oriented building blocks). It has been difficult to discriminate these two objects (vegetation against oriented buildings) in the decomposed images since the HV component is assigned to the volume scattering due to vegetation only. We propose to extend the volume scattering model suited for two physical scattering models. It is shown that a vegetation area and an oriented urban building area are well discriminated compared to those resulting from the implementation of the existing four-component scattering power decomposition.

Index Terms—Polarimetric synthetic aperture radar (SAR) (POLSAR), radar polarimetry, scattering power decomposition, volume scattering.

I. INTRODUCTION

POLARIMETRIC scattering power decompositions based on physical scattering models have been attracting attention for target classification, detection, and land parameter retrieval for fully polarimetric synthetic aperture radar (SAR) (POLSAR) data analysis [1]–[4]. Model-based approaches are simple and straightforward to implement on fully POLSAR data take analyses [5], [6]. The three-component [3] and four-component [4], [5] decompositions are well known in the case of employing physical scattering models as typical target classification and detection. There are nine real independent polarimetric parameters in the coherency or covariance matrix for the most general scattering case. The four-component scattering power decomposition method [4] accounts for six terms out of nine parameters by adding a helicity term to the three-component method [3]. After decomposing of total power into four components, it has been found that the method [4] has problems of overestimation in the volume scattering in urban areas and in negative power occurrence in the processing caused by a large cross-polarization component [5], [6]. Recently, Yamaguchi *et al.* [6] have proposed a rotation method of the coherency matrix for more accurate POLSAR image decomposition and target classification. This is essentially identical with

the deorientation method [7], [8] and the polarimetric orientation angle compensation method [9]. In addition, Arie *et al.* has recently proposed an adaptive model-based decomposition method [10] to deal with volume scattering, which ensures nonnegative power decomposition.

This rotation of coherency matrix significantly improved the decomposition results by minimizing the cross-polarized (HV) components and reduced the negative power occurrence in the surface scattering and the double-bounce scattering components. The scattering powers are calculated easily and are used to compose full color images with red–green–blue color coding, red for the double-bounce power, green for the volume scattering power, and blue for the surface scattering power, for which each color brightness corresponds to the magnitude. They have been successfully applied to POLSAR image analysis because color-coded images are easier to understand and because each color represents a specific scattering mechanism. This method accounts for six parameters out of eight independent observable polarimetric parameters.

However, there still exists a problem of discrimination between vegetation and oriented buildings within the same volume scattering (green) area. This ambiguity is caused by the assignment of the HV component. So far, the volume scattering power is evaluated by the HV component due to vegetation only in the existing methods [3]–[6]. To resolve the discrimination ambiguity, we have to assign the HV component whether it is created by vegetation (single-bounce object) or by edges of oriented buildings (double-bounce structures). In this letter, we propose to use a new volume scattering model that accounts for the HV component caused by double-bounce structures versus vegetation scatter and to improve the four-component scattering power decomposition [6] for additional more accurate classification.

II. FOUR-COMPONENT SCATTERING DECOMPOSITION WITH ROTATION OF COHERENCY MATRIX

Once the scattering matrix $[S]$ is acquired from fully polarimetric radar data sets, the Pauli vector can be defined as

$$\mathbf{k}_p = \frac{1}{\sqrt{2}} \begin{bmatrix} S_{HH} + S_{VV} \\ S_{HH} - S_{VV} \\ 2S_{HV} \end{bmatrix} \quad (1)$$

where S_{HH} , S_{VV} , and S_{HV} are elements of the scattering matrix, assuming the reciprocal condition of $S_{HV} = S_{VH}$.

The coherency matrix is given as

$$\langle [T] \rangle = \langle \mathbf{k}_p \mathbf{k}_p^\dagger \rangle = \begin{bmatrix} T_{11} & T_{12} & T_{13} \\ T_{21} & T_{22} & T_{23} \\ T_{31} & T_{32} & T_{33} \end{bmatrix} \quad (2)$$

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where \dagger denotes the complex conjugation and transposition, and $\langle \cdot \rangle$ denotes the ensemble average in an imaging window.

The coherency matrix after rotation by angle θ can be obtained as

$$\langle [T'] \rangle = [R(\theta)] \langle [T] \rangle [R(\theta)]^\dagger = \begin{bmatrix} T'_{11} & T'_{12} & T'_{13} \\ T'_{21} & T'_{22} & T'_{23} \\ T'_{31} & T'_{32} & T'_{33} \end{bmatrix} \quad (3)$$

where

$$[R(\theta)] = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos 2\theta & \sin 2\theta \\ 0 & -\sin 2\theta & \cos 2\theta \end{bmatrix}. \quad (4)$$

The rotation angle θ is determined so as to minimize the T_{33} or, equivalently, the HV component [6]

$$\theta = \frac{1}{4} \tan^{-1} \left(\frac{2\text{Re}\{T_{23}\}}{T_{22} - T_{33}} \right). \quad (5)$$

The rotated measured coherency matrix $\langle [T'] \rangle$ is expanded into four submatrices which correspond to surface, double-bounce, volume, and helix scattering mechanisms [6]

$$\langle [T'] \rangle = f_s \langle [T] \rangle_{\text{surface}} + f_d \langle [T] \rangle_{\text{double}} + f_\nu \langle [T] \rangle_{\text{vol}} + f_c \langle [T] \rangle_{\text{helix}} \quad (6)$$

where f_s , f_d , f_ν , and f_c are the coefficients to be determined. $\langle [T] \rangle_{\text{surface}}$, $\langle [T] \rangle_{\text{double}}$, $\langle [T] \rangle_{\text{vol}}$, and $\langle [T] \rangle_{\text{helix}}$ are expansion matrices corresponding to the surface, double-bounce, volume, and helix scattering mechanisms, respectively [2]–[5].

The expansion matrices for $\langle [T] \rangle_{\text{surface}}$, $\langle [T] \rangle_{\text{double}}$, and $\langle [T] \rangle_{\text{helix}}$ are given in [6], and not repeated here.

For the volume scattering mechanism, we must confirm from the experimental evidence that $\text{Re}\{\langle S_{HH}S_{VV}^* \rangle\} > 0$ is corresponding to the surface scattering. Under the condition $\text{Re}\{\langle S_{HH}S_{VV}^* \rangle\} > 0$, the HV component is assigned to a cloud of randomly oriented dipole scatterers such as vegetation. We choose one of the following matrices according to the magnitude balance of $\langle |S_{HH}|^2 \rangle$ and $\langle |S_{VV}|^2 \rangle$.

For $10 \log(\langle |S_{VV}|^2 \rangle / \langle |S_{HH}|^2 \rangle) > 2$ dB

$$\langle [T] \rangle_{\text{vol}}^{\text{dipole}} = \frac{1}{30} \begin{bmatrix} 15 & -5 & 0 \\ -5 & 7 & 0 \\ 0 & 0 & 8 \end{bmatrix}. \quad (7)$$

For $|10 \log(\langle |S_{VV}|^2 \rangle / \langle |S_{HH}|^2 \rangle)| < 2$ dB

$$\langle [T] \rangle_{\text{vol}}^{\text{dipole}} = \frac{1}{4} \begin{bmatrix} 2 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}. \quad (8)$$

For $10 \log(\langle |S_{VV}|^2 \rangle / \langle |S_{HH}|^2 \rangle) < -2$ dB

$$\langle [T] \rangle_{\text{vol}}^{\text{dipole}} = \frac{1}{30} \begin{bmatrix} 15 & 5 & 0 \\ 5 & 7 & 0 \\ 0 & 0 & 8 \end{bmatrix}. \quad (9)$$

However, for the case of $\text{Re}\{\langle S_{HH}S_{VV}^* \rangle\} < 0$, we have to use a different alternate expansion matrix as shown in the next section.

The corresponding scattering powers (the surface scattering power P_s , the double-bounce scattering P_d , the volume scattering power P_ν , and the helix scattering power P_c) are directly obtained from the expansion coefficients by applying

the decomposition algorithm. Details of this method are given in [6].

III. NEW FOUR-COMPONENT SCATTERING POWER DECOMPOSITION

In this section, a new four-component scattering power decomposition method is explained using a new volume scattering model.

A. New Expansion Matrix for Dihedral Scattering

The experimental evidence indicates that $\text{Re}\{\langle S_{HH}S_{VV}^* \rangle\} < 0$ corresponds to the double-bounce scattering case. This situation typically occurs for man-made objects with right angle structures such as building walls and roads, and river surfaces and bridges orthogonal to radar illumination. We call this kind of structure dihedral structures. When the directions of the main surface of buildings are oriented with respect to radar illumination, the HV component is generated and becomes a relatively large term in the acquired scattering matrix. We must consider this physical situation in more details and must incorporate the HV component by the oriented dihedral structures for more accurate modeling.

Since the rotation of coherency matrix minimizes the HV component [6], most of the orientations of dihedral structures are centered about zero degree with respect to the direction of radar illumination after the rotation (3). In order to derive a new expansion coherency matrix for the HV component, we implement ensemble averaging of dihedral corner reflectors using a probability density function $p(\theta)$ with its peak at zero degree,

$$p(\theta) = \frac{1}{2} \cos \theta, \quad \text{for } -\frac{\pi}{2} < \theta < \frac{\pi}{2}. \quad (10)$$

The theoretical ensemble matrix for a dihedral corner reflector can be derived from

$$\langle [T(\theta)] \rangle = \int_{-\pi/2}^{\pi/2} [T(\theta)] p(\theta) d\theta \quad (11)$$

yielding

$$\langle [T] \rangle_{\text{vol}}^{\text{dihedral}} = \frac{1}{15} \begin{bmatrix} 0 & 0 & 0 \\ 0 & 7 & 0 \\ 0 & 0 & 8 \end{bmatrix}. \quad (12)$$

This matrix was derived using (3) for dihedral in a similar way explained in [4] and was set so that the trace becomes unity. We use this matrix as a new volume scattering expansion matrix that accounts for the HV component for dihedral structures.

B. Branch Condition

After the rotation of coherency matrix, we first discriminate the scattering mechanism using the sign of $\text{Re}\{\langle S_{HH}S_{VV}^* \rangle\}$. It is known from the experimental evidence that the double scattering by dihedrals causes the parameter $\text{Re}\{\langle S_{HH}S_{VV}^* \rangle\}$ to be negative. On the other hand, a nondihedral structure causes $\text{Re}\{\langle S_{HH}S_{VV}^* \rangle\} > 0$. These relations can be explained in more details by rigorously using the covariance matrix formulation [4]. By the expansion of the C_{13} component for

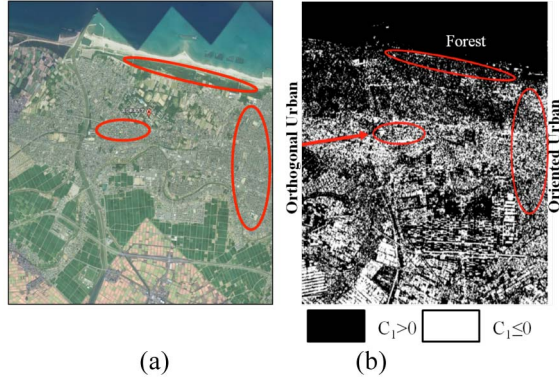


Fig. 1. (a) Google Earth optical image. (b) C_1 binary image of Pi-SAR data over Niigata University and the surrounding area derived by (17).

randomly distributed dipoles [4], we can derive the equation as follows:

$$\text{Re}\{f_s\beta + f_d\alpha^*\} + \frac{1}{8}f_\nu - \frac{1}{4}f_c = \text{Re}\{\langle S_{HH}S_{VV}^* \rangle\}. \quad (13)$$

This equation can be rearranged to

$$\begin{aligned} C_1 &= 2\text{Re}\{f_s\beta + f_d\alpha^*\} \\ &= T'_{11} - T'_{22} + \frac{1}{2}f_c. \end{aligned} \quad (14)$$

We have omitted the term $(1/8)f_\nu$ in (14) because the volume scattering coefficient is not decided at this stage. The sign of (14) determines the dominant scattering mechanism, i.e., surface scattering versus double-bounce scattering. According to the sign of C_1 , we assign the volume scattering (the HV component) to surface scattering (vegetation) or double bounce scattering (oriented dihedral structure) according to the following condition:

$$C_1 > 0 : \text{volume scattering by vegetation} \quad (15)$$

$$C_1 \leq 0 : \text{volume scattering by dihedral.} \quad (16)$$

It should be noted that (15) and (16) are the first-stage criteria. It happens that (16) assigns the volume scattering as dihedral scattering when double bounce by forest trunk and ground is too strong in vegetation area. However, it is easy to recognize vegetation scattering area in this case because the tree trunk and ground scattering (double bounce) points appear randomly and sparsely in the surrounding area and the volume scattering dominates in the final decomposed forest image.

To check the validity of the criteria, L-band Polarimetric and Interferometric SAR (Pi-SAR) data sets have been analyzed. These Pi-SAR data sets were acquired over Niigata University and the surrounding area. Fig. 1(a) shows the optical Google Earth image over the area. Fig. 1(b) shows the C_1 binary image of Pi-SAR data, corresponding to Fig. 1(a). We can clearly see $C_1 > 0$ in pine forest areas in the upper part of Fig. 1(a) and $C_1 \leq 0$ in orthogonal urban areas (middle left) to radar illumination direction. On the other hand, we see a mixture of $C_1 > 0$ and $C_1 \leq 0$ in oriented urban area (right). This result confirms the validity of the criteria, which will be applied in a new decomposition method.

C. New Four-Component Scattering Power Decomposition

New decomposition can be explained in (17) and in Fig. 2 by using the volume scattering component from vegetation and/or

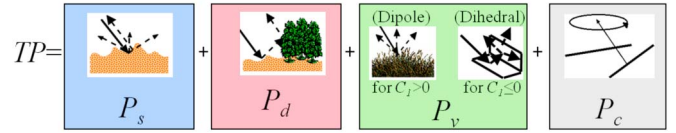


Fig. 2. New four-component scattering mechanism with rotation of coherency matrix.

oriented dihedral structures

$$\begin{aligned} \langle [T'] \rangle &= f_s \langle [T] \rangle_{\text{surface}} + f_d \langle [T] \rangle_{\text{double}} + f_c \langle [T] \rangle_{\text{helix}} \\ &+ \begin{cases} f_\nu \langle [T] \rangle_{\text{vol}}^{\text{dipole}}, & \text{for } C_1 > 0 \\ f_{\nu d} \langle [T] \rangle_{\text{vol}}^{\text{dihedral}}, & \text{for } C_1 \leq 0 \end{cases} \end{aligned} \quad (17)$$

where f_s , f_d , f_ν , $f_{\nu d}$, and f_c are coefficients to be determined. $\langle [T] \rangle_{\text{surface}}$, $\langle [T] \rangle_{\text{double}}$, and $\langle [T] \rangle_{\text{helix}}$ are expansion matrices corresponding to surface, double-bounce, and helix scattering mechanisms, respectively. $\langle [T] \rangle_{\text{vol}}^{\text{dihedral}}$ and $\langle [T] \rangle_{\text{vol}}^{\text{dipole}}$ are the volume scattering matrices for oriented dihedral structure and wire (ensemble average of dipole), respectively.

The total power (TP) decomposes into surface scattering power P_s , double-bounce scattering power P_d , volume scattering power P_ν , from dipole and/or oriented dihedral, and helix scattering power P_c , as shown in Fig. 2.

D. Decomposition Algorithm

The new decomposition algorithm is shown in Fig. 3. The first important step is the rotation of coherency matrix to minimize the HV component before the decomposition. It should be noted that (5) assumes $\arctan 2$ for obtaining the rotation angle in the computer algorithm. This rotation forces the real part of T_{23} to be zero, so that it ensures a reduction of the number of independent polarimetric parameters from nine to eight. Then, we check the sign of C_1 to assign the prevalent scattering mechanism. Once assigned to double-bounce scattering, we use a new expansion matrix (12) for volume scattering. Since most of the double-bounce structures are faced to radar illumination direction by implementation of the rotation of coherency matrix, the HV component comes from this dihedral structure.

On the other hand, if surface scattering is assigned, we use the same decomposition procedure as derived in [6]. Once the character of volume scattering power has been determined, it is possible to critically determine the dominant scattering mechanism within the volume scattering generated due to dipoles at this stage. We check the second branch condition using (13) to confirm the scattering mechanism again

$$2\text{Re}\{f_s\beta + f_d\alpha^*\} = 2\text{Re}\{\langle S_{HH}S_{VV}^* \rangle\} - 2\langle |S_{HV}|^2 \rangle + P_c. \quad (18)$$

This expression is equivalent to the following equation in terms of coherency matrix elements:

$$C_0 = T'_{11} - T'_{22} - T'_{33} + P_c. \quad (19)$$

The sign of C_0 determines the dominant scattering mechanism precisely within the volume scattering

$$\begin{aligned} C_0 > 0 &: \text{surface scattering} \\ C_0 \leq 0 &: \text{double-bounce scattering.} \end{aligned} \quad (20)$$

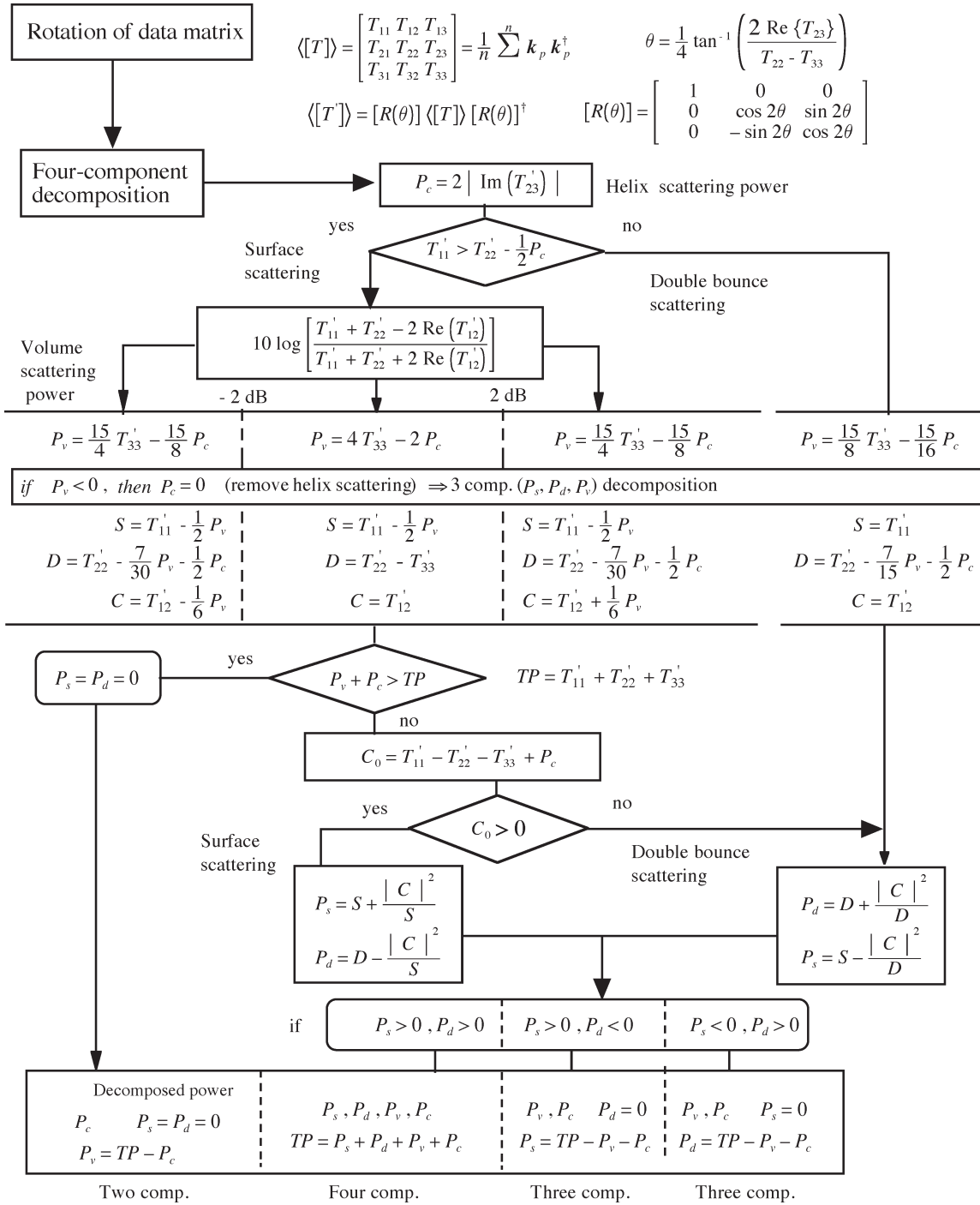


Fig. 3. Four-component decomposition with a rotation about the line of sight to set $\operatorname{Re}(T_{23}) = 0$ and the remaining HV contribution due to scattering from either a dihedral or a dipole distribution.

All of these physical branch conditions are included in a flowchart of the decomposition algorithm in Fig. 3.

IV. DECOMPOSITION RESULTS

Both existing [6] (four-component decomposition with rotation of coherency matrix) and the proposed procedures of decomposition methods have been applied to Pi-SAR images collected over Niigata University and environs. The volume scattering component due to [6] and the proposed method are shown in Fig. 4. Using the proposed extended volume scattering

model, significant reduction of volume scattering can be seen in dihedral structures (oriented urban and orthogonal urban areas) compared to the method developed in [6]. This fact serves for regulating the overestimation problem of the volume scattering in urban areas. On the other hand, the volume scattering remains almost the same in forested vegetation areas. Therefore, the proposed dihedral volume scattering model is expected to contribute for improving the decomposition accuracy in volume scattering.

Fig. 5 shows the decomposed color-coded composite images by both the existing and the proposed new methods for the sake

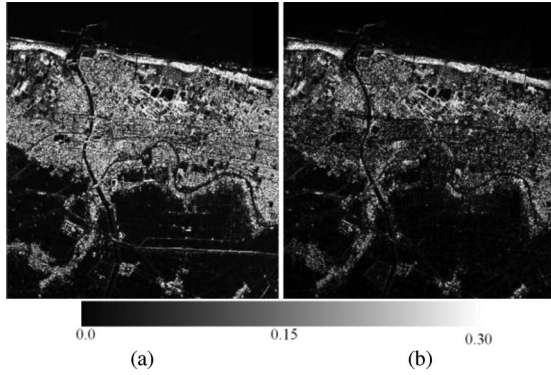


Fig. 4. (a) Volume scattering component derived from [6]. (b) Volume scattering component derived by the new method.

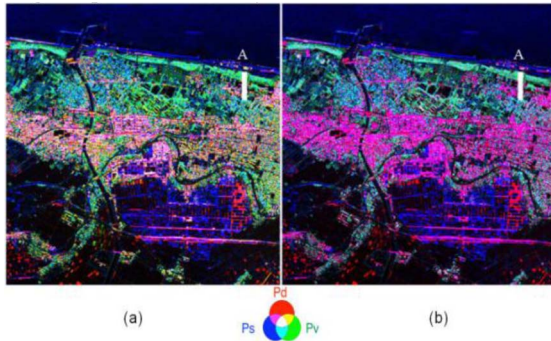


Fig. 5. Color-coded decomposed image of the Niigata University area by L-band Pi-SAR polarimetric data. (a) By the existing method [6]. (b) By the proposed method.

of comparison. The differences are clearly seen between the images in Fig. 5(a) and (b). We can see more “red” in Fig. 5(b) than in Fig. 5(a). This indicates that the dihedral structures are enhanced, while “green” caused by volume scattering within dihedral structures is reduced. In order to examine these results quantitatively, the decomposed power profiles along a transect over sandy ground, forest, and urban area [white line A in Figs. 5 and 6(a)] are shown in Fig. 6(b) and (c). The dashed lines in Fig. 6(b) and (c) show the boundaries of sand, forest, and urban areas. It is seen that the proposed method does not change the amount of volume scattering in sandy and forest areas as compared to the method in [6]. The presented new method reduces the amount of volume scattering power in urban areas where many residential houses exist, as shown in the red circle in Fig. 6(b) and (c).

V. CONCLUSION

In this letter, we have proposed a new volume scattering model that accounts for the HV component caused by double-bounce structures. This model better describes the HV component induced by rotated dihedral scattering and reduces the volume scattering power and enhances the double-bounce scattering power within man-made structures, leading to an improvement in the four-component scattering power decomposition. This decomposition method accounts for six terms out of eight real independent polarimetric parameters for the most general scattering case with the help of rotation of coherency

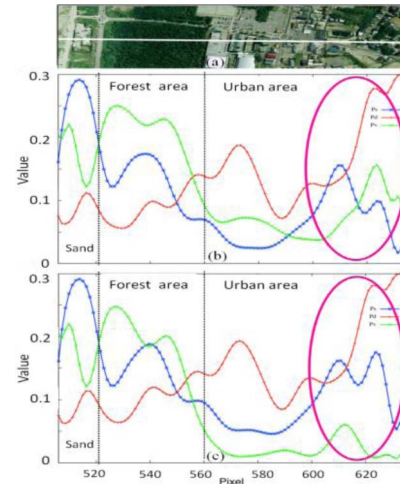


Fig. 6. (a) Photograph by Google Earth. (b) Profile of scattering components derived based on the existing method [6]. (c) Profile of scattering components derived based on the proposed method.

matrix. The decomposition results are in good agreement with Google Earth optical images.

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