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## Transient UV and Visible Luminescent Dynamics of Si-Rich SiOx Metal–Oxide–Semiconductor Light-Emitting Diodes

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Abstract: The transient ultraviolet (UV) and visible luminescent dynamics of the metal– oxide–semiconductor light-emitting diodes (MOSLEDs) made on Si-rich  $SiO_x$  with its O/Si composition ratio detuning from 0.75 to 1.62 are investigated. The size and luminescent wavelength of the buried Si quantum dots (Si-QDs) are controlled by adjusting the O/Si composition ratio of the Si-rich  $SiO<sub>x</sub>$ . Time-resolved photoluminescence shows a lifetime decaying from 11.5  $\mu$ s to 67 ns with the Si-QD size reducing from 4.5 to > 1.7 nm. The shorter lifetime for smaller Si-QDs is due to the increased nonphonon-assisted carrier recombination rate in smaller Si-QDs. The Si-QD size shrinkage is obtained by enlarging the O/Si composition ratio via the increase in the  $N_2O/SiH_4$  fluence ratio during synthesis, which makes the  $SiO<sub>x</sub>$  matrix approaching a standard dioxide with a higher turn-on threshold field under Fowler–Nordheim tunneling. By increasing the O/Si composition ratio from 1.15 to 1. 54, the obtained EL pattern changes its color from red to blue, which is associated with the turn-on voltage increasing from 40 to 175 V. Decreasing the Si-QD size to 1.7 nm inevitably attenuates the EL power to 100 nW and reduces the P/I slope to 0.63 mW/A. The UV EL patterns of MOSLEDs made on the  $SiO<sub>1.62</sub>$  film are demonstrated with an EL power of 40 nW, and the decay of ITO transmittance to  $<$  30% at an EL wavelength of  $<$  375 nm also contributes to the power attenuation.

Index Terms: Silicon nanophotonics, quantum dots and single molecules, light-emitting diodes.

## 1. Introduction

Near-infrared (NIR) photoluminescence (PL) and electroluminescence (EL) of Si quantum dots (Si-QDs) have been extensively studied due to their potential applications in Si-based light-emitting devices (LEDs) for the optical interconnects between electronic integrated circuits [1]–[6]. Up to now, Si-QDs can only be generated in a  $SiO<sub>x</sub>$  host matrix by employing a high-temperature annealing process. Indeed, the high-temperature annealing needed to form the Si-QDs in  $SiO<sub>x</sub>$  $(SiO<sub>x</sub>:Si-QDs)$  films is somehow a puzzling problem with regard to the processing compatibility of the complementary metal–oxide–semiconductor (CMOS) technology. Nevertheless, there are new observations in recent works that the as-grown Si-QDs can be embedded in Si-rich SiN<sub>x</sub>-based LEDs without annealing [7], [8]. Although different synthesizing and annealing conditions have been comprehensively studied for improving the light emission from a  $Si$ -rich  $SiO<sub>x</sub>$  film, plasma-enhanced chemical vapor deposition (PECVD) is still the most intriguing method to produce dense Si-QDs in a robust and stable  $SiO<sub>2</sub>$  matrix. To date, the record for the operational stability of nanoporous Si-based LEDs of up to 2 h has been reported by Gelloz et al. [9]. In comparison, high-temperature annealed Si-rich  $SiO<sub>x</sub>$  with size-tunable Si-QDs has also shown a comparable lifetime of operation, which is due to the relatively stable and saturated oxygen environment within the surrounding  $SiO<sub>2</sub>$ matrix [10].

In addition, versatile synthesizing and annealing recipes have also been proposed to obtain band gap tunable Si-QD and to improve the internal and external quantum efficiencies [11]–[14]. Previously, a Si nanostructure with a PL-related internal quantum efficiency (IQE) of 30%–50% has been measured by Atwater's group [11], [12]. Moreover, the maximum IQE of Si-QDs embedded in a SiO $_2$  host matrix is as high as 59  $\pm$  9% [12]. A PL-related IQE of similar Si nanocrystals of several tens of percent has also been reported by the other groups [13], [14]. However, the EL-related external quantum efficiency (EQE) reported in previous works is much lower than the PL-related IQE. In the  $SiO_x:Si-QD-based metal–oxide–semiconductor light-emitting diodes (MOSLEDs), most$ of the electron–hole pairs contribute to the EL are injected into the Si-QDs through either impact excitation or the Fowler–Nordheim (F-N) tunneling process [15]. Although electron–hole recombination provides relatively similar EL and PL spectra, the EL-related EQE suffers from a limited tunneling probability of the carriers when passing through the  $SiO<sub>x</sub>$  dielectric matrix [15]. As a result, only a few groups have reported an acceptable EL-related EQE ranging between  $10^{-5}$  and  $10^{-3}$  for Si-QD embedded MOSLEDs [4], [6], [16]–[18].

Nowadays, the fabricating technology of the MOSLEDs with red and infrared EL by tuning the Si-QD size is relatively mature [3], [19]–[21]. Marconi et al. employed a Si-QD/SiO<sub>2</sub> multilayer to form a MOSLED structure and observed EL at a wavelength of around 880 nm [20]. EL at shorter wavelengths of 700–780 nm was also demonstrated with Si-QD embedded  $SiO<sub>x</sub>$  MOSLEDs [3]. However, only a few studies have emphasized obtaining the Si-QD-dependent PL emission at ultraviolet (UV) and blue regions [22], [23]. The difficulty arose by not only the precise control of Si-QD size and density but also the efficient pumping of small Si-QDs with a finite density of states [24]–[26]. Valenta et al. observed nearly coincident PL and EL spectra of an  $n^+$ -Si/SiO<sub>x</sub>/p<sup>+</sup>-Si diode [27]. Franzo *et al.* subsequently confirmed the luminescence caused by conduction-to-valence band recombination within a Si-QD [28]. Few current studies have addressed UV or blue EL emitted from any type of LED with buried Si-QDs.

This study preliminarily demonstrates blue and yellow EL emitted from MOSLEDs made on a Sirich  $SiO<sub>x</sub>$  film with buried Si-QDs. This approach is based on the precise control of quantumconfined Si-QD size by detuning the recipe of substrate temperature and  $N_2O/SiH_4$  fluence ratio during PECVD synthesis. The variation of the O/Si composition ratio in  $SiO<sub>x</sub>$  determines the quantity of excessive Si atoms in one  $SiO<sub>2</sub>$  unit cell so that the finite diffusion length of Si atoms in a Si-rich oxide matrix can be obtained. The excessive Si atom density and the diffusion length play important roles in confining the size Si-QDs after high-temperature annealing. Specifically, the variable diffusion lengths of Si atoms in Si-rich  $SiO_x$  films with changing O/Si composition ratios are correlated with the wavelength of EL produced by MOSLEDs.

## 2. Experimental Details

## 2.1. Preparation of Si-QD-Based Mosleds

The Si-rich SiO<sub>x</sub> films were grown on a (100)-oriented p-type Si wafer using PECVD at a substrate temperature, chamber pressure, and RF plasma power of 250  $^{\circ}$ C–550  $^{\circ}$ C with 100  $^{\circ}$ C increments, 10 Pa, and 60 W, respectively. The  $SH_4$  and  $N_2O$  gaseous mixtures with changing fluence ratios were used as the reactant gas to deposit Si-rich  $SiO_x$  films with different O/Si composition ratios. The O/Si composition ratio is a nonlinear function of  $N_2O/SiH_4$  fluence ratio. For





TABLE 1 O/Si composition ratio and Si-QD size versus the  $N<sub>2</sub>O/SiH<sub>4</sub>$  fluence ratio and substrate temperature



Fig. 1. (a) HRTEM image for the 1.7  $\pm$  0.2 nm large Si-QDs embedded in a SiO<sub>x</sub> sample with blue PL emission. (b) Size distribution of Si-QDs embedded in  $SiO<sub>x</sub>$  films with blue PL emission. Inset: the magnified HRTEM image of the 1.7-nm large Si-QD.

samples grown with the (a) recipe with O/Si composition ratio of 0.73–1.15, the flowing rates of SiH<sub>4</sub> and  $N_2O$  gases were 11 sccm and 75 sccm, respectively. For Si-rich SiO<sub>x</sub> samples grown with the (b) recipe with larger O/Si composition ratios of 1.02–1.37, the flowing rates of SiH<sub>4</sub> and N<sub>2</sub>O gases were 11 sccm and 50 sccm, respectively. To obtain the Si-rich SiO<sub>x</sub> samples with O/Si composition ratios of up to 1.42 or larger, the (c) recipe with  $N_2O$  gas fluence further decreased to 25 sccm at a constant SiH<sub>4</sub> fluence is used. This N<sub>2</sub>O decreasing recipe facilitates Si-rich SiO<sub>x</sub> growth for greater control on the size of Si-QDs. Table 1 lists the synthesizing recipes for all samples used in this work.

After deposition, the SiO<sub>x</sub> films were annealed at 1100 °C in flowing N<sub>2</sub> (250 sccm) and H<sub>2</sub> (5 sccm) mixed gases for 90 min to induce Si-QD precipitation in the  $SiO<sub>x</sub>$  matrix. The Si-QD size in Si-rich  $SiO<sub>x</sub>$  is determined by using the high-resolution transmission electron microscopy (HRTEM, JEOL 2010). By taking the SiO<sub>x</sub>:Si-QD film with blue PL as an example, the bright-field-view HRTEM image for the embedded Si-QDs is shown in Fig. 1(a). The size distribution of Si-QDs shown in Fig. 1(b) reveals that the Si-QD diameter is ranged between 1.2 and 2.2 nm, which exhibits a nearly Gaussian-like distribution function with its peak located at 1.7 nm and a full-width at halfmaximum of 0.2 nm. Under illumination with a He-Cd laser at a wavelength and average intensity of 325 nm and 5 W/cm<sup>2</sup>, respectively, the PL of Si-QDs buried in Si-rich SiO<sub>x</sub> films ranging between 300 and 900 nm was resolved using a monochromator (CVI, DK240) connected to a photomultiplier (Hamamatsu R928) and a digital multimeter (HP 34401A). X-ray photoelectron spectroscopy (XPS) with a Mg  $K_{\alpha}$  line illumination at 1253.6 eV was employed to analyze the O/Si composition ratio of Si-rich SiO<sub>x</sub> films. A 200-nm-thick ITO film with a contact diameter of 0.8 mm was sputtered on the top of a Si-rich SiO<sub>x</sub> surface, and a 500-nm-thick AI film was evaporated on the bottom of a p-Si substrate to form a MOSLED. The MOSLED was operated by a programmable electrometer (Keithley, model 6517) for EL characterization.



Fig. 2. Thickness of postannealed  $SiO<sub>x</sub>$  samples grown with different recipes and substrate temperatures.

## 3. Results and Discussion

#### 3.1. Thickness of Si-Rich SiO<sub>x</sub> Films

The slight reduction in Si-rich  $SiO_x$  film thickness after annealing is mainly attributed to dehydrogenation, which becomes significant with lengthened annealing duration. With the same deposition and annealing durations, the  $SiO_{0.75-1.14}$  film was thinner than the  $SiO_{1.02-1.37}$  and  $SiO_{1.42-1.62}$  films because of their increasing O/Si composition ratios. When increasing the deposition temperature under the same deposition time, the thickness of Si-rich  $SiO<sub>x</sub>$  films grown with recipes (a), (b), and (c) slightly increased from 270 nm to 300 nm, 310 nm to 335 nm, and 340 nm to 370 nm, respectively, as shown in Fig. 2.

The Si-rich SiO<sub>x</sub> thickness was also varied by changing the O/Si composition ratios because the sizes of Si and the O atom differentiate from each other significantly. Therefore, the thickness of Si-rich  $SiO<sub>x</sub>$  films is a function of the volume ratio between self-assembled Si-QDs and standard  $SiO<sub>2</sub>$  so that the volumes of excessive Si atoms and standard  $SiO<sub>2</sub>$  in one mole of a Si-rich  $SiO<sub>x</sub>$ film can be individually calculated to obtain the thickness of Si-rich  $SiO<sub>x</sub>$  films. For example, the Si-rich  $SiO<sub>x</sub>$  sample can be theoretically described as a function of self-assembled Si-QDs and standard  $SiO<sub>2</sub>$ 

$$
SiO_x \rightarrow \left(1-\frac{x}{2}\right)Si + \frac{x}{2}SiO_2. \tag{1}
$$

By setting the O/Si composition ratio at 1.62 for Si-rich  $SiO<sub>x</sub>$  films grown with recipe (c) at a substrate temperature of 550 °C, (1) shows that 0.19 moles  $(1 - x/2)$  of excessive Si atoms are buried in the 0.81 mole  $(x/2)$  of standard SiO<sub>2</sub> molecules. Assuming that the unit cell volumes of  $SiO<sub>2</sub>$  and Si-QD are 13.996 and 0.16 nm<sup>3</sup>, respectively, the volume of one mole Si-rich SiO<sub>1.62</sub> molecule is approximately 23.6  $cm<sup>3</sup>$ , as determined by

$$
V_{\text{mole}}(SiO_x) \approx V_{\text{mole}}(SiO_2) + V_{\text{mole}}(Si) \approx \frac{\frac{X}{2} \times N_0}{N_{SiO_2}} \times V_{SiO_2} + \frac{\left(1-\frac{X}{2}\right) \times N_0}{N_{Si}} \times V_{Si}. \hspace{1cm} (2)
$$

Equation (2) shows both the estimated volume and the thickness of one mole Si-rich SiO<sub>x</sub> film increase with the enlarging O/Si composition ratio (see Table 2). The standard variation of  $SiO<sub>x</sub>$ thickness increases with the decreasing  $\text{SiH}_4/\text{N}_2\text{O}$  fluence ratio because of the enhanced decomposition process under the same RF plasma power. This causes the Si-rich SiO<sub>x</sub> grown with recipe (c) to thicken more than the other samples grown with recipes (a) and (b) because of their higher O/Si composition ratios.

**Temperature** 

 $250^{\circ}$ C

 $(c): 2.3$ 

 $217$ 



#### TABLE 2

Volume of one mole of Si-rich SiO<sub>x</sub> molecule versus the N<sub>2</sub>O/SiH<sub>4</sub> fluence ratio and substrate temperature during PECVD growth

 $(a): 6.8$ 

 $173$ 

Volume of One mole  $SiO<sub>x</sub>$  molecules (cm<sup>3</sup>)

 $(b): 4.5$ 

19.3



Fig. 3. (a) The actual and normalized (inset) PL of Si-rich SiO<sub>x</sub> samples grown by changing the recipes (a) to (c) at a constant substrate temperature of 450 °C. (b) Top: PL emission patterns changing from red to blue corresponding to the change in average Si-QD size from 4.5 nm to 1.7 nm. Bottom: normalized PL intensity of a  $SiO<sub>x</sub>$  sample grown with different recipes and substrate temperatures [250 °C (black), 350 °C (red), 450 °C (green), 550 °C (blue)] as a function of O/Si composition ratio. (All samples are annealed at 1100  $^{\circ}$ C for 90 min.)

## 3.2. PL Analysis of Si-QDs Embedded in Si-Rich SiO<sub>x</sub> Films

Fig. 2 shows the PL spectra and related intensities of Si-rich  $SiO<sub>x</sub>$  samples grown at different substrate temperatures (from 250  $\degree$ C to 550  $\degree$ C at 100  $\degree$ C increment) and annealed at a constant duration of 90 min. At a constant substrate temperature, the PL peak wavelength exhibits a blueshifted phenomenon with the growth recipe of  $N_2O/SiH_4$  fluence ratio changing from (a) to (c), as Fig. 3(a) shows. As the O/Si composition increased from 1.42 to 1.62, the room-temperature PL peak wavelength of the Si-rich SiO<sub>x</sub> sample blue-shifts to 350–410 nm and increases its intensity from 65 to 70 counts/nm. This is mainly attributed to the quantum confinement effect of Si-QDs as their size shrinks to 1.7–2.1 nm. The blue PL of the Si-rich  $SiO_x$  sample grown with recipe (c) is distorted by a sharp cut at the wavelength of approximately 350 nm caused by the high-pass filter in front of a monochromator, as Fig. 3(a) shows. Specifically, the PL peak wavelength centered at 760 nm for the Si-rich SiO<sub>x</sub> sample grown with recipe (c) appears as an artificial signal because of the second-order harmonics of blue PL at 380 nm caused by a monochromator grating. Conversely, the PL peak wavelength further red-shifts to 580 nm with its linewidth greatly broadened to 200 nm as the O/Si composition ratio decreases to 1.37. The red PL of the Si-rich SiO<sub>x</sub> sample grown with recipe (a) is slightly distorted because of the rapidly decayed spectral response of photomultiplier tubes (PMTs) at a wavelength exceeding 800 nm. When keeping the  $N_2O/SiH_4$  fluence ratio constant, the Si-rich SiO<sub>x</sub> samples grown at substrate temperature decreasing from 550 °C to 250  $\degree$ C creates a red-shifted PL wavelength accompanied with a reduced PL peak intensity (normalized to Si-rich  $SiO<sub>x</sub>$  film thickness), as Fig. 3(b) shows.

Fig. 3(b) shows the actual PL intensity normalized to the  $SiO<sub>x</sub>$  thickness when comparing different samples because several samples are with slightly different thickness. Fig. 3(b) reveals that



Fig. 4. TRPL traces for Si-QD embedded Si-rich  $SiO<sub>x</sub>$  films grown with recipes (a), (b), and (c).

a low-temperature growth usually results in both the decreased O/Si composition ratio and the redshifted PL wavelength. Furthermore, the Si-rich  $SiO<sub>x</sub>$  sample grown at recipe (b) exhibits the O/Si composition ratio of 1.15–1.37 and the buried Si-QDs with size of 3.4–3.6 nm, which provides a white-yellow emission pattern with a weak PL peak intensity of 15–40 counts/nm. By further reducing the O/Si composition ratio to 0.73–1.15 with increasing  $N_2O$  gas fluence to 75 sccm, the 4.0–4.5-nm large Si-QDs contributed to an NIR PL with its peak wavelength and intensity redshifted to 760–780 nm and decayed to 5–15 counts/nm, respectively. HRTEM images and PL patterns confirm that the PL color of the Si-rich  $SiO<sub>x</sub>$  film blue-shifts (red, orange, yellow, and blue) as the Si-QD size decreases (4.5  $\pm$  0.3, 3.5  $\pm$  0.2, 2.8  $\pm$  0.3, and 2.1  $\pm$  0.2 nm), which is relatively in good agreement with the theoretical prediction [29]. Although the  $SiO<sub>x</sub>$  samples have the identical O/Si composition ratio, the stronger PL of  $SiO<sub>x</sub>$  samples appears under higher  $N_2O/SiH_4$  fluence ratio and lower substrate temperature. Because the SiO<sub>x</sub> samples become thicker when growing with lower  $N_2O/SiH_4$  fluence ratio and lower substrate temperature, in which more Si-QDs can be precipitated to serve as the radiative luminescence center.

The invariant PL power from each Si-rich  $Si\alpha_s$  sample during a long-term test corroborates the stability of embedded Si-QDs. The luminescence contributed by most of the radiative defects can be ignored, including the weak oxygen bond (WOB, O–O) at 415 nm, the neutral oxygen vacancy (NOV,  $O_3 \equiv Si-Si \equiv O_3$ ) at 455 nm, the  $E_\delta'$  (Si  $\uparrow$  Si  $\equiv$  Si) at 520 m, and the nonbridged oxygen hole center (NOBHC,  $O_3 \equiv Si-O$ ) at 630 nm [30]–[32]. In addition, the lack of nitrogen in a Si-rich  $SiO<sub>x</sub>$  matrix confirmed by XPS analysis also excludes other radiative recombination mechanisms, such as the Si–N bond transition at 428–477 nm, and the conduction-band to  $\mathsf{N}_2^0$ -level (hole trap) or the N $_4^+$ -level to valence-band (electron trap) transition at 413 nm [33], [34]. In general, these defects remain their PL wavelengths as constant and never vary with different  $SiO<sub>x</sub>$  sample recipes. In this case, the shift of PL wavelength (from PL analysis) in samples with different Si-QD size (from HRTEM analysis) is observed, and the proportionality between the PL wavelength and Si-QD is coincident with the theoretical simulation. The PL peak wavelengths of the samples are different from those of the radiative defects. These experimental results elucidate that even the radiative defects may have contributed to the PL, which is less compared with the contribution of the Si-QDs as the peak intensities of the defect-related PL component is insufficiently large to affect the whole PL spectrum. In addition, the blue PL exhibits a central peak wavelength at 380 nm in our case, which is apparently different from those wavelengths of the PL contributed by the radiative defects. From the average Si-QD size of 1.7  $\pm$  0.2 nm, we can also correlate the PL peak wavelength and corresponding spectral linewidth with an empirical formula of  $\lambda = 1.24/(1.12 + 5.83/d^{1.78})$  set for the Si-QDs buried in the  $SiO_x$  matrix.

## 3.3. TRPL Analysis of Si-QDs Embedded in Si-Rich SiO<sub>x</sub> Films

The other evidence to distinguish the luminescence of Si-QDs from those of the radiative defects is the luminescent lifetime diagnosis of Si-QDs obtained from the time-resolved PL (TRPL) analysis. Fig. 4 shows the TRPL traces of three Si-rich  $SiO<sub>x</sub>$  samples pumped by a 10-ns gain-switched GaN laser diode pulse at 405 nm, which is detected by a PMT (Hamamatsu R928) with a switching response of 2–3 ns. In principle, the PL lifetime of the Si-QD embedded in the  $SiO<sub>x</sub>$  sample is strongly correlated with the radiative recombination rate of the electron–hole pairs inside the Si-QD.

Moreover, the Si-QD size distribution affects the dispersion factor  $(\beta)$  of the stretched exponential decay function in the fitting equation  $\mathsf{I}(t)=\mathit{l}_0$ exp $[-(t/\tau_\mathsf{PL})^\beta]$ .  $\tau_{\mathsf{PL}}$  denotes PL lifetime, and  $\beta$  denotes dispersion factor correlated with Si-QD size distribution. With increasing Si-QD size distribution, the PL spectra become broadened and result in the enhanced lifetime dispersion phenomenon. Therefore, the broadened Si-QD size distribution contributes to the smaller dispersion factor  $(\beta)$ . In this case, the single-stage stretched retention of the TRPL with a lower dispersion factor is employed to fit the oscilloscope trace of the time-resolved response of the full PL spectrum. The exponential decay constants of Si-QD-related luminescence for Si-rich  $SiO_x$  samples grown with recipes (a), (b), and (c) are determined as 10.5  $\mu$ s, 2.25  $\mu$ s, and 67 ns, respectively. The PL decay time constant is decreased by more than two orders of magnitude when the Si-QD size shrinks from 4.5 nm to  $\leq$  1.7 nm. The dispersion factor of Si-QD-related luminescence for Si-rich SiO<sub>x</sub> samples grown with recipes (a), (b), and (c) are determined as 0.51, 0.62, and 0.73, respectively. By shrinking the Si-QD size of smaller than 8 nm, the enhanced momentum overlapping can be observed with further decreasing Si-QD size due to an increasing probability of nonphonon-assisted electron–hole recombination. A faster nonphonon-assisted electron–hole PL lifetime is mainly attributed to a larger momentum overlap between the electron and hole wave function in smaller Si-QDs, thus providing a higher recombination rate through the enhanced probability of direct transition [35]. We find that the lifetimes are consistent with electron–hole wave function overlap and momentum uncertainty.

The lifetime is usually contributed by nonradiative, radiative, and Auger recombinations. However, the TRPL can only determine the radiative lifetime and cannot be correlated with the nonradiative recombination process. During TRPL analysis, the radiative lifetimes of PL contributed by radiative defects and Si-QDs are different from each other. By decomposing the TRPL traces, we can distinguish the lifetime of radiative defects in the Si-rich  $SiO<sub>x</sub>$  film from that of the Si-QDs. The defect-related luminescence has been deconvoluted from the full-band TRPL traces shown in Fig. 4. Previously, the blue-band PL emission at 400–500 nm with lifetime of 5–37 ns due to oxygen-related defects at the  $Si-QD/SiO<sub>x</sub>$  interface was reported by Pi et al. and de Boer et al. [36], [37]. Moreover, the PL lifetime of  $\mathsf{E}_\delta'$  defect-related luminescence around 10 ns was reported by Nishikawa et al. [38]. In comparison, the PL lifetime of 67 ns obtained from the blue-color PL is much longer than those of defect-related luminescence. Therefore, the contribution of radiative defects has been taken into account, and the TRPL lifetimes reported in our work are not attributed to those defects [39]. The contribution of nonradiative defects is currently left as an unknown issue. This evidence further supports the direct radiative recombination in Si-QD core with a strong quantum confinement effect. For Si-rich  $SiO<sub>x</sub>$  samples with buried Si-QDs at different sizes, the strengthened PL intensity is observed and correlated with the reduced PL decay time constant. This result is straightforward by considering the PL intensity described as a function of the PL decay time and volume of Si-QDs [40]

$$
I_{\rm PL} = \eta \sigma N_{\rm Si\text{-}QD} \Phi(t) \frac{\tau_{\rm pump}}{\tau_{\rm PL}} \propto N_{\rm Si\text{-}QD} \Phi(t) \tau_{\rm PL}^{-1}
$$
(3)

where  $I_{PL}$  is the Si-QD-related PL intensity,  $\eta$  is a relative coefficient,  $\sigma$  is an emission (absorption) cross section of radiative Si-QD,  $\Phi(t)$  is the pumping flux density,  $\tau_{\mathsf{pump}}$  is the lifetime of pumping source,  $\tau_{PL}$  is the radiative lifetime of Si-QD, and N<sub>Si-QD</sub> is the volume density of recombination centers in Si-QDs.

## 3.4. XPS of Si-QDs Embedded in Si-Rich  $SiO<sub>x</sub>$  Films

The XPS analysis in Fig. 5 shows a linear variation of Si–O-related electron energy peak moving toward higher binding energy, indicating that the Si-rich  $SiO<sub>x</sub>$  matrix gradually changes to a pure  $SiO<sub>2</sub>$  structure after increasing the deposition temperature and N<sub>2</sub>O fluence during growth [41]–[43]. The fitted XPS binding energy from 99.7 eV to 103.35 eV is attributed to the  $Si_{(2p)}$  electrons from Si– Si bonds and O–Si–O bonds, respectively [44]. In this case, the XPS-binding peak of the Si-rich  $SiO<sub>x</sub>$  film with buried Si-QDs is caused by varying the O/Si composition ratio of the Si-rich SiO<sub>x</sub>



Fig. 5. XPS spectral peak energy of electrons from Si<sub>sp</sub> core level versus the O/Si composition ratio of the SiO<sub>x</sub> film grown by detuning the N<sub>2</sub>O/SiH<sub>4</sub> fluence ratio and substrate temperature [250 °C (black), 350 °C (red), 450 °C (green), 550 °C (blue)].

matrix. Compared with  $SiO<sub>1.42-1.62</sub>$  and  $SiO<sub>1.15-1.37</sub>$  samples, a critical discrepancy between the  $N_2$ O fluence and total fluence occurs to vary the O/Si composition ratio of the Si-rich SiO<sub>x</sub> film. The decomposition of a Si atom from SiH<sub>4</sub> is generally easier than that of an oxygen atom from  $N_2O$ because the desorption energies of  $N_2O$  and  $SiH_4$  are 101.5 kcal/mol and 75.6 kcal/mol, respectively. When the PECVD chamber is filled with a lower molecule density of  $N_2O/SiH_4$  mixture, the RF energy for each molecule given by the plasma in the PECVD chamber is high enough that the decomposition rates of O and Si atoms do not deviate significantly from each other. As the molecule density of the  $N_2O/SiH_4$  mixture doubles by enlarging the reactant gas fluence, the reduced plasma energy on each molecule forces the decomposition of O and Si atoms to differentiate from each other. As a result, the anomalous growth of the Si-rich SiO<sub>x</sub> film causes a higher excessive Si content. Thus, the O/Si composition ratio can be linearly detuned by changing the total N<sub>2</sub>O/SiH<sub>4</sub> mixture fluence instead of the adjustment on the N<sub>2</sub>O/SiH<sub>4</sub> fluence ratio. The increasing O/Si composition ratio inevitably results in a linear blue-shift of  $Si_{2p}$  X-ray photoelectron peak energy. Specifically, the deconvoluted peak intensity of Si-Si bonds within the  $Si<sub>2p</sub>$  electronrelated XPS signal also shows a decreasing trend as the O/Si composition ratio increases. This corroborates the reduced Si-QD density in a Si-rich  $SiO<sub>x</sub>$  film grown with a lower N<sub>2</sub>O/SiH<sub>4</sub> fluence ratio at a higher substrate temperature.

## 3.5. Voltage–Current Analysis of Si-QD-Based MOSLEDs

In addition, the voltage–current  $(V-I)$  analysis of Si-rich SiO<sub>x</sub> MOSLEDs grown with different recipes and substrate temperatures is performed. The turn-on voltage and current defines the voltage required to turn on the tunneling the current and the EL emission, respectively. The V–I analysis shows that the threshold current for turning on MOSLEDs grown with recipe (b) is greatly reduced from 110  $\mu$ A to 40  $\mu$ A (see Fig. 6). The *V–I* plots illustrate the correlation of turn-on voltage and turn-on current of Si-rich  $SiO<sub>x</sub>$  MOSLEDs with the O/Si composition ratio of the Si-QD-doped  $SiO<sub>x</sub>$  film, as shown in the lower and upper part of Fig. 6. In Fig. 6, the short-dashed line with same color links the data of samples grown with a constant  $N_2O/SiH_4$  fluence ratio, and the different patterns with same color indicate that the samples were grown at identical temperature. Consequently, the corresponding turn-on voltage of MOSLEDs grown with recipe (b) inevitably increases from 60 V to 140 V by enhancing the substrate temperature from 250  $\degree$ C to 550  $\degree$ C. In addition, the turn-on voltage and current of the MOSLEDs grown with recipes (a) and (c) have the same trends. In contrast, the Si-rich  $Si\chi$ -based MOSLEDs grown with recipe (a) exhibit the relatively low turn-on voltage and high injection current because of their extremely low O/Si composition ratios. However, the current passing across  $SiO<sub>x</sub>$ -based MOSLEDs grown with recipe (c) decays significantly from 60  $\mu$ A to 20  $\mu$ A and is accompanied with increasing compliance voltages from 120 V to 225 V. The turn-on voltage and turn-on current of devices are dominated by the O/Si composition ratio in the Si-rich SiO<sub>x</sub> film. The Si-rich SiO<sub>x</sub> film gradually increases its



Fig. 6. Turn-on current (upper) and turn-on voltage (lower) of the  $SiO<sub>x</sub>$  MOSLED versus the O/Si composition ratio of the Si-rich SiO<sub>x</sub> film grown by detuning the N<sub>2</sub>O/SiH<sub>4</sub> fluence ratio as 6.8 (triangle patterns linked with red short-dashed line), 4.5 (square patterns linked with green short-dashed line), and 2.3 (circle patterns linked with blue short-dashed line) and by changing the substrate temperature as 250 °C (black patterns), 350 °C (red patterns), 450 °C (green patterns), and 550 °C (blue patterns).



Fig. 7. (a) The P–I curves and (b) actual and normalized (inset) EL spectra of  $SiO<sub>x</sub>$ -based MOSLEDs grown with recipes (a), (b), and (c) at substrate temperature 450  $\degree$ C.

resistance with increasing the O/Si composition ratio because the Si-rich SiO<sub>x</sub> film gradually becomes stoichiometric to approach a pure  $SiO<sub>2</sub>$  matrix. A higher forward-biased condition is required to inject the carriers through the Si-rich  $SiO<sub>x</sub>$  film with a higher O/Si composition ratio, and the tunneling-based carrier transport mechanism is dominated by exponential-like V-I behavior. In addition, the Si-rich SiO<sub>x</sub> film with higher O/Si composition ratio easily contributes to smaller Si-QDs [45]. A smaller Si-QD decreases the effective dielectric constant and enhances the barrier height of F-N tunneling to degrade overall tunneling probability [46]. Therefore, the MOSLEDs decrease its turn-on current with increasing the O/Si composition ratio.

## 3.6. EL and Power–Current Analyses of Si-QD-Based MOSLEDs

The maximum output power of MOSLEDs made by  $SiO<sub>x</sub>$  grown with recipes (a) and (b) are both up to 400 nW, and the corresponding P/I slopes are 0.63 mW/A and 0.91 mW/A, respectively [see Fig. 7(a)]. In addition, the SiO<sub>1.54</sub>-based MOSLED has an EL power of 100 nW at 40  $\mu$ A with a P/I slope of 1.01 mW/A. The EL emission power of the Si-QD-doped  $SiO<sub>x</sub>$  MOSLEDs grown with changing the recipe from (a) to (c) degrades because the carriers are hardly injected through the Si-rich SiO<sub>x</sub> film with a higher O/Si composition ratio. Another reason is the limited ITO transmittance



Fig. 8. PL (dashed line) and EL (solid line) spectra of Si-QD-based MOSLEDs with blue, yellow, and red color emissions grown with recipes (a), (b), and (c).

of 40%–50% in the blue region. The EQE is defined as the ratio of the output photon number and input electron number described as

$$
\eta_{\text{ext}} = \int_{t_0}^{t_1} \frac{P(t)}{I(t)} \frac{e}{h\nu} dt = \frac{P_{\text{opt}}}{I_{\text{bias}}(1.24/\lambda)}
$$
(4)

where  $P_{oot}$  defines the optical output power,  $\lambda$  defines the EL wavelength, and  $I_{bias}$  defines the bias current. The EQE of devices grown with recipe (a) increases from  $5.4 \times 10^{-3}$ % to  $2.5 \times 10^{-2}$ % with increasing substrate temperature from 250 °C to 550 °C. In addition, the EQE of devices grown with recipes (b) and (c) has same trends. The EQE of devices grown with recipes (b) and (c) enhances from  $8.9 \times 10^{-3}$ % to  $3.6 \times 10^{-2}$ % and from  $3.2 \times 10^{-3}$ % to  $5.2 \times 10^{-2}$ %, respectively. The enhancement of EQE is mainly attributed to the stronger quantum confinement in smaller Si-QD. The stronger quantum confinement contributes to carriers confined in smaller Si-QDs to promote a higher probability of recombination. Therefore, the smaller Si-QD embedded  $SiO_x$ -based MOSLEDs have a higher EQE. Fig. 7(b) shows the EL spectra of MOSLEDs grown with the growth recipes changing from (a) to (c) at the biased electric field of 3.3 MV/cm, 6.7 MV/cm, and 8.8 MV/cm, respectively. It shows the primary EL peak wavelengths at 400–450 nm for blue MOSLEDs, at 600–650 nm for yellow MOSLEDs, and at 700–750 nm for red MOSLEDs. The blue-shifted phenomenon appears because the embedded Si-QD size decreases when increasing the O/Si composition ratio in the Si-rich SiO<sub>x</sub> film. In addition, the EL spectra of MOSLEDs grown with recipes (a) and (b) are stronger and broader than those of the devices made with recipe (c) because the Si-rich  $SiO_x$  film with lower O/Si composition ratio has a broader Si-QD size distribution. Therefore, the devices with lower O/Si composition ratio easily have the secondary EL peak wavelengths due to the emission of small Si-QDs.

The EL and PL spectra of three different samples have been compared in Fig. 8. The central emission wavelength of EL and PL for all samples is not deviated from each other; however, there are some specific peaks observed in the EL spectrum. For the MOSLED with blue EL/PL, the EL peak wavelength centered at 380 nm is similar to PL peak wavelength. The sharpened PL at the short-wavelength side is due to the notch filter set for blocking the pumping laser spectrum. However, the EL linewidth of 126 nm is more broadened than the PL linewidth of 71 nm because another  $\mathsf{E}_\delta'$  defect (the precursor of Si-QDs) is also pumped under high biases [47]. In addition, the excessive carriers under a highly biased condition also excite the EL from larger Si-QDs with smaller volume density. Such a spectral broadening phenomenon has ever been discussed in previous work [47]. Similar trend is also observed in the other two samples. The yellow EL from the MOSLED driven at relatively high bias broadens its linewidth to 258 nm. A secondary EL peak decomposed with its peak located at around 400–405 nm results from the excitation of WOB defects in the incomplete  $SiO<sub>x</sub>$  matrix with a small O/Si composition ratio. The highly biased



Fig. 9. EL emission patterns and biased field of  $SiO_x$ -based MOSLEDs grown at 450 °C with recipe (a) in upper row, (b) in middle row, and (c) in lower row.



Fig. 10. UV EL patterns of  $SiO<sub>1.62</sub>$ -based MOSLEDs grown with recipe (c) at a substrate temperature of 550 °C (biased at 225–265 V from left to right at a 10-V increment).

operation also excites the larger Si-QDs in the yellow-EL MOSLED, whereas the weak  $\mathsf{E}_\delta'$  defect related radiation maybe covered by the broadband Si-QD-related EL spectrum. In the MOSLED with red-EL spectrum, the secondary EL peak at nearly 500–520 nm originated from the  $\mathsf{E}_\delta'$  defects is also observable in the PL spectrum [38].

The Si-QD-doped SiO<sub>x</sub> MOSLEDs grown at a substrate temperature of 450 °C and the recipes (a), (b), and (c) show EL patterns in the upper, middle, and lower rows of Fig. 9, respectively. The EL pattern of the MOSLEDs changes its color from red to blue with increasing the  $N_2O/SiH_4$  fluence ratio because the Si-QD size decreases by increasing the O/Si composition ratio. To obtain sufficiently bright blue-EL emission patterns (immediately above the soft-breakdown regime), the forward bias was increased from 8.5 MV/cm to 9.1 MV/cm to elucidate that a higher turn-on voltage is required to tunnel the carriers through a Si-rich  $Si\sigma_x$  film with a high O/Si composition ratio. In comparison, the EL patterns of devices grown with recipe (b) at a forward bias changing from 6.4 MV/cm to 7.0 MV/cm. In these MOSLEDs, the EL patterns become from dark yellow to bright yellow as the biased electrical field from 6.4 MV/cm to 6.7 MV/cm or larger. However, the yellow devices show an obvious soft-breakdown phenomenon at  $>$  7.0 MV/cm, which breaks down soon after a few minutes (1–3 min) when biasing above 7.5 MV/cm. A similar development of EL patterns from dark red to orange appears in the  $SiO<sub>x</sub>$ -based MOSLEDs grown with recipe (a). Nevertheless, their bias electric field from 3.0 MV/cm to 3.6 MV/cm is substantially lower than those of the MOSLEDs made by other  $SiO<sub>x</sub>$  samples with higher O/Si composition ratios. This phenomenon is corroborated by the decreased oxide resistance when growing the Si-rich SiO<sub>x</sub> film with larger Si-QD size and lower O/Si composition ratio.

At last, Fig. 10 shows the UV EL patterns of  $SiO<sub>1.62</sub>$ -based MOSLEDs at a biasing voltage of up to 225–265 V. However, the top ITO contact degrades its transparency at UV-blue wavelengths to attenuate the EL power of UV MOSLEDs. The maximum EL power of a  $SiO<sub>1.62</sub>$ -based UV MOSLED is greatly attenuated to 35-40 nW when the ITO transmittance further decays to  $<$  30% when the EL wavelength blue-shifts to 375 nm or shorter [48]. The device lifetime is only 1–3 min for emitting EL at UV wavelength. The EL spectra of  $SiO<sub>1.62</sub>$ -based MOSLEDs at different biases are unavailable because the overheating problem of UV MOSLEDs becomes severe at extremely high biases.

## 4. Conclusion

This study investigates the transient UV and visible luminescent dynamics of  $SiO_x$ -based MOSLEDs with different O/Si composition ratios grown by changing the  $N_2O/SiH_4$  fluence ratio and substrate temperatures. The size and luminescent wavelength of Si-QDs can be controlled by adjusting the O/Si composition ratio precisely. By decreasing the N<sub>2</sub>O fluence and the N<sub>2</sub>O/SiH<sub>4</sub> fluence ratio from 75 to 25 sccm and from 6.8 to 2.3, respectively, TEM and XPS analyses reveal that the Si-QD size decreases from 4.5 to 1.7 nm as the O/Si composition ratio increases from 1.15 to 1.62. Substrate heating facilitates the growth of a Si-rich  $SiO<sub>x</sub>$  film with a larger O/Si composition ratio, thus providing a small distribution range for Si-QD size and its PL wavelength. By shrinking the average size of Si-QDs from 4.5 to  $<$  1.7 nm, the TRPL lifetime decay constant is significantly shortened from 11.5  $\mu$ s to 67 ns. This is because of the larger overlap between electron and hole wave functions, which essentially leads to a faster nonphonon-assisted carrier recombination rate in smaller Si-QDs. Observations of both the continuous-wave and transient PL show that the direct radiative recombination process in Si-QD is the result of the strong quantum confinement effect rather than the oxygen-related defects. With the O/Si composition ratio between 0.8 and 1.15, the red EL pattern can be obtained from  $SiO<sub>x</sub>$ -based MOSLEDs with relatively low turn-on voltage and high injecting current. The turn-on electric field of MOSLEDs inevitably increases when they are fabricated on a Si-rich SiO<sub>x</sub> film with a larger O/Si composition ratio. The small Si-QDs with blueshifted PL are obtained as the Si-rich  $SiO<sub>x</sub>$  film gradually becomes stoichiometric to approach the standard SiO<sub>2</sub>. Nevertheless, the density of Si-QDs decreases, and the F-N carrier tunneling threshold increases accordingly. To obtain a UV or blue EL from MOSLEDs made on a Si-rich  $SiO<sub>x</sub>$ film with  $x > 1.4$ , the turn-on voltage and biased current are greatly increased to 225 V and decreased to 20  $\mu$ A, respectively. To shorten the EL wavelength from 780 to < 370 nm, the turn-on electric field incredibly increases from 3 to 9 MV/cm or higher. The maximum EL powers of 100 nW, 380 nW, and 400 nW for MOSLEDs with blue-, yellow-, and red-emission, respectively, are obtained with P/I slopes of 1.01 mW/A, 0.91 mW/A, and 0.63 mW/A, respectively. With a bias voltage of up to 225–265 V, the UV EL pattern can be observed from MOSLEDs made by  $SiO<sub>1.62</sub>$ with a buried Si-QD size of  $< 1.7$  nm. However, the EL power greatly attenuates to 40 nW or less because of the greatly attenuated transmittance of the top transparent ITO contact at UV wavelengths. The device lifetime is only 1–3 min for emitting EL at the UV wavelength. In the future, a multicolor MOSLED-based light source could be used in applications such as an alternative transmitter in an all-Si-based optical interconnect to improve the chip-to-chip transmission performance of Si-based microelectronics. With further enhancement on IQE and EQE, this type of light source could be used in other applications in Si-based microdisplay and microillumination chips.

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